

GROMACS

Groningen Machine for Chemical Simulations



USER MANUAL

Version 3.3

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More information can be found on our website: www.gromacs.org.

Preface & Disclaimer

This manual is not complete and has no pretention to be so due to lack of time of the contributors – our first priority is to improve the software. It is worked on continuously, which in some cases might mean the information is not entirely correct.

Comments are welcome, please send them by e-mail to gromacs@gromacs.org, or to one of the mailing lists (see www.gromacs.org).

We try to release an updated version of the manual whenever we release a new version of the software, so in general it is a good idea to use a manual with the same major and minor release number as your GROMACS installation. Any revision numbers (like 3.1.1) are however independent, to make it possible to implement bugfixes and manual improvements if necessary.

Online Resources

You can find more documentation and other material at our homepage www.gromacs.org. Among other things there is an online reference, several GROMACS mailing lists with archives and contributed topologies/force fields.

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However, we prefer that you cite (some of) the GROMACS papers [1, 2, 3, 4] when you publish your results. Any future development depends on academic research grants, since the package is distributed as free software!

Current development

Gromacs is a joint effort, with contributions from lots of developers around the world. The core development is currently taking place at

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GROMACS is *Free Software*

The entire GROMACS package is available under the GNU General Public License. This means it's free as in free speech, not just that you can use it without paying us money. For details, check the COPYING file in the source code or consult www.gnu.org/copyleft/gpl.html.

The GROMACS source code and selected set of binary packages are available on our homepage, www.gromacs.org. Have fun.

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Chapter 1

Introduction

1.1 Computational Chemistry and Molecular Modeling

GROMACS is an engine to perform molecular dynamics simulations and energy minimization. These are two of the many techniques that belong to the realm of computational chemistry and molecular modeling. *Computational Chemistry* is just a name to indicate the use of computational techniques in chemistry, ranging from quantum mechanics of molecules to dynamics of large complex molecular aggregates. *Molecular modeling* indicates the general process of describing complex chemical systems in terms of a realistic atomic model, with the aim to understand and predict macroscopic properties based on detailed knowledge on an atomic scale. Often molecular modeling is used to design new materials, for which the accurate prediction of physical properties of realistic systems is required.

Macroscopic physical properties can be distinguished in (a) *static equilibrium properties*, such as the binding constant of an inhibitor to an enzyme, the average potential energy of a system, or the radial distribution function in a liquid, and (b) *dynamic or non-equilibrium properties*, such as the viscosity of a liquid, diffusion processes in membranes, the dynamics of phase changes, reaction kinetics, or the dynamics of defects in crystals. The choice of technique depends on the question asked and on the feasibility of the method to yield reliable results at the present state of the art. Ideally, the (relativistic) time-dependent Schrödinger equation describes the properties of molecular systems with high accuracy, but anything more complex than the equilibrium state of a few atoms cannot be handled at this *ab initio* level. Thus approximations are necessary; the higher the complexity of a system and the longer the time span of the processes of interest is, the more severe the required approximations are. At a certain point (reached very much earlier than one would wish) the *ab initio* approach must be augmented or replaced by *empirical* parameterization of the model used. Where simulations based on physical principles of atomic interactions still fail due to the complexity of the system molecular modeling is based entirely on a similarity analysis of known structural and chemical data. The QSAR methods (Quantitative Structure-Activity Relations) and many homology-based protein structure predictions belong to the latter category.

Macroscopic properties are always ensemble averages over a representative statistical ensemble

(either equilibrium or non-equilibrium) of molecular systems. For molecular modeling this has two important consequences:

- The knowledge of a single structure, even if it is the structure of the global energy minimum, is not sufficient. It is necessary to generate a representative ensemble at a given temperature, in order to compute macroscopic properties. But this is not enough to compute thermodynamic equilibrium properties that are based on free energies, such as phase equilibria, binding constants, solubilities, relative stability of molecular conformations, etc. The computation of free energies and thermodynamic potentials requires special extensions of molecular simulation techniques.
- While molecular simulations in principle provide atomic details of the structures and motions, such details are often not relevant for the macroscopic properties of interest. This opens the way to simplify the description of interactions and average over irrelevant details. The science of statistical mechanics provides the theoretical framework for such simplifications. There is a hierarchy of methods ranging from considering groups of atoms as one unit, describing motion in a reduced number of collective coordinates, averaging over solvent molecules with potentials of mean force combined with stochastic dynamics [5], to *mesoscopic dynamics* describing densities rather than atoms and fluxes as response to thermodynamic gradients rather than velocities or accelerations as response to forces [6].

For the generation of a representative equilibrium ensemble two methods are available: (a) *Monte Carlo simulations* and (b) *Molecular Dynamics simulations*. For the generation of non-equilibrium ensembles and for the analysis of dynamic events, only the second method is appropriate. While Monte Carlo simulations are more simple than MD (they do not require the computation of forces), they do not yield significantly better statistics than MD in a given amount of computer time. Therefore MD is the more universal technique. If a starting configuration is very far from equilibrium, the forces may be excessively large and the MD simulation may fail. In those cases a robust *energy minimization* is required. Another reason to perform an energy minimization is the removal of all kinetic energy from the system: if several 'snapshots' from dynamic simulations must be compared, energy minimization reduces the thermal noise in the structures and potential energies, so that they can be compared better.

1.2 Molecular Dynamics Simulations

MD simulations solve Newton's equations of motion for a system of N interacting atoms:

$$m_i \frac{\partial^2 \mathbf{r}_i}{\partial t^2} = \mathbf{F}_i, \quad i = 1 \dots N. \quad (1.1)$$

The forces are the negative derivatives of a potential function $V(\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_N)$:

$$\mathbf{F}_i = -\frac{\partial V}{\partial \mathbf{r}_i} \quad (1.2)$$

The equations are solved simultaneously in small time steps. The system is followed for some time, taking care that the temperature and pressure remain at the required values, and the coordinates are written to an output file at regular intervals. The coordinates as a function of time

type of bond	type of vibration	wavenumber (cm ⁻¹)
C-H, O-H, N-H	stretch	3000–3500
C=C, C=O,	stretch	1700–2000
HOH	bending	1600
C-C	stretch	1400–1600
H ₂ CX	sciss, rock	1000–1500
CCC	bending	800–1000
O-H···O	libration	400– 700
O-H···O	stretch	50– 200

Table 1.1: Typical vibrational frequencies (wavenumbers) in molecules and hydrogen-bonded liquids. Compare $kT/h = 200 \text{ cm}^{-1}$ at 300 K.

represent a *trajectory* of the system. After initial changes, the system will usually reach an *equilibrium state*. By averaging over an equilibrium trajectory many macroscopic properties can be extracted from the output file.

It is useful at this point to consider the limitations of MD simulations. The user should be aware of those limitations and always perform checks on known experimental properties to assess the accuracy of the simulation. We list the approximations below.

The simulations are classical

Using Newton's equation of motion automatically implies the use of *classical mechanics* to describe the motion of atoms. This is all right for most atoms at normal temperatures, but there are exceptions. Hydrogen atoms are quite light and the motion of protons is sometimes of essential quantum mechanical character. For example, a proton may *tunnel* through a potential barrier in the course of a transfer over a hydrogen bond. Such processes cannot be properly treated by classical dynamics! Helium liquid at low temperature is another example where classical mechanics breaks down. While helium may not deeply concern us, the high frequency vibrations of covalent bonds should make us worry! The statistical mechanics of a classical harmonic oscillator differs appreciably from that of a real quantum oscillator, when the resonance frequency ν approximates or exceeds $k_B T/h$. Now at room temperature the wavenumber $\sigma = 1/\lambda = \nu/c$ at which $h\nu = k_B T$ is approximately 200 cm^{-1} . Thus all frequencies higher than, say, 100 cm^{-1} may misbehave in classical simulations. This means that practically all bond and bond-angle vibrations are suspect, and even hydrogen-bonded motions as translational or librational H-bond vibrations are beyond the classical limit (see Table 1.1). What can we do?

Well, apart from real quantum-dynamical simulations, we can do one of two things:

(a) If we perform MD simulations using harmonic oscillators for bonds, we should make corrections to the total internal energy $U = E_{kin} + E_{pot}$ and specific heat C_V (and to entropy S and free energy A or G if those are calculated). The corrections to the energy and specific heat of a one-dimensional oscillator with frequency ν are: [7]

$$U^{QM} = U^{cl} + kT \left(\frac{1}{2}x - 1 + \frac{x}{e^x - 1} \right) \quad (1.3)$$

$$C_V^{QM} = C_V^{cl} + k \left(\frac{x^2 e^x}{(e^x - 1)^2} - 1 \right), \quad (1.4)$$

where $x = h\nu/kT$. The classical oscillator absorbs too much energy (kT), while the high-frequency quantum oscillator is in its ground state at the zero-point energy level of $\frac{1}{2}h\nu$.

(b) We can treat the bonds (and bond angles) as *constraints* in the equation of motion. The rationale behind this is that a quantum oscillator in its ground state resembles a constrained bond more closely than a classical oscillator. A good practical reason for this choice is that the algorithm can use larger time steps when the highest frequencies are removed. In practice the time step can be made four times as large when bonds are constrained than when they are oscillators [8]. GROMACS has this option for the bonds and bond angles. The flexibility of the latter is rather essential to allow for the realistic motion and coverage of configurational space [8].

Electrons are in the ground state

In MD we use a *conservative* force field that is a function of the positions of atoms only. This means that the electronic motions are not considered: the electrons are supposed to adjust their dynamics instantly when the atomic positions change (the *Born-Oppenheimer* approximation), and remain in their ground state. This is really all right, almost always. But of course, electron transfer processes and electronically excited states can not be treated. Neither can chemical reactions be treated properly, but there are other reasons to shy away from reactions for the time being.

Force fields are approximate

Force fields provide the forces. They are not really a part of the simulation method and their parameters can be user-modified as the need arises or knowledge improves. But the form of the forces that can be used in a particular program is subject to limitations. The force field that is incorporated in GROMACS is described in Chapter 4. In the present version the force field is pair-additive (apart from long-range coulomb forces), it cannot incorporate polarizabilities, and it does not contain fine-tuning of bonded interactions. This urges the inclusion of some limitations in this list below. For the rest it is quite useful and fairly reliable for bio macro-molecules in aqueous solution!

The force field is pair-additive

This means that all *non-bonded* forces result from the sum of non-bonded pair interactions. Non pair-additive interactions, the most important example of which is interaction through atomic polarizability, are represented by *effective pair potentials*. Only average non pair-additive contributions are incorporated. This also means that the pair interactions are not pure, *i.e.*, they are not valid for isolated pairs or for situations that differ appreciably from the test systems on which the models were parameterized. In fact, the effective pair potentials are not that bad in practice. But the omission of polarizability also means that electrons in atoms do not provide a dielectric constant as they should. For example, real liquid alkanes have a dielectric constant of slightly more than 2, which reduce the long-range electrostatic interaction between (partial) charges. Thus the simulations will exaggerate the long-range Coulomb terms. Luckily, the next item compensates this effect a bit.

Long-range interactions are cutoff

In this version GROMACS always uses a cutoff radius for the Lennard-Jones interactions

and sometimes for the Coulomb interactions as well. Due to the minimum-image convention (only one image of each particle in the periodic boundary conditions is considered for a pair interaction), the cutoff range can not exceed half the box size. That is still pretty big for large systems, and trouble is only expected for systems containing charged particles. But then truly bad things can happen, like accumulation of charges at the cutoff boundary or very wrong energies! For such systems you should consider using one of the implemented long-range electrostatic algorithms, such as particle-mesh Ewald [9, 10].

Boundary conditions are unnatural

Since system size is small (even 10,000 particles is small), a cluster of particles will have a lot of unwanted boundary with its environment (vacuum). This we must avoid if we wish to simulate a bulk system. So we use periodic boundary conditions, to avoid real phase boundaries. But liquids are not crystals, so something unnatural remains. This item is mentioned last because it is the least of the evils. For large systems the errors are small, but for small systems with a lot of internal spatial correlation, the periodic boundaries may enhance internal correlation. In that case, beware and test the influence of system size. This is especially important when using lattice sums for long-range electrostatics, since these are known to sometimes introduce extra ordering.

1.3 Energy Minimization and Search Methods

As mentioned in sec. 1.1, in many cases energy minimization is required. GROMACS provides a number of methods for local energy minimization, as detailed in sec. 3.10.

The potential energy function of a (macro)molecular system is a very complex landscape (or *hyper surface*) in a large number of dimensions. It has one deepest point, the *global minimum* and a very large number of *local minima*, where all derivatives of the potential energy function with respect to the coordinates are zero and all second derivatives are nonnegative. The matrix of second derivatives, which is called the *Hessian matrix*, has nonnegative eigenvalues; only the collective coordinates that correspond to translation and rotation (for an isolated molecule) have zero eigenvalues. In between the local minima there are *saddle points*, where the Hessian matrix has only one negative eigenvalue. These points are the mountain passes through which the system can migrate from one local minimum to another.

Knowledge of all local minima, including the global one, and of all saddle points would enable us to describe the relevant structures and conformations and their free energies, as well as the dynamics of structural transitions. Unfortunately, the dimensionality of the configurational space and the number of local minima is so high that it is impossible to sample the space at a sufficient number of points to obtain a complete survey. In particular, no minimization method exists that guarantees the determination of the global minimum in any practical amount of time [Impractical methods exist, some much faster than others [11]]. However, given a starting configuration, it is possible to find the *nearest local minimum*. Nearest in this context does not always imply nearest in a geometrical sense (*i.e.*, the least sum of square coordinate differences), but means the minimum that can be reached by systematically moving down the steepest local gradient. Finding this nearest local minimum is all that GROMACS can do for you, sorry! If you want to find other minima and hope to discover the global minimum in the process, the best advice is to experiment

with temperature-coupled MD: run your system at a high temperature for a while and then quench it slowly down to the required temperature; do this repeatedly! If something as a melting or glass transition temperature exists, it is wise to stay for some time slightly below that temperature and cool down slowly according to some clever scheme, a process called *simulated annealing*. Since no physical truth is required, you can use your imagination to speed up this process. One trick that often works is to make hydrogen atoms heavier (mass 10 or so): although that will slow down the otherwise very rapid motions of hydrogen atoms, it will hardly influence the slower motions in the system while enabling you to increase the time step by a factor of 3 or 4. You can also modify the potential energy function during the search procedure, *e.g.* by removing barriers (remove dihedral angle functions or replace repulsive potentials by *soft core* potentials [12]), but always take care to restore the correct functions slowly. The best search method that allows rather drastic structural changes is to allow excursions into four-dimensional space [13], but this requires some extra programming beyond the standard capabilities of GROMACS.

Three possible energy minimization methods are:

- Those that require only function evaluations. Examples are the simplex method and its variants. A step is made on the basis of the results of previous evaluations. If derivative information is available, such methods are inferior to those that use this information.
- Those that use derivative information. Since the partial derivatives of the potential energy with respect to all coordinates are known in MD programs (these are equal to minus the forces) this class of methods is very suitable as modification of MD programs.
- Those that use second derivative information as well. These methods are superior in their convergence properties near the minimum: a quadratic potential function is minimized in one step! The problem is that for N particles a $3N \times 3N$ matrix must be computed, stored and inverted. Apart from the extra programming to obtain second derivatives, for most systems of interest this is beyond the available capacity. There are intermediate methods building up the Hessian matrix on the fly, but they also suffer from excessive storage requirements. So GROMACS will shy away from this class of methods.

The *steepest descent* method, available in GROMACS, is of the second class. It simply takes a step in the direction of the negative gradient (hence in the direction of the force), without any consideration of the history built up in previous steps. The step size is adjusted such that the search is fast but the motion is always downhill. This is a simple and sturdy, but somewhat stupid, method: its convergence can be quite slow, especially in the vicinity of the local minimum! The faster converging *conjugate gradient method* (see *e.g.* [14]) uses gradient information from previous steps. In general, steepest descents will bring you close to the nearest local minimum very quickly, while conjugate gradients brings you *very* close to the local minimum, but performs worse far away from the minimum. GROMACS also supports the L-BFGS minimizer, which is mostly comparable to *conjugate gradient method*, but in some cases converges faster.

Chapter 2

Definitions and Units

2.1 Notation

The following conventions for mathematical typesetting are used throughout this document:

Item	Notation	Example
Vector	Bold italic	\mathbf{r}_i
Vector Length	Italic	r_i

We define the *lowercase* subscripts i, j, k and l to denote particles: \mathbf{r}_i is the *position vector* of particle i , and using this notation:

$$\mathbf{r}_{ij} = \mathbf{r}_j - \mathbf{r}_i \quad (2.1)$$

$$r_{ij} = |\mathbf{r}_{ij}| \quad (2.2)$$

The force on particle i is denoted by \mathbf{F}_i and

$$\mathbf{F}_{ij} = \text{force on } i \text{ exerted by } j \quad (2.3)$$

Please note that we changed notation as of ver. 2.0 to $\mathbf{r}_{ij} = \mathbf{r}_j - \mathbf{r}_i$ since this is the notation commonly used. If you encounter an error, let us know.

2.2 MD units

GROMACS uses a consistent set of units that produce values in the vicinity of unity for most relevant molecular quantities. Let us call them *MD units*. The basic units in this system are nm, ps, K, electron charge (e) and atomic mass unit (u), see Table 2.1.

Consistent with these units are a set of derived units, given in Table 2.2.

The **electric conversion factor** $f = \frac{1}{4\pi\epsilon_0} = 138.935\,485(9) \text{ kJ mol}^{-1} \text{ nm e}^{-2}$. It relates the mechanical quantities to the electrical quantities as in

$$V = f \frac{q^2}{r} \text{ or } F = f \frac{q^2}{r^2} \quad (2.4)$$

Quantity	Symbol	Unit
length	r	nm = 10^{-9} m
mass	m	u (atomic mass unit) = $1.6605402(10) \times 10^{-27}$ kg (1/12 the mass of a ^{12}C atom) $1.6605402(10) \times 10^{-27}$ kg
time	t	ps = 10^{-12} s
charge	q	e = electronic charge = $1.60217733(49) \times 10^{-19}$ C
temperature	T	K

Table 2.1: Basic units used in GROMACS. Numbers in parentheses give accuracy.

Quantity	Symbol	Unit
energy	E, V	kJ mol^{-1}
Force	\mathbf{F}	$\text{kJ mol}^{-1} \text{ nm}^{-1}$
pressure	p	$\text{kJ mol}^{-1} \text{ nm}^{-3} = 10^{30}/N_{AV}$ Pa 1.66054×10^6 Pa = 16.6054 Bar
velocity	v	$\text{nm ps}^{-1} = 1000$ m/s
dipole moment	μ	e nm
electric potential	Φ	$\text{kJ mol}^{-1} e^{-1} = 0.010364272(3)$ Volt
electric field	E	$\text{kJ mol}^{-1} \text{ nm}^{-1} e^{-1} = 1.0364272(3) \times 10^7$ V/m

Table 2.2: Derived units

Electric potentials Φ and electric fields \mathbf{E} are intermediate quantities in the calculation of energies and forces. They do not occur inside GROMACS. If they are used in evaluations, there is a choice of equations and related units. We recommend strongly to follow the usual practice to include the factor f in expressions that evaluate Φ and \mathbf{E} :

$$\Phi(\mathbf{r}) = f \sum_j \frac{q_j}{|\mathbf{r} - \mathbf{r}_j|} \quad (2.5)$$

$$\mathbf{E}(\mathbf{r}) = f \sum_j q_j \frac{(\mathbf{r} - \mathbf{r}_j)}{|\mathbf{r} - \mathbf{r}_j|^3} \quad (2.6)$$

With these definitions $q\Phi$ is an energy and $q\mathbf{E}$ is a force. The units are those given in Table 2.2: about 10 mV for potential. Thus the potential of an electronic charge at a distance of 1 nm equals $f \approx 140$ units ≈ 1.4 V. (exact value: 1.439965 V)

Note that these units are mutually consistent; changing any of the units is likely to produce inconsistencies and is therefore *strongly discouraged!* In particular: if \AA are used instead of nm, the unit of time changes to 0.1 ps. If the kcal/mol (= 4.184 kJ/mol) is used instead of kJ/mol for energy, the unit of time becomes 0.488882 ps and the unit of temperature changes to 4.184 K. But in both cases all electrical energies go wrong, because they will still be computed in kJ/mol, expecting nm as the unit of length. Although careful rescaling of charges may still yield consistency, it is clear that such confusions must be rigidly avoided.

In terms of the MD units the usual physical constants take on different values, see Table 2.3. All quantities are per mol rather than per molecule. There is no distinction between Boltzmann's constant k and the gas constant R : their value is $0.00831451 \text{ kJ mol}^{-1} \text{ K}^{-1}$.

Symbol	Name	Value
N_{AV}	Avogadro's number	$6.022\,136\,7(36) \times 10^{23} \text{ mol}^{-1}$
R	gas constant	$8.314\,510(70) \times 10^{-3} \text{ kJ mol}^{-1} \text{ K}^{-1}$
k_B	Boltzmann's constant	idem
h	Planck's constant	$0.399\,031\,32(24) \text{ kJ mol}^{-1} \text{ ps}$
\hbar	Dirac's constant	$0.063\,507\,807(38) \text{ kJ mol}^{-1} \text{ ps}$
c	velocity of light	$299\,792.458 \text{ nm/ps}$

Table 2.3: Some Physical Constants

Quantity	Symbol	Relation to SI
Length	r^*	$r \sigma^{-1}$
Mass	m^*	$m M^{-1}$
Time	t^*	$t \sigma^{-1} \sqrt{\epsilon/M}$
Temperature	T^*	$k_B T \epsilon^{-1}$
Energy	E^*	$E \epsilon^{-1}$
Force	F^*	$F \sigma \epsilon^{-1}$
Pressure	P^*	$P \sigma^3 \epsilon^{-1}$
Velocity	v^*	$v \sqrt{M/\epsilon}$
Density	ρ^*	$N \sigma^3 V^{-1}$

Table 2.4: Reduced Lennard-Jones quantities

2.3 Reduced units

When simulating Lennard-Jones (LJ) systems it might be advantageous to use reduced units (i.e., setting $\epsilon_{ii} = \sigma_{ii} = m_i = k_B = 1$ for one type of atoms). This is possible. When specifying the input in reduced units, the output will also be in reduced units. There is one exception: the *temperature*, which is expressed in 0.008 314 51 reduced units. This is a consequence of the use of Boltzmann's constant in the evaluation of temperature in the code. Thus not T , but $k_B T$ is the reduced temperature. A GROMACS temperature $T = 1$ means a reduced temperature of 0.008... units; if a reduced temperature of 1 is required, the GROMACS temperature should be 120.2717.

In Table 2.4 quantities are given for LJ potentials:

$$V_{LJ} = 4\epsilon \left[\left(\frac{\sigma}{r} \right)^{12} - \left(\frac{\sigma}{r} \right)^6 \right] \quad (2.7)$$

Chapter 3

Algorithms

3.1 Introduction

In this chapter we first give describe some general concepts used in GROMACS: *periodic boundary conditions* (sec. 3.2) and the *group concept* (sec. 3.3). The MD algorithm is described in sec. 3.4: first a global form of the algorithm is given, which is refined in subsequent subsections. The (simple) EM (Energy Minimization) algorithm is described in sec. 3.10. Some other algorithms for special purpose dynamics are described after this.

A few issues are of general interest. In all cases the *system* must be defined, consisting of molecules. Molecules again consist of particles with defined interaction functions. The detailed description of the *topology* of the molecules and of the *force field* and the calculation of forces is given in chapter 4. In the present chapter we describe other aspects of the algorithm, such as pair list generation, update of velocities and positions, coupling to external temperature and pressure, conservation of constraints. The *analysis* of the data generated by an MD simulation is treated in chapter 8.

3.2 Periodic boundary conditions

The classical way to minimize edge effects in a finite system is to apply *periodic boundary conditions*. The atoms of the system to be simulated are put into a space-filling box, which is surrounded by translated copies of itself (Fig. 3.1). Thus there are no boundaries of the system; the artifact caused by unwanted boundaries in an isolated cluster is now replaced by the artifact of periodic conditions. If a crystal is simulated, such boundary conditions are desired (although motions are naturally restricted to periodic motions with wavelengths fitting into the box). If one wishes to simulate non-periodic systems, as liquids or solutions, the periodicity by itself causes errors. The errors can be evaluated by comparing various system sizes; they are expected to be less severe than the errors resulting from an unnatural boundary with vacuum.

There are several possible shapes for space-filling unit cells. Some, as the *rhombic dodecahedron* and the *truncated octahedron* [15] are closer to a sphere than a cube is and are therefore more

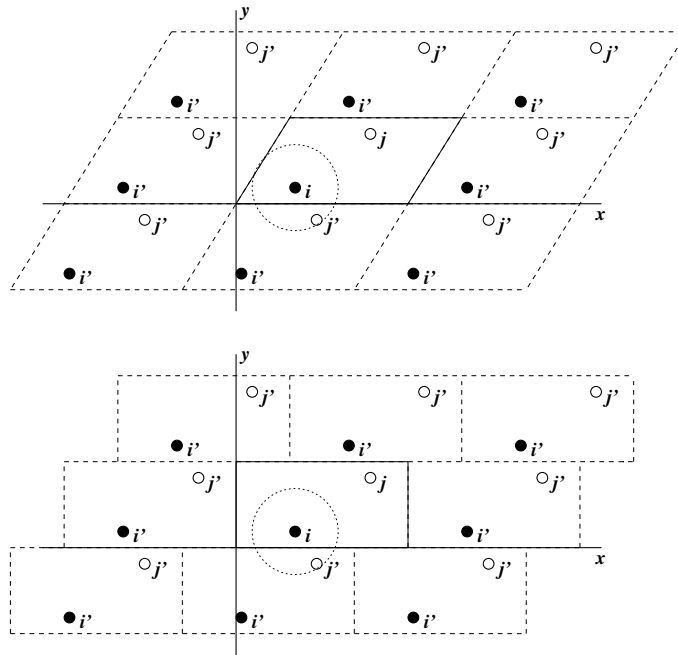


Figure 3.1: Periodic boundary conditions in two dimensions.

economical for studying an (approximately spherical) macromolecule in solution, since fewer solvent molecules are required to fill the box given a minimum distance between macromolecular images. However, a periodic system based on the rhombic dodecahedron or truncated octahedron is equivalent to a periodic system based on a *triclinic* unit cell. The latter shape is the most general space-filling unit cell; it comprises all possible space-filling shapes [16]. Therefore GROMACS is based on the triclinic unit cell.

GROMACS uses periodic boundary conditions, combined with the *minimum image convention*: only one - the nearest - image of each particle is considered for short-range non-bonded interaction terms. For long-range electrostatic interactions this is not always accurate enough, and GROMACS therefore also incorporates lattice sum methods like Ewald Sum, PME and PPPM.

Gromacs supports triclinic boxes of any shape. The box is defined by the 3 box vectors \mathbf{a} , \mathbf{b} and \mathbf{c} . The box vectors must satisfy the following conditions:

$$a_y = a_z = b_z = 0 \quad (3.1)$$

$$a_x > 0, \quad b_y > 0, \quad c_z > 0 \quad (3.2)$$

$$|b_x| \leq \frac{1}{2} a_x, \quad |c_x| \leq \frac{1}{2} a_x, \quad |c_y| \leq \frac{1}{2} b_y \quad (3.3)$$

Equations 3.1 can always be satisfied by rotating the box. Inequalities (3.2) and (3.3) can always be satisfied by adding and subtracting box vectors.

Even when simulating using a triclinic box, GROMACS always puts the particles in a brick shaped volume, for efficiency reasons. This is illustrated in Fig. 3.1 for a 2-dimensional system. So from the output trajectory it might seem like the simulation was done in a rectangular box. The program `trjconv` can be used to convert the trajectory to a different unit-cell representation.

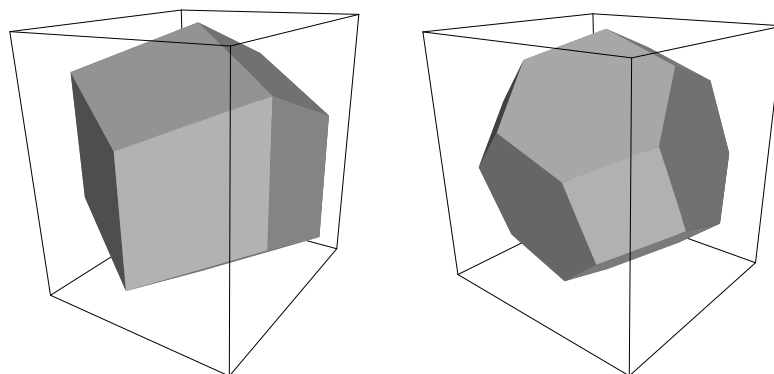


Figure 3.2: A rhombic dodecahedron and truncated octahedron (arbitrary orientations).

box type	image distance	box volume	box vectors			box vector angles		
			a	b	c	$\angle bc$	$\angle ac$	$\angle ab$
cubic	d	d^3	d 0 0	0 d 0	0 0 d	90°	90°	90°
rhombic dodecahedron (xy-square)	d	$\frac{1}{2}\sqrt{2}d^3$ $0.707d^3$	d 0 0	0 d 0	$\frac{1}{2}d$ $\frac{1}{2}d$ $\frac{1}{2}\sqrt{2}d$	60°	60°	90°
rhombic dodecahedron (xy-hexagon)	d	$\frac{1}{2}\sqrt{2}d^3$ $0.707d^3$	d 0 0	$\frac{1}{2}d$ $\frac{1}{2}\sqrt{3}d$ 0	$\frac{1}{2}d$ $\frac{1}{6}\sqrt{3}d$ $\frac{1}{3}\sqrt{6}d$	60°	60°	60°
truncated octahedron	d	$\frac{4}{9}\sqrt{3}d^3$ $0.770d^3$	d 0 0	$\frac{1}{3}d$ $\frac{2}{3}\sqrt{2}d$ 0	$-\frac{1}{3}d$ $\frac{1}{3}\sqrt{2}d$ $\frac{1}{3}\sqrt{6}d$	71.53°	109.47°	71.53°

Table 3.1: The cubic box, the rhombic dodecahedron and the truncated octahedron.

It is also possible to simulate without periodic boundary conditions, but it is more efficient to simulate an isolated cluster of molecules in a large periodic box, since fast grid searching can only be used in a periodic system.

3.2.1 Some useful box types

The three most useful box types for simulations of solvated systems are described in Table 3.1. The rhombic dodecahedron (Fig. 3.2) is the smallest and most regular space-filling unit cell. Each of the 12 image cells is at the same distance. The volume is 71% of the volume of a cube having the same image distance. This saves about 29% of CPU-time when simulating a spherical or flexible molecule in solvent. There are two different orientations of a rhombic dodecahedron that satisfy equations 3.1, 3.2 and 3.3. The program `editconf` produces the orientation which has a square intersection with the xy-plane. This orientation was chosen because the first two box vectors coincide with the x and y-axis, which is easier to comprehend. The other orientation can

be useful for simulations of membrane proteins. In this case the cross-section with the xy-plane is a hexagon, which has an area which is 14% smaller than the area of a square with the same image distance. The height of the box (c_z) should be changed to obtain an optimal spacing. This box shape not only saves CPU-time, it also results in a more uniform arrangement of the proteins.

3.2.2 Cutoff restrictions

The minimum image convention implies that the cutoff radius used to truncate non-bonded interactions must not exceed half the shortest box vector:

$$R_c < \frac{1}{2} \min(\|\mathbf{a}\|, \|\mathbf{b}\|, \|\mathbf{c}\|), \quad (3.4)$$

otherwise more than one image would be within the cutoff distance of the force. When a macromolecule, such as a protein, is studied in solution, this restriction does not suffice. In principle a single solvent molecule should not be able to ‘see’ both sides of the macromolecule. This means that the length of each box vector must exceed the length of the macromolecule in the direction of that edge *plus* two times the cutoff radius R_c . It is common to compromise in this respect, and make the solvent layer somewhat smaller in order to reduce the computational cost. For efficiency reasons the cutoff with triclinic boxes is more restricted. For grid search the extra restriction is weak:

$$R_c < \min(a_x, b_y, c_z) \quad (3.5)$$

For simple search the extra restriction is stronger:

$$R_c < \frac{1}{2} \min(a_x, b_y, c_z) \quad (3.6)$$

Each unit cell (cubic, rectangular or triclinic) is surrounded by 26 translated images. Thus a particular image can always be identified by an index pointing to one of 27 *translation vectors* and constructed by applying a translation with the indexed vector (see 3.4.3). Restriction (3.5) ensures that only 26 images need to be considered.

3.3 The group concept

In the GROMACS MD and analysis programs one uses *groups* of atoms to perform certain actions on. The maximum number of groups is 256, but each atom can only belong to six different groups, one each of the following:

T-coupling group The temperature coupling parameters (reference temperature, time constant, number of degrees of freedom, see 3.4.4) can be defined for each T-coupling group separately. For example, in a solvated macromolecule the solvent (that tends to generate more heating by force and integration errors) can be coupled with a shorter time constant to a bath than is a macromolecule, or a surface can be kept cooler than an adsorbing molecule. Many different T-coupling groups may be defined. See also center of mass groups below.

Freeze group Atoms that belong to a freeze group are kept stationary in the dynamics. This is useful during equilibration, *e.g.* to avoid badly placed solvent molecules giving unreasonable kicks to protein atoms, although the same effect can also be obtained by putting a restraining potential on the atoms that must be protected. The freeze option can be used, if desired, on just one or two coordinates of an atom, thereby freezing the atoms in a plane or on a line. When an atom is partially frozen, constraints will still be able to move it, even in a frozen direction. A fully frozen atom can not be moved by constraints. Many freeze groups can be defined. Frozen coordinates are unaffected by pressure scaling, in some cases this can produce unwanted results, in particular when constraints are used as well (in this case you will get very large pressures). Because of this it is recommended to not combine freeze groups with constraints and pressure coupling. For the sake of equilibration it could suffice to start with freezing in a constant volume simulation, and afterwards use position restraints in conjunction with constant pressure.

Accelerate group On each atom in an 'accelerate group' an acceleration a^g is imposed. This is equivalent to an external force. This feature makes it possible to drive the system into a non-equilibrium state and enables the performance of non-equilibrium MD and hence to obtain transport properties.

Energy monitor group Mutual interactions between all energy monitor groups are compiled during the simulation. This is done separately for Lennard-Jones and Coulomb terms. In principle up to 256 groups could be defined, but that would lead to 256×256 items! Better use this concept sparingly.

All non-bonded interactions between pairs of energy monitor groups can be excluded (see sec. 7.3.1). Pairs of particles from excluded pairs of energy monitor groups are not put into the pair list. This can result in a significant speedup for simulations where interactions within or between parts of the system are not required.

Center of mass group In GROMACS the center of mass (COM) motion can be removed, for either the complete system or for groups of atoms. The latter is useful, *e.g.* for systems where there is limited friction (*e.g.* gas systems) to prevent center of mass motion to occur. It makes sense to use the same groups for Temperature coupling and center of mass motion removal.

XTC output group In order to reduce the size of the XTC trajectory file, only a subset of all particles can be stored. All XTC groups that are specified are saved, the rest is not. If no XTC groups are specified, than all atoms are saved to the XTC file.

The use of groups in analysis programs is described in chapter 8.

3.4 Molecular Dynamics

A global flow scheme for MD is given in Fig. 3.3. Each MD or EM run requires as input a set of initial coordinates and - optionally - initial velocities of all particles involved. This chapter does not describe how these are obtained; for the setup of an actual MD run check the online manual at www.gromacs.org.

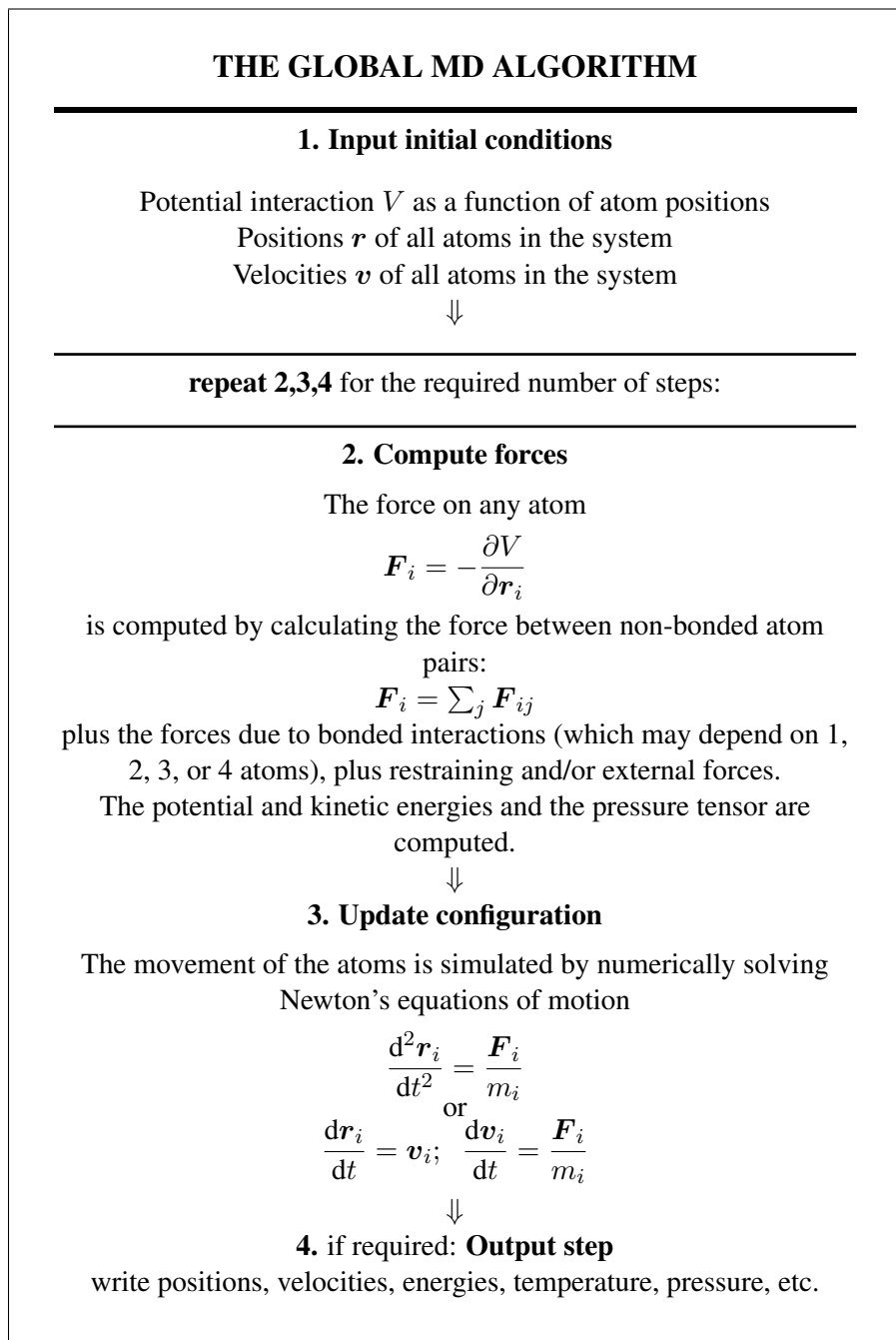


Figure 3.3: The global MD algorithm

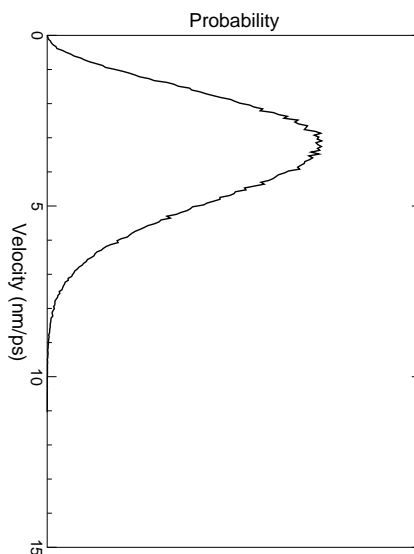


Figure 3.4: A Maxwellian distribution, generated from random numbers.

3.4.1 Initial conditions

Topology and force field

The system topology, including a description of the force field, must be loaded. These items are described in chapter 4. All this information is static; it is never modified during the run.

Coordinates and velocities

Then, before a run starts, the box size and the coordinates and velocities of all particles are required. The box size is determined by three vectors (nine numbers) $\mathbf{b}_1, \mathbf{b}_2, \mathbf{b}_3$, which represent the three basis vectors of the periodic box. While in the present version of GROMACS only rectangular boxes are allowed, three numbers suffice, but the use of three vectors already prepares for arbitrary triclinic boxes to be implemented in a later version.

If the run starts at $t = t_0$, the coordinates at $t = t_0$ must be known. The *leap-frog algorithm*, used to update the time step with Δt (see 3.4.4), requires that the velocities at $t = t_0 - \frac{\Delta t}{2}$ are known. If velocities are not available, the program can generate initial atomic velocities $v_i, i = 1 \dots 3N$ with a Maxwellian distribution (Fig. 3.4) at a given absolute temperature T :

$$p(v_i) = \sqrt{\frac{m_i}{2\pi kT}} \exp\left(-\frac{m_i v_i^2}{2kT}\right) \quad (3.7)$$

where k is Boltzmann's constant (see chapter 2). To accomplish this, normally distributed random numbers are generated by adding twelve random numbers R_k in the range $0 \leq R_k < 1$ and subtracting 6.0 from their sum. The result is then multiplied by the standard deviation of the velocity distribution $\sqrt{kT/m_i}$. Since the resulting total energy will not correspond exactly to the

required temperature T , a correction is made: first the center-of-mass motion is removed and then all velocities are scaled so that the total energy corresponds exactly to T (see eqn. 3.13).

Center-of-mass motion

The center-of-mass velocity is normally set to zero at every step. Normally there is no net external force acting on the system and the center-of-mass velocity should remain constant. In practice, however, the update algorithm develops a very slow change in the center-of-mass velocity, and thus in the total kinetic energy of the system, specially when temperature coupling is used. If such changes are not quenched, an appreciable center-of-mass motion develops eventually in long runs, and the temperature will be significantly misinterpreted. The same may happen due to overall rotational motion, but only when an isolated cluster is simulated. In periodic systems with filled boxes, the overall rotational motion is coupled to other degrees of freedom and does not give any problems.

3.4.2 Neighbor searching

As mentioned in chapter 4, internal forces are either generated from fixed (static) lists, or from dynamics lists. The latter concern non-bonded interactions between any pair of particles. When calculating the non-bonded forces, it is convenient to have all particles in a rectangular box. As shown in Fig. 3.1, it is possible to transform a triclinic box into a rectangular box. The output coordinates are always in a rectangular box, even when a dodecahedron or triclinic box was used for the simulation. Equation 3.1 ensures that we can reset particles in a rectangular box by first shifting them with box vector \mathbf{c} , then with \mathbf{b} and finally with \mathbf{a} . Equations 3.3, 3.4 and 3.5 ensure that we can find the 14 nearest triclinic images within a linear combination which does not involve multiples of box vectors.

Pair lists generation

The non-bonded pair forces need to be calculated only for those pairs i, j for which the distance r_{ij} between i and the nearest image of j is less than a given cutoff radius R_c . Some of the particle pairs that fulfill this criterion are excluded, when their interaction is already fully accounted for by bonded interactions. GROMACS employs a *pair list* that contains those particle pairs for which non-bonded forces must be calculated. The pair list contains atoms i , a displacement vector for atom i , and all particles j that are within `rshort` of this particular image of atom i . The list is updated every `nstlist` steps, where `nstlist` is typically 10 for the GROMACS forcefield and 5 for the GROMOS-96 forcefield. There is an option to calculate the total non-bonded force on each particle due to all particle in a shell around the list-cutoff, *i.e.* at a distance between `rshort` and `rlong`. This force is calculated during the pair list update and retained during `nstlist` steps.

To make the neighbor list all particles that are close (*i.e.* within the cutoff) to a given particle must be found. This searching, usually called neighbor searching (NS), involves periodic boundary conditions and determining the *image* (see sec. 3.2). Without periodic boundary conditions a simple $O(N^2)$ algorithm must be used. With periodic boundary conditions a grid search can be

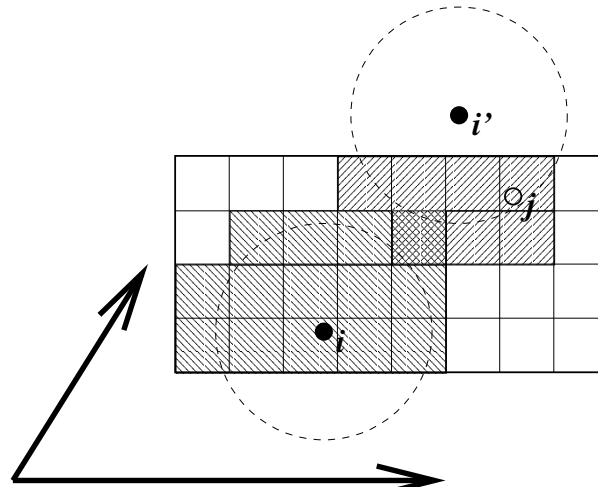


Figure 3.5: Grid search in two dimensions. The arrows are the box vectors.

used, which is $O(N)$.

Simple search

Due to eqns. 3.1 and 3.6, the vector \mathbf{r}_{ij} connecting images within the cutoff R_c can be found by constructing:

$$\mathbf{r}''' = \mathbf{r}_j - \mathbf{r}_i \quad (3.8)$$

$$\mathbf{r}'' = \mathbf{r}''' - \mathbf{c} * \text{round}(r_z'''/c_z) \quad (3.9)$$

$$\mathbf{r}' = \mathbf{r}'' - \mathbf{b} * \text{round}(r_y''/b_y) \quad (3.10)$$

$$\mathbf{r}_{ij} = \mathbf{r}' - \mathbf{a} * \text{round}(r_x'/a_x) \quad (3.11)$$

When distances between two particles in a triclinic box are needed that do not obey eqn. 3.1, many shifts of combinations of box vectors need to be considered to find the nearest image.

Grid search

The grid search is schematically depicted in Fig. 3.5. All particles are put on the NS grid, with the smallest spacing $\geq R_c/2$ in each of the directions. In the direction of each box vector, a particle i has three images. For each direction the image may be -1, 0 or 1, corresponding to a translation over -1, 0 or +1 box vector. We do not search the surrounding NS grid cells for neighbors of i and then calculate the image, but rather construct the images first and then search neighbors corresponding to that image of i . As Fig. 3.5 shows, some grid cells may be searched more than once for different images of i . This is not a problem, since, due to the minimum image convention, at most one image will “see” the j -particle. For every particle, fewer than 125 (5^3) neighboring cells are searched. Therefore, the algorithm scales linearly with the number of particles. Although the prefactor is large, the scaling behavior makes the algorithm far superior over the standard $O(N^2)$ algorithm when there are more than a few hundred particles. The grid search is equally

fast for rectangular and triclinic boxes. Thus for most protein and peptide simulations the rhombic dodecahedron will be the preferred box shape.

Charge groups

Where applicable, neighbor searching is carried out on the basis of *charge groups*. A charge group is a small set of nearby atoms with an integer net charge. Charge groups are defined in the molecular topology. If the nearest image distance between the *geometrical centers* of the atoms of two charge groups is less than the cutoff radius, all atom pairs between the charge groups are included in the pair list. This procedure avoids the creation of charges due to the use of a cutoff (when one charge of a dipole is within range and the other not), which can have disastrous consequences for the behavior of the Coulomb interaction function at distances near the cutoff radius. If molecular groups have full charges (ions), charge groups do not avoid adverse cutoff effects, and you should consider using one of the lattice sum methods supplied by GROMACS [17].

If appropriately constructed shift functions are used for the electrostatic forces, no charge groups are needed. Such shift functions are implemented in GROMACS (see chapter 4) but must be used with care: in principle, they should be combined with a lattice sum for long-range electrostatics.

Charge groups also provide a speed up of the neighbor search. The neighbor searching for a water system, for instance, is $3^2 = 9$ times faster when each molecule is treated as a charge group.

3.4.3 Compute forces

Potential energy

When forces are computed, the potential energy of each interaction term is computed as well. The total potential energy is summed for various contributions, such as Lennard-Jones, Coulomb, and bonded terms. It is also possible to compute these contributions for *groups* of atoms that are separately defined (see sec. 3.3).

Kinetic energy and temperature

The temperature is given by the total kinetic energy of the N -particle system:

$$E_{kin} = \frac{1}{2} \sum_{i=1}^N m_i v_i^2 \quad (3.12)$$

From this the absolute temperature T can be computed using:

$$\frac{1}{2} N_{df} kT = E_{kin} \quad (3.13)$$

where k is Boltzmann's constant and N_{df} is the number of degrees of freedom which can be computed from:

$$N_{df} = 3N - N_c - N_{com} \quad (3.14)$$

Here N_c is the number of *constraints* imposed on the system. When performing molecular dynamics $N_{com} = 3$ additional degrees of freedom must be removed, because the three center-of-mass

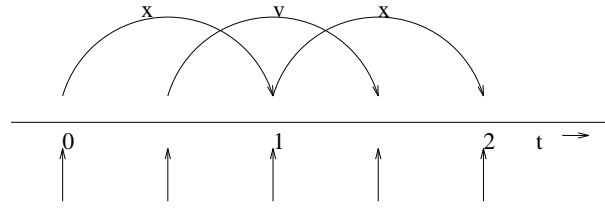


Figure 3.6: The Leap-Frog integration method. The algorithm is called Leap-Frog because r and v are leaping like frogs over each others back.

velocities are constants of the motion, which are usually set to zero. When simulating in vacuo, the rotation around the center of mass can also be removed, in this case $N_{com} = 6$. When more than one temperature coupling group is used, the number of degrees of freedom for group i is:

$$N_{df}^i = (3N^i - N_c^i) \frac{3N - N_c - N_{com}}{3N - N_c} \quad (3.15)$$

The kinetic energy can also be written as a tensor, which is necessary for pressure calculation in a triclinic system, or systems where shear forces are imposed:

$$\mathbf{E}_{kin} = \frac{1}{2} \sum_i^N m_i \mathbf{v}_i \otimes \mathbf{v}_i \quad (3.16)$$

Pressure and virial

The pressure tensor \mathbf{P} is calculated from the difference between kinetic energy E_{kin} and the virial Ξ

$$\mathbf{P} = \frac{2}{V} (\mathbf{E}_{kin} - \Xi) \quad (3.17)$$

where V is the volume of the computational box. The scalar pressure P , which can be used for pressure coupling in the case of isotropic systems, is computed as:

$$P = \text{trace}(\mathbf{P})/3 \quad (3.18)$$

The virial Ξ tensor is defined as

$$\Xi = -\frac{1}{2} \sum_{i < j} \mathbf{r}_{ij} \otimes \mathbf{F}_{ij} \quad (3.19)$$

The GROMACS implementation of the virial computation is described in sec. B.1.

3.4.4 Update configuration

The GROMACS MD program utilizes the so-called *leap-frog* algorithm [18] for the integration of the equations of motion. The leap-frog algorithm uses positions \mathbf{r} at time t and velocities \mathbf{v} at

time $t - \frac{\Delta t}{2}$; it updates positions and velocities using the forces $\mathbf{F}(t)$ determined by the positions at time t :

$$\mathbf{v}(t + \frac{\Delta t}{2}) = \mathbf{v}(t - \frac{\Delta t}{2}) + \frac{\mathbf{F}(t)}{m} \Delta t \quad (3.20)$$

$$\mathbf{r}(t + \Delta t) = \mathbf{r}(t) + \mathbf{v}(t + \frac{\Delta t}{2}) \Delta t \quad (3.21)$$

The algorithm is visualized in Fig. 3.6. It is equivalent to the Verlet [19] algorithm:

$$\mathbf{r}(t + \Delta t) = 2\mathbf{r}(t) - \mathbf{r}(t - \Delta t) + \frac{\mathbf{F}(t)}{m} \Delta t^2 + O(\Delta t^4) \quad (3.22)$$

The algorithm is of third order in \mathbf{r} and is time-reversible. See ref. [20] for the merits of this algorithm and comparison with other time integration algorithms.

The equations of motion are modified for temperature coupling and pressure coupling, and extended to include the conservation of constraints, all of which are described below.

3.4.5 Temperature coupling

For several reasons (drift during equilibration, drift as a result of force truncation and integration errors, heating due to external or frictional forces), it is necessary to control the temperature of the system. GROMACS can use either the *weak coupling* scheme of Berendsen [21] or the extended ensemble Nosé-Hoover scheme [22, 23].

Berendsen temperature coupling

The Berendsen algorithm mimics weak coupling with first-order kinetics to an external heat bath with given temperature T_0 . See ref. [24] for a comparison with the Nosé-Hoover scheme. The effect of this algorithm is that a deviation of the system temperature from T_0 is slowly corrected according to

$$\frac{dT}{dt} = \frac{T_0 - T}{\tau} \quad (3.23)$$

which means that a temperature deviation decays exponentially with a time constant τ . This method of coupling has the advantage that the strength of the coupling can be varied and adapted to the user requirement: for equilibration purposes the coupling time can be taken quite short (*e.g.* 0.01 ps), but for reliable equilibrium runs it can be taken much longer (*e.g.* 0.5 ps) in which case it hardly influences the conservative dynamics.

The heat flow into or out of the system is effected by scaling the velocities of each particle every step with a time-dependent factor λ , given by

$$\lambda = \left[1 + \frac{\Delta t}{\tau_T} \left\{ \frac{T_0}{T(t - \frac{\Delta t}{2})} - 1 \right\} \right]^{1/2} \quad (3.24)$$

The parameter τ_T is close to, but not exactly equal to the time constant τ of the temperature coupling (eqn. 3.23):

$$\tau = 2C_V \tau_T / N_d f k \quad (3.25)$$

where C_V is the total heat capacity of the system, k is Boltzmann's constant, and N_{df} is the total number of degrees of freedom. The reason that $\tau \neq \tau_T$ is that the kinetic energy change caused by scaling the velocities is partly redistributed between kinetic and potential energy and hence the change in temperature is less than the scaling energy. In practice, the ratio τ/τ_T ranges from 1 (gas) to 2 (harmonic solid) to 3 (water). When we use the term 'temperature coupling time constant', we mean the parameter τ_T . **Note** that in practice the scaling factor λ is limited to the range of $0.8 \leq \lambda \leq 1.25$, to avoid scaling by very large numbers which may crash the simulation. In normal use, λ will always be much closer to 1.0.

Strictly, for computing the scaling factor the temperature T is needed at time t , but this is not available in the algorithm. In practice, the temperature at the previous time step is used (as indicated in eqn. 3.24), which is perfectly all right since the coupling time constant is much longer than one time step. The Berendsen algorithm is stable up to $\tau_T \approx \Delta t$. [A simple steepest-descent minimizer can be implemented by setting $T = 0$ and $\tau_T \ll \delta t$.]

Nosé-Hoover temperature coupling

The Berendsen weak coupling algorithm is extremely efficient for relaxing a system to the target temperature, but once your system has reached equilibrium it might be more important to probe a correct canonical ensemble. This is unfortunately not the case for the weak coupling scheme, although the difference is usually negligible.

To enable canonical ensemble simulations, GROMACS also supports the extended-ensemble approach first proposed by Nosé[22] and later modified by Hoover[23]. The system Hamiltonian is extended by introducing a thermal reservoir and a friction term in the equations of motion. The friction force is proportional to the product of each particle's velocity and a friction parameter ξ . This friction parameter (or 'heat bath' variable) is a fully dynamic quantity with its own equation of motion; the time derivative is calculated from the difference between the current kinetic energy and the reference temperature.

In Hoover's formulation, the particles' equations of motion in Fig. 3.3 are replaced by

$$\frac{d^2 \mathbf{r}_i}{dt^2} = \frac{\mathbf{F}_i}{m_i} - \xi \frac{d\mathbf{r}_i}{dt}, \quad (3.26)$$

where the equation of motion for the heat bath parameter ξ is

$$\frac{d\xi}{dt} = \frac{1}{Q} (T - T_0). \quad (3.27)$$

The reference temperature is denoted T_0 , while T is the current instantaneous temperature of the system. The strength of the coupling is determined by the constant Q (usually called the 'mass parameter' of the reservoir) in combination with the reference temperature.

In our opinion, the mass parameter is a somewhat awkward way of describing coupling strength, especially due to its dependence on reference temperature (and some implementations even include the number of degrees of freedom in your system when defining Q). To maintain the coupling strength, one would have to change Q in proportion to the change in reference temperature.

For this reason, we prefer to let the GROMACS user work instead with the period τ_T of the oscillations of kinetic energy between the system and the reservoir instead. It is directly related to Q and T_0 via

$$Q = \frac{\tau_T^2 T_0}{4\pi^2}. \quad (3.28)$$

This provides a much more intuitive way of selecting the Nosé-Hoover coupling strength (similar to the weak coupling relaxation), and in addition τ_T is independent of system size and reference temperature.

It is however important to keep the difference between the weak coupling scheme and the Nosé-Hoover algorithm in mind: Using weak coupling you get a strongly damped *exponential relaxation*, while the Nosé-Hoover approach produces an *oscillatory relaxation*. The actual time it takes to relax with Nosé-Hoover coupling is several times larger than the period of the oscillations that you select. These oscillations (in contrast to exponential relaxation) also means that the time constant normally should be 4–5 times larger than the relaxation time used with weak coupling, but your mileage may vary.

Group temperature coupling

In GROMACS temperature coupling can be performed on groups of atoms, typically a protein and solvent. The reason such algorithms were introduced is that energy exchange between different components is not perfect, due to different effects including cutoffs etc. If now the whole system is coupled to one heat bath, water (which experiences the largest cutoff noise) will tend to heat up and the protein will cool down. Typically 100 K differences can be obtained. With the use of proper electrostatic methods (PME) these difference are much smaller but still not negligible. The parameters for temperature coupling in groups are given in the `mdp` file. One special case should be mentioned: it is possible to T-couple only part of the system (or nothing at all obviously). This is done by specifying zero for the time constant τ_T for the group of interest.

3.4.6 Pressure coupling

In the same spirit as the temperature coupling, the system can also be coupled to a 'pressure bath'. GROMACS supports both the Berendsen algorithm [21] that scales coordinates and box vectors every step, and the extended ensemble Parrinello-Rahman approach. Both of these can be combined with any of the temperature coupling methods above.

Berendsen pressure coupling

The Berendsen algorithm rescales the coordinates and box vectors every step with a matrix μ , which has the effect of a first-order kinetic relaxation of the pressure towards a given reference pressure \mathbf{P}_0 :

$$\frac{d\mathbf{P}}{dt} = \frac{\mathbf{P}_0 - \mathbf{P}}{\tau_p} \quad (3.29)$$

The scaling matrix $\boldsymbol{\mu}$ is given by

$$\mu_{ij} = \delta_{ij} - \frac{\Delta t}{3\tau_p} \beta_{ij} \{P_{0ij} - P_{ij}(t)\} \quad (3.30)$$

Here β is the isothermal compressibility of the system. In most cases this will be a diagonal matrix, with equal elements on the diagonal, the value of which is generally not known. It suffices to take a rough estimate because the value of β only influences the non-critical time constant of the pressure relaxation without affecting the average pressure itself. For water at 1 atm and 300 K $\beta = 4.6 \times 10^{-10} \text{ Pa}^{-1} = 4.6 \times 10^{-5} \text{ Bar}^{-1}$, which is 7.6×10^{-4} MD units (see chapter 2). Most other liquids have similar values. When scaling completely anisotropically, the system has to be rotated in order to obey eqn. 3.1. This rotation is approximated in first order in the scaling, which is usually less than 10^{-4} . The actual scaling matrix $\boldsymbol{\mu}'$ is:

$$\boldsymbol{\mu}' = \begin{pmatrix} \mu_{xx} & \mu_{xy} + \mu_{yx} & \mu_{xz} + \mu_{zx} \\ 0 & \mu_{yy} & \mu_{yz} + \mu_{zy} \\ 0 & 0 & \mu_{zz} \end{pmatrix} \quad (3.31)$$

The velocities are neither scaled nor rotated.

In GROMACS, the Berendsen scaling can also be done isotropically, which means that instead of \boldsymbol{P} a diagonal matrix with elements of size $\text{trace}(\boldsymbol{P})/3$ is used. For systems with interfaces, semi-isotropic scaling can be useful. In this case the x/y -directions are scaled isotropically and the z direction is scaled independently. The compressibility in the x/y or z -direction can be set to zero, to scale only in the other direction(s).

If you allow full anisotropic deformations and use constraints you might have to scale more slowly or decrease your timestep to avoid errors from the constraint algorithms.

Parrinello-Rahman pressure coupling

In cases where the fluctuations in pressure or volume are important *per se* (e.g. to calculate thermodynamic properties) it might at least theoretically be a problem that the exact ensemble is not well-defined for the weak coupling scheme.

For this reason, GROMACS also supports constant-pressure simulations using the Parrinello-Rahman approach[25, 26], which is similar to the Nosé-Hoover temperature coupling. With the Parrinello-Rahman barostat, the box vectors as represented by the matrix \boldsymbol{b} obey the matrix equation of motion¹

$$\frac{d\boldsymbol{b}^2}{dt^2} = V\boldsymbol{W}^{-1}\boldsymbol{b}'^{-1}(\boldsymbol{P} - \boldsymbol{P}_{ref}). \quad (3.32)$$

The volume of the box is denoted V , and \boldsymbol{W} is a matrix parameter that determines the strength of the coupling. The matrices \boldsymbol{P} and \boldsymbol{P}_{ref} are the current and reference pressures, respectively.

The equations of motion for the particles are also changed, just as for the Nosé-Hoover coupling. In most cases you would combine the Parrinello-Rahman barostat with the Nosé-Hoover thermostat, but to keep it simple we only show the Parrinello-Rahman modification here:

¹The box matrix representation \boldsymbol{b} in GROMACS corresponds to the transpose of the box matrix representation \boldsymbol{h} in the paper by Nosé and Klein. Because of this, some of our equations will look slightly different.

$$\frac{d^2 \mathbf{r}_i}{dt^2} = \frac{\mathbf{F}_i}{m_i} - \mathbf{M} \frac{d\mathbf{r}_i}{dt}, \quad (3.33)$$

$$\mathbf{M} = \mathbf{b}^{-1} \left[\mathbf{b} \frac{d\mathbf{b}'}{dt} + \frac{d\mathbf{b}}{dt} \mathbf{b}' \right] \mathbf{b}'^{-1}. \quad (3.34)$$

The (inverse) mass parameter matrix \mathbf{W}^{-1} determines the strength of the coupling, and how the box can be deformed. The box restriction (3.1) will be fulfilled automatically if the corresponding elements of \mathbf{W}^{-1} are zero. Since the coupling strength also depends on the size of your box, we prefer to calculate it automatically in GROMACS. You only have to provide the approximate isothermal compressibilities β and the pressure time constant τ_p in the input file (L is the largest box matrix element):

$$\left(\mathbf{W}^{-1}\right)_{ij} = \frac{4\pi^2 \beta_{ij}}{3\tau_p^2 L}. \quad (3.35)$$

Just as for the Nosé-Hoover thermostat, you should realize that the Parrinello-Rahman time constant is *not* equivalent to the relaxation time used in the Berendsen pressure coupling algorithm. In most cases you will need to use a 4–5 times larger time constant with Parrinello-Rahman coupling. If your pressure is very far from equilibrium, the Parrinello-Rahman coupling may result in very large box oscillations that could even crash your run. In that case you would have to increase the time constant, or (better) use the weak coupling scheme to reach the target pressure, and then switch to Parrinello-Rahman coupling once the system is in equilibrium.

Surface tension coupling

When a periodic system consists of more than one phase, separated by surfaces which are parallel to the xy-plane, the surface tension and the z-component of the pressure can be coupled to a pressure bath. Presently, this only works with the Berendsen pressure coupling algorithm in GROMACS. The average surface tension $\gamma(t)$ can be calculated from the difference between the normal and the lateral pressure:

$$\gamma(t) = \frac{1}{n} \int_0^{L_z} \left\{ P_{zz}(z, t) - \frac{P_{xx}(z, t) + P_{yy}(z, t)}{2} \right\} dz \quad (3.36)$$

$$= \frac{L_z}{n} \left\{ P_{zz}(t) - \frac{P_{xx}(t) + P_{yy}(t)}{2} \right\} \quad (3.37)$$

where L_z is the height of the box and n is the number of surfaces. The pressure in the z-direction is corrected by scaling the height of the box with μ_z :

$$\Delta P_{zz} = \frac{\Delta t}{\tau_p} \{ P_{0zz} - P_{zz}(t) \} \quad (3.38)$$

$$\mu_{zz} = 1 + \beta_{zz} \Delta P_{zz} \quad (3.39)$$

This is similar to normal pressure coupling, except that the power of one third is missing. The pressure correction in the z-direction is then used to get the correct convergence for the surface

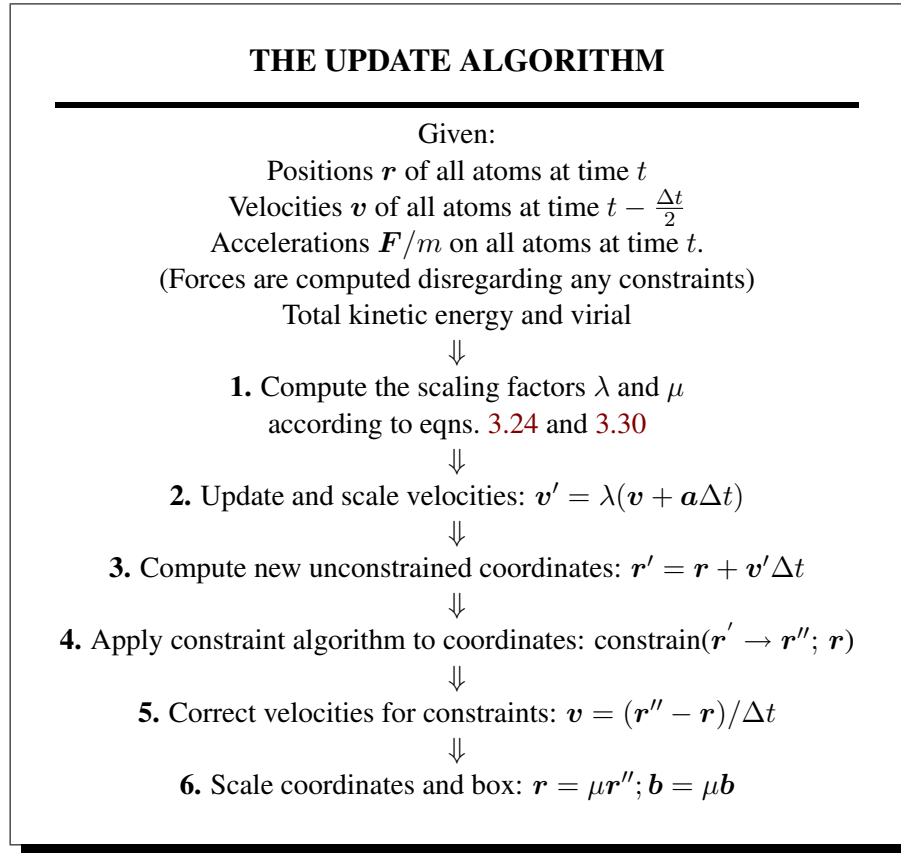


Figure 3.7: The MD update algorithm

tension to the reference value γ_0 . The correction factor for the box-length in the x/y-direction is:

$$\mu_{x/y} = 1 + \frac{\Delta t}{2\tau_p} \beta_{x/y} \left(\frac{n\gamma_0}{\mu_{zz}L_z} - \left\{ P_{zz}(t) + \Delta P_{zz} - \frac{P_{xx}(t) + P_{yy}(t)}{2} \right\} \right) \quad (3.40)$$

The value of β_{zz} is more critical than with normal pressure coupling. Normally an incorrect compressibility will just scale τ_p , but with surface tension coupling it affects the convergence of the surface tension. When β_{zz} is set to zero (constant box height), ΔP_z is also set to zero, which is necessary for obtaining the correct surface tension.

The complete update algorithm

The complete algorithm for the update of velocities and coordinates is given in Fig. 3.7. The SHAKE algorithm of step 4 is explained below.

GROMACS has a provision to "freeze" (prevent motion of) selected particles, which must be defined as a 'freeze group'. This is implemented using a *freeze factor* \mathbf{f}_g , which is a vector, and differs for each *freeze group* (see sec. 3.3). This vector contains only zero (freeze) or one (don't freeze). When we take this freeze factor and the external acceleration \mathbf{a}_h into account the update

algorithm for the velocities becomes:

$$\mathbf{v}(t + \frac{\Delta t}{2}) = \mathbf{f}_g * \lambda * \left[\mathbf{v}(t - \frac{\Delta t}{2}) + \frac{\mathbf{F}(t)}{m} \Delta t + \mathbf{a}_h \Delta t \right] \quad (3.41)$$

where g and h are group indices which differ per atom.

3.4.7 Output step

The important output of the MD run is the *trajectory file* `name.trj` which contains particle coordinates and -optionally- velocities at regular intervals. Since the trajectory files are lengthy, one should not save every step! To retain all information it suffices to write a frame every 15 steps, since at least 30 steps are made per period of the highest frequency in the system, and Shannon's sampling theorem states that two samples per period of the highest frequency in a band-limited signal contain all available information. But that still gives very long files! So, if the highest frequencies are not of interest, 10 or 20 samples per ps may suffice. Be aware of the distortion of high-frequency motions by the *stroboscopic effect*, called *aliasing*: higher frequencies are mirrored with respect to the sampling frequency and appear as lower frequencies.

3.5 Shell molecular dynamics

GROMACS can simulate polarizability using the shell model of Dick and Overhauser [27]. In such models a shell particle representing the electronic degrees of freedom is attached to a nucleus by a spring. The potential energy is minimized with respect to the shell position at every step of the simulation (see below). Successful applications of shell models in GROMACS have been published for N_2 [28] and water [29].

3.5.1 Optimization of the shell positions

The force \mathbf{F}_S on a shell particle S can be decomposed into two components:

$$\mathbf{F}_S = \mathbf{F}_{bond} + \mathbf{F}_{nb} \quad (3.42)$$

where \mathbf{F}_{bond} denotes the component representing the polarization energy, usually represented by a harmonic potential and \mathbf{F}_{nb} is the sum of Coulomb and Van der Waals interactions. If we assume that \mathbf{F}_{nb} is almost constant we can analytically derive the optimal position of the shell, i.e. where $\mathbf{F}_S = 0$. If we have the shell S connected to atom A we have

$$\mathbf{F}_{bond} = k_b (\mathbf{x}_S - \mathbf{x}_A) \quad (3.43)$$

In an iterative solver, we have positions $\mathbf{x}_S(n)$ where n is the iteration count. We now have it iteration n :

$$\mathbf{F}_{nb} = \mathbf{F}_S - k_b (\mathbf{x}_S(n) - \mathbf{x}_A) \quad (3.44)$$

and the optimal position for the shells $\mathbf{x}_S(n+1)$ thus follows from

$$\mathbf{F}_S - k_b (\mathbf{x}_S(n) - \mathbf{x}_A) + k_b (\mathbf{x}_S(n+1) - \mathbf{x}_A) = 0 \quad (3.45)$$

if we write

$$\Delta \mathbf{x}_S = \mathbf{x}_S(n+1) - \mathbf{x}_S(n) \quad (3.46)$$

we finally obtain

$$\Delta \mathbf{x}_S = \mathbf{F}_S/k_b \quad (3.47)$$

which then yields the algorithm to compute the next trial in the optimization of shell positions:

$$\mathbf{x}_S(n+1) = \mathbf{x}_S(n) + \mathbf{F}_S/k_b \quad (3.48)$$

3.6 Constraint algorithms

Constraints can be imposed in GROMACS using LINCS (default) or the traditional SHAKE method.

3.6.1 SHAKE

The SHAKE [30] algorithm changes a set of unconstrained coordinates \mathbf{r}' to a set of coordinates \mathbf{r}'' that fulfill a list of distance constraints, using a set \mathbf{r} as reference:

$$\text{SHAKE}(\mathbf{r}' \rightarrow \mathbf{r}''; \mathbf{r})$$

This action is consistent with solving a set of Lagrange multipliers in the constrained equations of motion. SHAKE needs a *tolerance* TOL; it will continue until all constraints are satisfied within a *relative* tolerance TOL. An error message is given if SHAKE cannot reset the coordinates because the deviation is too large, or if a given number of iterations is surpassed.

Assume the equations of motion must fulfill K holonomic constraints, expressed as

$$\sigma_k(\mathbf{r}_1 \dots \mathbf{r}_N) = 0; \quad k = 1 \dots K \quad (3.49)$$

(e.g. $(\mathbf{r}_1 - \mathbf{r}_2)^2 - b^2 = 0$). Then the forces are defined as

$$-\frac{\partial}{\partial \mathbf{r}_i} \left(V + \sum_{k=1}^K \lambda_k \sigma_k \right) \quad (3.50)$$

where λ_k are Lagrange multipliers which must be solved to fulfill the constraint equations. The second part of this sum determines the *constraint forces* \mathbf{G}_i , defined by

$$\mathbf{G}_i = - \sum_{k=1}^K \lambda_k \frac{\partial \sigma_k}{\partial \mathbf{r}_i} \quad (3.51)$$

The displacement due to the constraint forces in the leap frog or Verlet algorithm is equal to $(\mathbf{G}_i/m_i)(\Delta t)^2$. Solving the Lagrange multipliers (and hence the displacements) requires the solution of a set of coupled equations of the second degree. These are solved iteratively by SHAKE. For the special case of rigid water molecules, that often make up more than 80% of the simulation system we have implemented the SETTLE algorithm [31] (sec. 5.4).

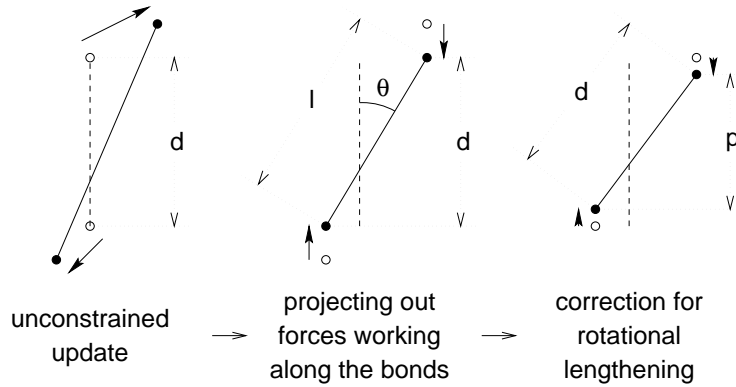


Figure 3.8: The three position updates needed for one time step. The dashed line is the old bond of length d , the solid lines are the new bonds. $l = d \cos \theta$ and $p = (2d^2 - l^2)^{\frac{1}{2}}$.

3.6.2 LINCS

The LINCS algorithm

LINCS is an algorithm that resets bonds to their correct lengths after an unconstrained update [32]. The method is non-iterative, as it always uses two steps. Although LINCS is based on matrices, no matrix-matrix multiplications are needed. The method is more stable and faster than SHAKE, but it can only be used with bond constraints and isolated angle constraints, such as the proton angle in OH. Because of its stability LINCS is especially useful for Brownian dynamics. LINCS has two parameters, which are explained in the subsection parameters.

The LINCS formulas

We consider a system of N particles, with positions given by a $3N$ vector $\mathbf{r}(t)$. For molecular dynamics the equations of motion are given by Newton's law

$$\frac{d^2 \mathbf{r}}{dt^2} = \mathbf{M}^{-1} \mathbf{F} \quad (3.52)$$

where \mathbf{F} is the $3N$ force vector and \mathbf{M} is a $3N \times 3N$ diagonal matrix, containing the masses of the particles. The system is constrained by K time-independent constraint equations

$$g_i(\mathbf{r}) = |\mathbf{r}_{i_1} - \mathbf{r}_{i_2}| - d_i = 0 \quad i = 1, \dots, K \quad (3.53)$$

In a numerical integration scheme LINCS is applied after an unconstrained update, just like SHAKE. The algorithm works in two steps (see figure Fig. 3.8). In the first step the projections of the new bonds on the old bonds are set to zero. In the second step a correction is applied for the lengthening of the bonds due to rotation. The numerics for the first step and the second step are very similar. A complete derivation of the algorithm can be found in [32]. Only a short description of the first step is given here.

A new notation is introduced for the gradient matrix of the constraint equations which appears on the right hand side of the equation

$$B_{hi} = \frac{\partial g_h}{\partial r_i} \quad (3.54)$$

Notice that \mathbf{B} is a $K \times 3N$ matrix, it contains the directions of the constraints. The following equation shows how the new constrained coordinates \mathbf{r}_{n+1} are related to the unconstrained coordinates \mathbf{r}_{n+1}^{unc}

$$\begin{aligned} \mathbf{r}_{n+1} &= (\mathbf{I} - \mathbf{T}_n \mathbf{B}_n) \mathbf{r}_{n+1}^{unc} + \mathbf{T}_n \mathbf{d} = \\ &\mathbf{r}_{n+1}^{unc} - \mathbf{M}^{-1} \mathbf{B}_n (\mathbf{B}_n \mathbf{M}^{-1} \mathbf{B}_n^T)^{-1} (\mathbf{B}_n \mathbf{r}_{n+1}^{unc} - \mathbf{d}) \end{aligned} \quad (3.55)$$

where $\mathbf{T} = \mathbf{M}^{-1} \mathbf{B}^T (\mathbf{B} \mathbf{M}^{-1} \mathbf{B}^T)^{-1}$. The derivation of this equation from eqns. 3.52 and 3.53 can be found in [32].

This first step does not set the real bond lengths to the prescribed lengths, but the projection of the new bonds onto the old directions of the bonds. To correct for the rotation of bond i , the projection of the bond on the old direction is set to

$$p_i = \sqrt{2d_i^2 - l_i^2} \quad (3.56)$$

where l_i is the bond length after the first projection. The corrected positions are

$$\mathbf{r}_{n+1}^* = (\mathbf{I} - \mathbf{T}_n \mathbf{B}_n) \mathbf{r}_{n+1} + \mathbf{T}_n \mathbf{p} \quad (3.57)$$

This correction for rotational effects is actually an iterative process, but during MD only one iteration is applied. The relative constraint deviation after this procedure will be less than 0.0001 for every constraint. In energy minimization this might not be accurate enough, so the number of iterations is equal to the order of the expansion (see below).

Half of the CPU time goes to inverting the constraint coupling matrix $\mathbf{B}_n \mathbf{M}^{-1} \mathbf{B}_n^T$, which has to be done every time step. This $K \times K$ matrix has $1/m_{i_1} + 1/m_{i_2}$ on the diagonal. The off-diagonal elements are only non-zero when two bonds are connected, then the element is $\cos \phi / m_c$, where m_c is the mass of the atom connecting the two bonds and ϕ is the angle between the bonds.

The matrix \mathbf{T} is inverted through a power expansion. A $K \times K$ matrix \mathbf{S} is introduced which is the inverse square root of the diagonal of $\mathbf{B}_n \mathbf{M}^{-1} \mathbf{B}_n^T$. This matrix is used to convert the diagonal elements of the coupling matrix to one

$$\begin{aligned} (\mathbf{B}_n \mathbf{M}^{-1} \mathbf{B}_n^T)^{-1} &= \mathbf{S} \mathbf{S}^{-1} (\mathbf{B}_n \mathbf{M}^{-1} \mathbf{B}_n^T)^{-1} \mathbf{S}^{-1} \mathbf{S} \\ &= \mathbf{S} (\mathbf{S} \mathbf{B}_n \mathbf{M}^{-1} \mathbf{B}_n^T \mathbf{S})^{-1} \mathbf{S} = \mathbf{S} (\mathbf{I} - \mathbf{A}_n)^{-1} \mathbf{S} \end{aligned} \quad (3.58)$$

The matrix \mathbf{A}_n is symmetric and sparse and has zeros on the diagonal. Thus a simple trick can be used to calculate the inverse

$$(\mathbf{I} - \mathbf{A}_n)^{-1} = \mathbf{I} + \mathbf{A}_n + \mathbf{A}_n^2 + \mathbf{A}_n^3 + \dots \quad (3.59)$$

This inversion method is only valid if the absolute values of all the eigenvalues of \mathbf{A}_n are smaller than one. In molecules with only bond constraints the connectivity is so low that this will always be true, even if ring structures are present. Problems can arise in angle-constrained molecules. By constraining angles with additional distance constraints multiple small ring structures are introduced. This gives a high connectivity, leading to large eigenvalues. Therefore LINCS should NOT be used with coupled angle-constraints.

The LINCS Parameters

The accuracy of LINCS depends on the number of matrices used in the expansion eqn. 3.59. For MD calculations a fourth order expansion is enough. For Brownian dynamics with large time steps an eighth order expansion may be necessary. The order is a parameter in the input file for `mdrun`. The implementation of LINCS is done in such a way that the algorithm will never crash. Even when it is impossible to reset the constraints LINCS will generate a conformation which fulfills the constraints as well as possible. However, LINCS will generate a warning when in one step a bond rotates over more than a predefined angle. This angle is set by the user in the input file for `mdrun`.

3.7 Simulated Annealing

The well known simulated annealing (SA) protocol is supported in GROMACS, and you can even couple multiple groups of atoms separately with an arbitrary number of reference temperatures that change during the simulation. The annealing is implemented by simply changing the current reference temperature for each group in the temperature coupling, so the actual relaxation and coupling properties depends on the type of thermostat you use and how hard you are coupling it. Since we are changing the reference temperature it is important to remember that the system will NOT instantaneously reach this value - you need to allow for the inherent relaxation time in the coupling algorithm too. If you are changing the annealing reference temperature faster than the temperature relaxation you will probably end up with a crash when the difference becomes too large.

The annealing protocol is specified as a series of corresponding times and reference temperatures for each group, and you can also choose whether you only want a single sequence (after which the temperature will be coupled to the last reference value), or if the annealing should be periodic and restart at the first reference point once the sequence is completed. You can mix and match both types of annealing and non-annealed groups in your simulation.

3.8 Stochastic Dynamics

Stochastic or velocity Langevin dynamics adds a friction and a noise term to Newton's equations of motion:

$$m_i \frac{d^2 \mathbf{r}_i}{dt^2} = -m_i \xi_i \frac{d\mathbf{r}_i}{dt} + \mathbf{F}_i(\mathbf{r}) + \dot{\mathbf{r}}_i \quad (3.60)$$

where ξ_i is the friction constant [1/ps] and $\dot{\mathbf{r}}_i(t)$ is a noise process with $\langle \dot{\mathbf{r}}_i(t) \dot{\mathbf{r}}_j(t+s) \rangle = 2m_i \xi_i k_B T \delta(s) \delta_{ij}$. When $1/\xi_i$ is large compared to the time scales present in the system, one could see stochastic dynamics as molecular dynamics with stochastic temperature-coupling. The advantage compared to MD with Berendsen temperature-coupling is that in case of SD the generated ensemble is known. For simulating a system in vacuum there is the additional advantage that there is no accumulation of errors for the overall translational and rotational degrees of freedom. When $1/\xi_i$ is small compared to the time scales present in the system, the dynamics will be completely different from MD, but the sampling is still correct.

GROMACS uses a complicated third-order leap-frog algorithm [33] to integrate equation (3.60). When constraints are present in the system, two constraint steps are performed per time step. Exact continuation of a stochastic dynamics simulation is not possible, since apart from the coordinates and the velocities one random term of the previous step is required, however, the error will be very small.

3.9 Brownian Dynamics

In the limit of high friction stochastic dynamics reduces to Brownian dynamics, also called position Langevin dynamics. This applies to over-damped systems, *i.e.* systems in which the inertia effects are negligible. The equation is:

$$\frac{d\mathbf{r}_i}{dt} = \frac{1}{\gamma_i} \mathbf{F}_i(\mathbf{r}) + \dot{\mathbf{r}}_i \quad (3.61)$$

where γ_i is the friction coefficient [amu/ps] and $\dot{\mathbf{r}}_i(t)$ is a noise process with $\langle \dot{\mathbf{r}}_i(t) \dot{\mathbf{r}}_j(t+s) \rangle = 2\delta(s)\delta_{ij}k_B T/\gamma_i$. In GROMACS the equations are integrated with a simple, explicit scheme:

$$\mathbf{r}_i(t + \Delta t) = \mathbf{r}_i(t) + \frac{\Delta t}{\gamma_i} \mathbf{F}_i(\mathbf{r}(t)) + \sqrt{2k_B T \frac{\Delta t}{\gamma_i}} \mathbf{r}_i^G \quad (3.62)$$

where \mathbf{r}_i^G is Gaussian distributed noise with $\mu = 0$, $\sigma = 1$. The friction coefficients γ_i can be chosen the same for all particles or as $\gamma_i = m_i/\xi_i$, where the friction constants ξ_i can be different for different groups of atoms. Because the system is assumed to be over damped, large time-steps can be used. LINCS should be used for the constraints since SHAKE will not converge for large atomic displacements. BD is an option of the `mdrun` program.

3.10 Energy Minimization

Energy minimization in GROMACS can be done using steepest descent, conjugate gradients, or l-bfgs (limited-memory Broyden-Fletcher-Goldfarb-Shanno quasi-Newtonian minimizer... we prefer the abbreviation). EM is just an option of the `mdrun` program.

3.10.1 Steepest Descent

Although steepest descent is certainly not the most efficient algorithm for searching, it is robust and easy to implement.

We define the vector \mathbf{r} as the vector of all $3N$ coordinates. Initially a maximum displacement h_0 (*e.g.* 0.01 nm) must be given.

First the forces \mathbf{F} and potential energy are calculated. New positions are calculated by

$$\mathbf{r}_{n+1} = \mathbf{r}_n + \frac{\mathbf{F}_n}{\max(|\mathbf{F}_n|)} h_n \quad (3.63)$$

where h_n is the maximum displacement and \mathbf{F}_n is the force, or the negative gradient of the potential V . The notation $\max(|\mathbf{F}_n|)$ means the largest of the absolute values of the force components. The forces and energy are again computed for the new positions

If ($V_{n+1} < V_n$) the new positions are accepted and $h_{n+1} = 1.2h_n$.

If ($V_{n+1} \geq V_n$) the new positions are rejected and $h_n = 0.2h_n$.

The algorithm stops when either a user specified number of force evaluations has been performed (*e.g.* 100), or when the maximum of the absolute values of the force (gradient) components is smaller than a specified value ϵ . Since force truncation produces some noise in the energy evaluation, the stopping criterion should not be made too tight to avoid endless iterations. A reasonable value for ϵ can be estimated from the root mean square force f a harmonic oscillator would exhibit at a temperature T . This value is

$$f = 2\pi\nu\sqrt{2mkT} \quad (3.64)$$

where ν is the oscillator frequency, m the (reduced) mass, and k Boltzmann's constant. For a weak oscillator with a wave number of 100 cm^{-1} and a mass of 10 atomic units, at a temperature of 1 K, $f = 7.7 \text{ kJ mol}^{-1} \text{ nm}^{-1}$. A value for ϵ between 1 and 10 is acceptable.

3.10.2 Conjugate Gradient

Conjugate gradient is slower than steepest descent in the early stages of the minimization, but becomes more efficient closer to the energy minimum. The parameters and stop criterion are the same as for steepest descent. In GROMACS conjugate gradient can not be used with constraints, including the SETTLE algorithm for water [31], as this has not been implemented. If water is present it must be of a flexible model, which can be specified in the mdp file by `define = -DFLEXIBLE`

This is not really a restriction, since the accuracy of conjugate gradient is only required for minimization prior to a normal mode analysis, which can not be performed with constraints. For most other purposes steepest descent is efficient enough.

3.10.3 L-BFGS

The original BFGS algorithm works by successively creating better approximations of the inverse Hessian matrix, and moving the system to the currently estimated minimum. The memory requirements for this are proportional to the square of the number of particles, so it is not practical for large systems like biomolecules. Instead, we use the L-BFGS algorithm of Nocedal, which approximates the inverse Hessian by a fixed number of corrections from previous steps. This sliding-window technique is almost as efficient as the original method, but the memory requirements are much lower - proportional to the number of particles multiplied with the correction steps. In practice we have found it to converge faster than conjugate gradients, but due to the correction steps it is not yet parallelized. It is also noteworthy that switched or shifted interactions usually improve the convergence, since sharp cut-offs means the potential function at the current coordinates is slightly different from the previous steps used to build the inverse Hessian approximation.

3.11 Normal Mode Analysis

Normal mode analysis [34, 35, 36] can be performed using GROMACS, by diagonalization of the mass-weighted Hessian H :

$$R^T M^{-1/2} H M^{-1/2} R = \text{diag}(\lambda_1, \dots, \lambda_{3N}) \quad (3.65)$$

$$\lambda_i = (2\pi\omega_i)^2 \quad (3.66)$$

where M contains the atomic masses, R is a matrix that contains the eigenvectors as columns, λ_i are the eigenvalues and ω_i are the corresponding frequencies.

First the Hessian matrix, which is a $3N \times 3N$ matrix where N is the number of atoms, needs to be calculated:

$$H_{ij} = \frac{\partial^2 V}{\partial x_i \partial x_j} \quad (3.67)$$

where x_i and x_j denote the atomic x, y or z coordinates. In practice, this equation is not used, but the Hessian is calculated numerically from the force as:

$$H_{ij} = -\frac{f_i(\mathbf{x} + h\mathbf{e}_j) - f_i(\mathbf{x} - h\mathbf{e}_j)}{2h} \quad (3.68)$$

$$f_i = -\frac{\partial V}{\partial x_i} \quad (3.69)$$

where \mathbf{e}_j is the unit vector in direction j . It should be noted that for a usual Normal Mode calculation, it is necessary to completely minimize the energy prior to computation of the Hessian. What tolerance is required depends on the type of system, but a rough indication is $0.001 \text{ kJ mol}^{-1}$. This should be done with conjugate gradients or l-bfgs in double precision.

A number of GROMACS programs are involved in these calculations. First the energy should be minimized using `mdrun`. Then `mdrun` computes the Hessian, note that for generating the run input file one should use the minimized conformation from the full precision trajectory file, as the structure file is not accurate enough. `gnmeig` does the diagonalization and the sorting of the normal modes according to their frequencies. Both `mdrun` and `gnmeig` should be run in double precision. The normal modes can be analyzed with the program `ganaeig`. Ensembles of structures at any temperature and for any subset of normal modes can be generated with `gnmens`. An overview of normal mode analysis and the related principal component analysis (see sec. 8.10) can be found in [37].

3.12 Free energy calculations

Free energy calculations can be performed in GROMACS using slow-growth methods. An example problem might be: calculate the difference in free energy of binding of an inhibitor **I** to an enzyme **E** and to a mutated enzyme **E'**. It is not feasible with computer simulations to perform a docking calculation for such a large complex, or even releasing the inhibitor from the enzyme in a reasonable amount of computer time with reasonable accuracy. However, if we consider the free energy cycle in (Fig. 3.9A) we can write

$$\Delta G_1 - \Delta G_2 = \Delta G_3 - \Delta G_4 \quad (3.70)$$

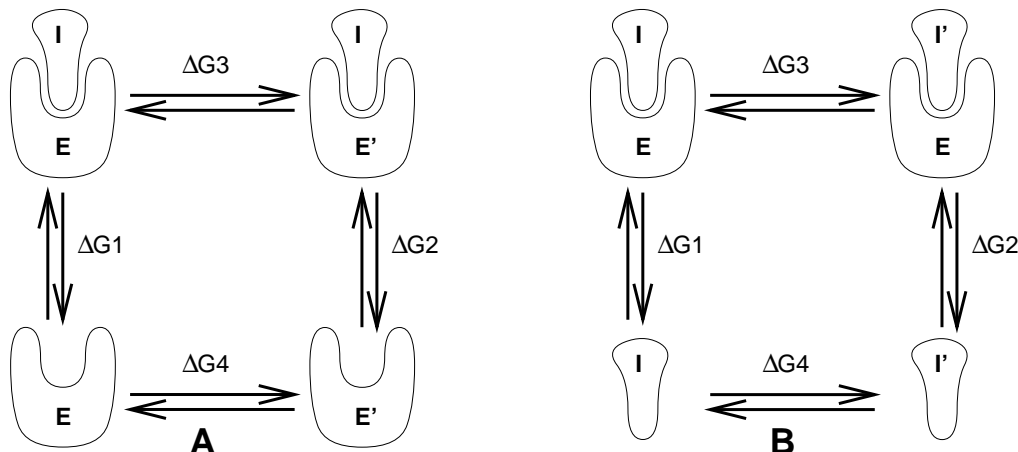


Figure 3.9: Free energy cycles. **A**: to calculate ΔG_{12} , the free energy difference between the binding of inhibitor **I** to enzymes **E** respectively **E'**. **B**: to calculate ΔG_{12} , the free energy difference for binding of inhibitors **I** respectively **I'** to enzyme **E**.

If we are interested in the left-hand term we can equally well compute the right-hand term.

If we want to compute the difference in free energy of binding of two inhibitors **I** and **I'** to an enzyme **E** (Fig. 3.9B) we can again use eqn. 3.70 to compute the desired property.

Free energy differences between two molecular species can be calculated in GROMACS using the “slow-growth” method. In fact, such free energy differences between different molecular species are physically meaningless, but they can be used to obtain meaningful quantities employing a thermodynamic cycle. The method requires a simulation during which the Hamiltonian of the system changes slowly from that describing one system (A) to that describing the other system (B). The change must be so slow that the system remains in equilibrium during the process; if that requirement is fulfilled, the change is reversible and a slow-growth simulation from B to A will yield the same results (but with a different sign) as a slow-growth simulation from A to B. This is a useful check, but the user should be aware of the danger that equality of forward and backward growth results does not guarantee correctness of the results.

The required modification of the Hamiltonian H is realized by making H a function of a *coupling parameter* λ : $H = H(p, q; \lambda)$ in such a way that $\lambda = 0$ describes system A and $\lambda = 1$ describes system B:

$$H(p, q; 0) = H^A(p, q); \quad H(p, q; 1) = H^B(p, q). \quad (3.71)$$

In GROMACS, the functional form of the λ -dependence is different for the various force-field contributions and is described in section sec. 4.5.

The Helmholtz free energy A is related to the partition function Q of an N, V, T ensemble, which is assumed to be the equilibrium ensemble generated by a MD simulation at constant volume and temperature. The generally more useful Gibbs free energy G is related to the partition function Δ of an N, p, T ensemble, which is assumed to be the equilibrium ensemble generated by a MD simulation at constant pressure and temperature:

$$A(\lambda) = -k_B T \ln Q \quad (3.72)$$

$$Q = c \int \int \exp[-\beta H(p, q; \lambda)] dp dq \quad (3.73)$$

$$G(\lambda) = -k_B T \ln \Delta \quad (3.74)$$

$$\Delta = c \int \int \int \exp[-\beta H(p, q; \lambda) - \beta pV] dp dq dV \quad (3.75)$$

$$G = A + pV, \quad (3.76)$$

where $\beta = 1/(k_B T)$ and $c = (N!h^{3N})^{-1}$. These integrals over phase space cannot be evaluated from a simulation, but it is possible to evaluate the derivative with respect to λ as an ensemble average:

$$\frac{dA}{d\lambda} = \frac{\int \int (\partial H / \partial \lambda) \exp[-\beta H(p, q; \lambda)] dp dq}{\int \int \exp[-\beta H(p, q; \lambda)] dp dq} = \left\langle \frac{\partial H}{\partial \lambda} \right\rangle_{NVT; \lambda}, \quad (3.77)$$

with a similar relation for $dG/d\lambda$ in the N, p, T ensemble. The difference in free energy between A and B can be found by integrating the derivative over λ :

$$A^B(V, T) - A^A(V, T) = \int_0^1 \left\langle \frac{\partial H}{\partial \lambda} \right\rangle_{NVT; \lambda} d\lambda \quad (3.78)$$

$$G^B(p, T) - G^A(p, T) = \int_0^1 \left\langle \frac{\partial H}{\partial \lambda} \right\rangle_{NpT; \lambda} d\lambda. \quad (3.79)$$

If one wishes to evaluate $G^B(p, T) - G^A(p, T)$, the natural choice is a constant-pressure simulation. However, this quantity can also be obtained from a slow-growth simulation at constant volume, starting with system A at pressure p and volume V and ending with system B at pressure p_B , by applying the following small correction:

$$G^B(p) - G^A(p) = A^B(V) - A^A(V) - \int_p^{p_B} [V^B(p') - V] dp' \quad (3.80)$$

Here we omitted the constant T from the notation. This correction is roughly equal to $-\frac{1}{2}(p^B - p)\Delta V = (\Delta V)^2/(2\kappa V)$, where ΔV is the volume change at p and κ is the isothermal compressibility. This is usually negligible. For example, the growth of a water molecule from nothing in a bath of 1000 water molecules at constant volume would produce an additional pressure of 22 bar and a correction to the Helmholtz free energy of -20 J/mol.

In cartesian coordinates, the kinetic energy term in the Hamiltonian depends only on the momenta, and can be separately integrated and in fact removed from the equations. When masses do not change, there is no contribution from the kinetic energy at all; otherwise the integrated contribution to the free energy is $-\frac{3}{2}k_B T \ln(m^B/m^A)$. This is no longer true in the presence of constraints.

GROMACS offers the possibility to integrate eq. 3.78 or eq. 3.79 in one simulation over the full range from A to B. However, if the change is large and sampling insufficiency can be expected, the user may prefer to determine the value of $\langle dG/d\lambda \rangle$ accurately at a number of well-chosen intermediate values of λ . This can be easily done by setting the stepsize **delta lambda** to zero. Each simulation can be equilibrated first, and a proper error estimate can be made for each value of $dG/d\lambda$ from the fluctuation of $\partial H/\partial \lambda$. The total free energy change is then determined afterwards by an appropriate numerical integration procedure.

The λ -dependence for the force-field contributions is described in section sec. 4.5.

3.13 Replica exchange

Replica exchange molecular dynamics (REMD) is a method which can be used to speed up the sampling of any type of simulation, especially if conformations are separated by relatively high energy barriers. It involves simulating multiple replicas of the same system at different temperatures and randomly exchanging the complete state of two replicas at regular intervals with the probability:

$$P(1 \leftrightarrow 2) = \min \left(1, \exp \left[\left(\frac{1}{k_B T_1} - \frac{1}{k_B T_2} \right) (U_1 - U_2) \right] \right) \quad (3.81)$$

where T_1 and T_2 are the reference temperatures and U_1 and U_2 are the instantaneous potential energies of replicas 1 and 2 respectively. After exchange the velocities are scaled by $(T_1/T_2)^{\pm 0.5}$ and a neighbor search is performed the next step. This combines the fast sampling and frequent barrier-crossing of the highest temperature with correct Boltzmann sampling at all the different temperatures [38, 39]. We only attempt exchanges for neighboring temperatures as the probability decreases very rapidly with the temperature difference. One should not attempt exchanges for all possible pairs in one step. If, for instance, replicas 1 and 2 would exchange, the chance of exchange for replicas 2 and 3 not only depends on the energies of replicas 2 and 3, but also on the energy of replica 1. In GROMACS this is solved by attempting exchange for all 'odd' pairs on 'odd' attempts and for all 'even' pairs on 'even' attempts. If we have four replicas: 0, 1, 2 and 3, ordered in temperature and we attempt exchange every 1000 steps, pairs 0-1 and 2-3 will be tried at steps 1000, 3000 etc. and pair 1-2 at steps 2000, 4000 etc.

How should one choose the temperatures? The energy difference can be written as:

$$U_1 - U_2 = N_{df} \frac{c}{2} k_B (T_1 - T_2) \quad (3.82)$$

where N_{df} is the total number of degrees of freedom of one replica and c is 1 for harmonic potentials and around 2 for protein/water systems. If $T_2 = (1 + \epsilon)T_1$ the probability becomes:

$$P(1 \leftrightarrow 2) = \exp \left(-\frac{\epsilon^2 c N_{df}}{2(1 + \epsilon)} \right) \approx \exp \left(-\epsilon^2 \frac{c}{2} N_{df} \right) \quad (3.83)$$

Thus for a probability of $e^{-2} \approx 0.135$ one obtains $\epsilon \approx 2/\sqrt{c N_{df}}$. With all bonds constrained one has $N_{df} \approx 2 N_{atoms}$ and thus for $c = 2$ one should choose ϵ as $1/\sqrt{N_{atoms}}$. However there is one problem when using pressure coupling. The density at higher temperatures will decrease, leading to higher energy [40] and this should be taken into account. The GROMACS website features a so-called "REMD" - calculator, that lets you type in the temperature range and the number of atoms, and based on that proposes a set of temperatures.

An extension to the REMD for the isobaric-isothermal ensemble was proposed by Okabe *et al.* [41]. In this work the exchange probability is modified to:

$$P(1 \leftrightarrow 2) = \min \left(1, \exp \left[\left(\frac{1}{k_B T_1} - \frac{1}{k_B T_2} \right) (U_1 - U_2) + \left(\frac{P_1}{k_B T_1} - \frac{P_2}{k_B T_2} \right) (V_1 - V_2) \right] \right) \quad (3.84)$$

where P_1 and P_2 are the respective reference pressures and V_1 and V_2 are the respective instantaneous volumes in the simulations. In most cases the differences in volume are so small that the second term is negligible. It only plays a role when the difference between P_1 and P_2 is large or in phase transitions.

Replica exchange is an option of the `mdrun` program. It will only work when MPI is installed, due to the inherent parallelism in the algorithm. For efficiency each replica can run on a separate node. See the manual page of `mdrun` on how to use it.

3.14 Essential Dynamics Sampling

The results from Essential Dynamics (see sec. 8.10) of a protein can be used to guide MD simulations. The idea is that from an initial MD simulation (or from other sources) a definition of the collective fluctuations with largest amplitude is obtained. The position along one or more of these collective modes can be constrained in a (second) MD simulation in a number of ways for several purposes. For example, the position along a certain mode may be kept fixed to monitor the average force (free-energy gradient) on that coordinate in that position. Another application is to enhance sampling efficiency with respect to usual MD [42, 43]. In this case, the system is encouraged to sample its available configuration space more systematically than in a diffusion-like path that proteins usually take.

All available constraint types are described in the appropriate chapter of the WHAT IF [44] manual.

Chapter 4

Interaction function and force field

A force field is built up from two distinct components:

- The set of equations (called the *potential functions*) used to generate the potential energies and their derivatives, the forces.
- The parameters used in this set of equations

Within one set of equations various sets of parameters can be used. Care must be taken that the combination of equations and parameters form a consistent set. It is in general dangerous to make *ad hoc* changes in a subset of parameters, because the various contributions to the total force are usually interdependent.

GROMACS 3.3 includes several force fields, and additional ones are available on the website. If you don't know which one to select we recommend Gromos96 for united-atom setups and OPLS-AA/L for all-atom parameters. The GROMACS forcefields is mainly included for historical reasons; it is based on GROMOS-87 [45], with a small modification concerning the interaction between water-oxygens and carbon atoms [46, 47], as well as 10 extra atom types [48, 49, 46, 47, 50]. However, the user is free to make her own modifications (beware!). This will be explained in details in chapter 5, which deals with the **Topology**. The organization of the force field files is described in sec. 5.7.

To accommodate the potential functions used in some popular force fields, GROMACS offers a choice of functions, both for non-bonded interaction and for dihedral interactions. They are described in the appropriate subsections.

The potential functions can be subdivided into three parts

1. *Non-bonded*: Lennard-Jones or Buckingham, and Coulomb or modified Coulomb. The non-bonded interactions are computed on the basis of a neighbor list (a list of non-bonded atoms within a certain radius), in which exclusions are already removed.

2. *Bonded*: covalent bond-stretching, angle-bending, improper dihedrals, and proper dihedrals. These are computed on the basis of fixed lists.
3. *Restraints*: position restraints, angle restraints, distance restraints, orientation restraints and dihedral restraints, all based on fixed lists.

4.1 Non-bonded interactions

Non-bonded interactions in GROMACS are pair-additive and centro-symmetric:

$$V(\mathbf{r}_1, \dots, \mathbf{r}_N) = \sum_{i < j} V_{ij}(\mathbf{r}_{ij}); \quad (4.1)$$

$$\mathbf{F}_i = - \sum_j \frac{dV_{ij}(r_{ij})}{dr_{ij}} \frac{\mathbf{r}_{ij}}{r_{ij}} = -\mathbf{F}_j \quad (4.2)$$

The non-bonded interactions contain a repulsion term, a dispersion term, and a Coulomb term. The repulsion and dispersion term are combined in either the Lennard-Jones (or 6-12 interaction), or the Buckingham (or exp-6 potential). In addition, (partially) charged atoms act through the Coulomb term.

4.1.1 The Lennard-Jones interaction

The Lennard-Jones potential V_{LJ} between two atoms equals

$$V_{LJ}(r_{ij}) = \frac{C_{ij}^{(12)}}{r_{ij}^{12}} - \frac{C_{ij}^{(6)}}{r_{ij}^6} \quad (4.3)$$

see also Fig. 4.1 The parameters $C_{ij}^{(12)}$ and $C_{ij}^{(6)}$ depend on pairs of *atom types*; consequently they are taken from a matrix of LJ-parameters.

The force derived from this potential is:

$$\mathbf{F}_i(\mathbf{r}_{ij}) = \left(12 \frac{C_{ij}^{(12)}}{r_{ij}^{13}} - 6 \frac{C_{ij}^{(6)}}{r_{ij}^7} \right) \frac{\mathbf{r}_{ij}}{r_{ij}} \quad (4.4)$$

The LJ potential may also be written in the following form :

$$V_{LJ}(\mathbf{r}_{ij}) = 4\epsilon_{ij} \left(\left(\frac{\sigma_{ij}}{r_{ij}} \right)^{12} - \left(\frac{\sigma_{ij}}{r_{ij}} \right)^6 \right) \quad (4.5)$$

In constructing the parameter matrix for the non-bonded LJ-parameters, two types of combination rules can be used within GROMACS:

$$\begin{aligned} C_{ij}^{(6)} &= \left(C_{ii}^{(6)} * C_{jj}^{(6)} \right)^{1/2} \\ C_{ij}^{(12)} &= \left(C_{ii}^{(12)} * C_{jj}^{(12)} \right)^{1/2} \end{aligned} \quad (4.6)$$

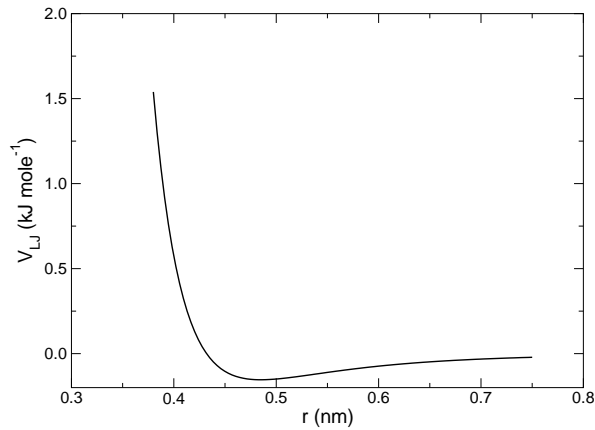


Figure 4.1: The Lennard-Jones interaction.

or, alternatively the Lorentz-Bertelot rules can be used. An arithmetic average is used for the sigma's, while a geometric average is used for the epsilon's,

$$\begin{aligned}\sigma_{ij} &= \frac{1}{2}(\sigma_{ii} + \sigma_{jj}) \\ \epsilon_{ij} &= (\epsilon_{ii}\epsilon_{jj})^{1/2}\end{aligned}\quad (4.7)$$

4.1.2 Buckingham potential

The Buckingham potential has a more flexible and realistic repulsion term than the Lennard-Jones interaction, but is also more expensive to compute. The potential form is:

$$V_{bh}(r_{ij}) = A_{ij} \exp(-B_{ij}r_{ij}) - \frac{C_{ij}}{r_{ij}^6} \quad (4.8)$$

see also Fig. 4.2, the force derived from this is:

$$\mathbf{F}_i(r_{ij}) = \left[A_{ij}B_{ij} \exp(-B_{ij}r_{ij}) - 6\frac{C_{ij}}{r_{ij}^7} \right] \frac{\mathbf{r}_{ij}}{r_{ij}} \quad (4.9)$$

4.1.3 Coulomb interaction

The Coulomb interaction between two charge particles is given by:

$$V_c(r_{ij}) = f \frac{q_i q_j}{\epsilon_r r_{ij}} \quad (4.10)$$

see also Fig. 4.3, where $f = \frac{1}{4\pi\epsilon_0} = 138.935485$ (see chapter 2)

The force derived from this potential is:

$$\mathbf{F}_i(r_{ij}) = f \frac{q_i q_j}{\epsilon_r r_{ij}^2} \frac{\mathbf{r}_{ij}}{r_{ij}} \quad (4.11)$$

In GROMACS the relative dielectric constant ϵ_r may be set in the in the input for `grompp`.

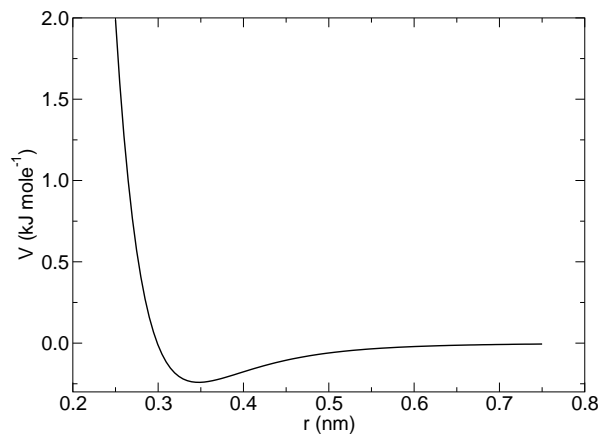


Figure 4.2: The Buckingham interaction.

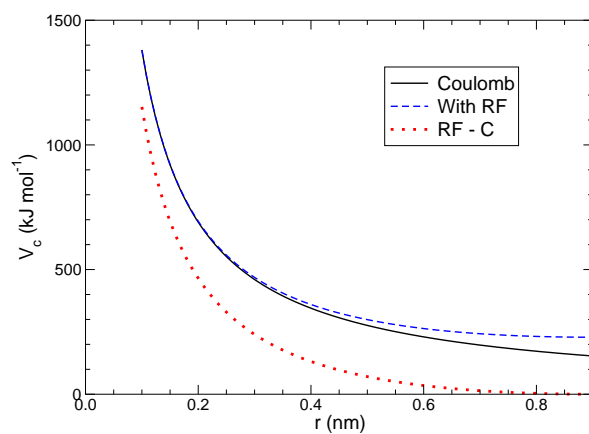


Figure 4.3: The Coulomb interaction (for particles with equal signed charge) with and without reaction field. In the latter case ϵ_r was 1, ϵ_{rf} was 78, and r_c was 0.9 nm. The dot-dashed line is the same as the dashed line, except for a constant.

4.1.4 Coulomb interaction with reaction field

The coulomb interaction can be modified for homogeneous systems, by assuming a constant dielectric environment beyond the cutoff r_c with a dielectric constant of ε_{rf} . The interaction then reads:

$$V_{crf} = f \frac{q_i q_j}{\varepsilon_r r_{ij}} \left[1 + \frac{\varepsilon_{rf} - \varepsilon_r}{2\varepsilon_{rf} + \varepsilon_r} \frac{r_{ij}^3}{r_c^3} \right] - f \frac{q_i q_j}{\varepsilon_r r_c} \frac{3\varepsilon_{rf}}{2\varepsilon_{rf} + \varepsilon_r} \quad (4.12)$$

in which the constant expression on the right makes the potential zero at the cutoff r_c . For charged cut-off spheres this corresponds to neutralization with a homogeneous background charge. We can rewrite eqn. 4.12 for simplicity as

$$V_{crf} = f \frac{q_i q_j}{\varepsilon_r} \left[\frac{1}{r_{ij}} + k_{rf} r_{ij}^2 - c_{rf} \right] \quad (4.13)$$

with

$$k_{rf} = \frac{1}{r_c^3} \frac{\varepsilon_{rf} - \varepsilon_r}{(2\varepsilon_{rf} + \varepsilon_r)} \quad (4.14)$$

$$c_{rf} = \frac{1}{r_c} + k_{rf} r_c^2 = \frac{1}{r_c} \frac{3\varepsilon_{rf}}{(2\varepsilon_{rf} + \varepsilon_r)} \quad (4.15)$$

For large ε_{rf} the k_{rf} goes to $r_c^{-3}/2$, while for $\varepsilon_{rf} = \varepsilon_r$ the correction vanishes. In Fig. 4.3 the modified interaction is plotted, and it is clear that the derivative with respect to r_{ij} (= -force) goes to zero at the cutoff distance. The force derived from this potential reads:

$$\mathbf{F}_i(\mathbf{r}_{ij}) = f \frac{q_i q_j}{\varepsilon_r} \left[\frac{1}{r_{ij}^2} - 2k_{rf} r_{ij} \right] \frac{\mathbf{r}_{ij}}{r_{ij}} \quad (4.16)$$

The reaction-field correction should also be applied to all excluded atoms pairs, including self pairs, in which case the normal Coulomb term in eqns. 4.12 and 4.16 is absent.

Tironi *et al.* have introduced a generalized reaction field in which the dielectric continuum beyond the cutoff r_c also has an ionic strength I [51]. In this case we can rewrite the constants k_{rf} and c_{rf} using the inverse Debye screening length κ :

$$\kappa^2 = \frac{2I F^2}{\varepsilon_0 \varepsilon_{rf} R T} = \frac{F^2}{\varepsilon_0 \varepsilon_{rf} R T} \sum_{i=1}^K c_i z_i \quad (4.17)$$

$$k_{rf} = \frac{1}{r_c^3} \frac{(\varepsilon_{rf} - \varepsilon_r)(1 + \kappa r_c) + \frac{1}{2} \varepsilon_{rf} (\kappa r_c)^2}{(2\varepsilon_{rf} + \varepsilon_r)(1 + \kappa r_c) + \varepsilon_{rf} (\kappa r_c)^2} \quad (4.18)$$

$$c_{rf} = \frac{1}{r_c} \frac{3\varepsilon_{rf}(1 + \kappa r_c + \frac{1}{2}(\kappa r_c)^2)}{(2\varepsilon_{rf} + \varepsilon_r)(1 + \kappa r_c) + \varepsilon_{rf} (\kappa r_c)^2} \quad (4.19)$$

where F is Faraday's constant, R is the ideal gas constant, T the absolute temperature, c_i the molar concentration for species i and z_i the charge number of species i where we have K different species. In the limit of zero ionic strength ($\kappa = 0$) eqns. 4.18 and 4.19 reduce to the simple forms of eqns. 4.14 and 4.15 respectively.

4.1.5 Modified non-bonded interactions

In the GROMACS force field the non-bonded potentials can be modified by a shift function. The purpose of this is to replace the truncated forces by forces that are continuous and have continuous derivatives at the cutoff radius. With such forces the time-step integration produces much smaller errors and there are no such complications as creating charges from dipoles by the truncation procedure. In fact, by using shifted forces there is no need for charge groups in the construction of neighbor lists. However, the shift function produces a considerable modification of the Coulomb potential. Unless the 'missing' long-range potential is properly calculated and added (through the use of PPPM, Ewald, or PME), the effect of such modifications must be carefully evaluated. The modification of the Lennard-Jones dispersion and repulsion is only minor, but it does remove the noise caused by cutoff effects.

There is *no* fundamental difference between a switch function (which multiplies the potential with a function) and a shift function (which adds a function to the force or potential) [52]. The switch function is a special case of the shift function, which we apply to the *force function* $F(r)$, related to the electrostatic or Van der Waals force acting on particle i by particle j as

$$\mathbf{F}_i = cF(r_{ij})\frac{\mathbf{r}_{ij}}{r_{ij}} \quad (4.20)$$

For pure Coulomb or Lennard-Jones interactions $F(r) = F_\alpha(r) = r^{-(\alpha+1)}$. The shifted force $F_s(r)$ can generally be written as:

$$\begin{aligned} F_s(r) &= F_\alpha(r) & r < r_1 \\ F_s(r) &= F_\alpha(r) + S(r) & r_1 \leq r < r_c \\ F_s(r) &= 0 & r_c \leq r \end{aligned} \quad (4.21)$$

When $r_1 = 0$ this is a traditional shift function, otherwise it acts as a switch function. The corresponding shifted coulomb potential then reads:

$$V_s(r_{ij}) = f\Phi_s(r_{ij})q_iq_j \quad (4.22)$$

where $\Phi(r)$ is the potential function

$$\Phi_s(r) = \int_r^\infty F_s(x) dx \quad (4.23)$$

The GROMACS shift function should be smooth at the boundaries, therefore the following boundary conditions are imposed on the shift function:

$$\begin{aligned} S(r_1) &= 0 \\ S'(r_1) &= 0 \\ S(r_c) &= -F_\alpha(r_c) \\ S'(r_c) &= -F'_\alpha(r_c) \end{aligned} \quad (4.24)$$

A 3^{rd} degree polynomial of the form

$$S(r) = A(r - r_1)^2 + B(r - r_1)^3 \quad (4.25)$$

fulfills these requirements. The constants A and B are given by the boundary condition at r_c :

$$\begin{aligned} A &= -\frac{(\alpha + 4)r_c - (\alpha + 1)r_1}{r_c^{\alpha+2} (r_c - r_1)^2} \\ B &= \frac{(\alpha + 3)r_c - (\alpha + 1)r_1}{r_c^{\alpha+2} (r_c - r_1)^3} \end{aligned} \quad (4.26)$$

Thus the total force function is

$$F_s(r) = \frac{\alpha}{r^{\alpha+1}} + A(r - r_1)^2 + B(r - r_1)^3 \quad (4.27)$$

and the potential function reads

$$\Phi(r) = \frac{1}{r^\alpha} - \frac{A}{3}(r - r_1)^3 - \frac{B}{4}(r - r_1)^4 - C \quad (4.28)$$

where

$$C = \frac{1}{r_c^\alpha} - \frac{A}{3}(r_c - r_1)^3 - \frac{B}{4}(r_c - r_1)^4 \quad (4.29)$$

When $r_1 = 0$, the modified Coulomb force function is

$$F_s(r) = \frac{1}{r^2} - \frac{5r^2}{r_c^4} + \frac{4r^3}{r_c^5} \quad (4.30)$$

identical to the *parabolic force* function recommended to be used as a short-range function in conjunction with a Poisson solver for the long-range part [17]. The modified Coulomb potential function is

$$\Phi(r) = \frac{1}{r} - \frac{5}{3r_c} + \frac{5r^3}{3r_c^4} - \frac{r^4}{r_c^5} \quad (4.31)$$

see also Fig. 4.4.

4.1.6 Modified short-range interactions with Ewald summation

When Ewald summation or particle-mesh Ewald is used to calculate the long-range interactions, the short-range coulomb potential must also be modified, similar to the switch function above. In this case the short range potential is given by

$$V(r) = f \frac{\operatorname{erfc}(\beta r_{ij})}{r_{ij}} q_i q_j, \quad (4.32)$$

where β is a parameter that determines the relative weight between the direct space sum and the reciprocal space sum and $\operatorname{erfc}(x)$ is the complementary error function. For further details on long-range electrostatics, see sec. 4.9.

4.2 Bonded interactions

Bonded interactions are based on a fixed list of atoms. They are not exclusively pair interactions, but include 3- and 4-body interactions as well. There are *bond stretching* (2-body), *bond angle* (3-body), and *dihedral angle* (4-body) interactions. A special type of dihedral interaction (called *improper dihedral*) is used to force atoms to remain in a plane or to prevent transition to a configuration of opposite chirality (a mirror image).

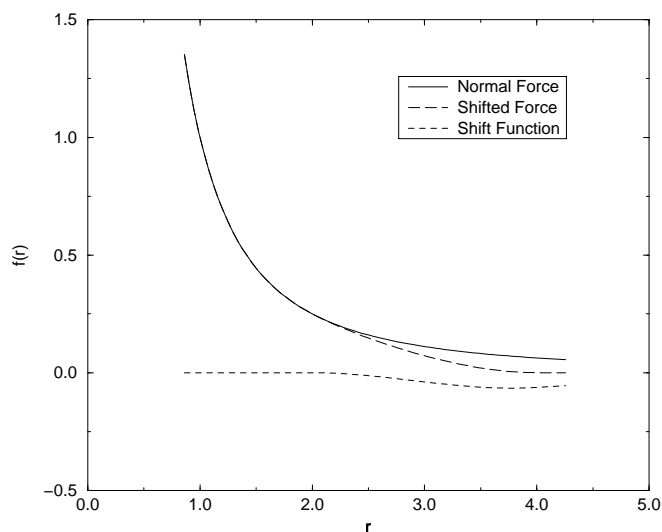


Figure 4.4: The Coulomb Force, Shifted Force and Shift Function $S(r)$, using $r_1 = 2$ and $r_c = 4$.

4.2.1 Bond stretching

Harmonic potential

The bond stretching between two covalently bonded atoms i and j is represented by a harmonic potential

$$V_b(r_{ij}) = \frac{1}{2}k_{ij}^b(r_{ij} - b_{ij})^2 \quad (4.33)$$

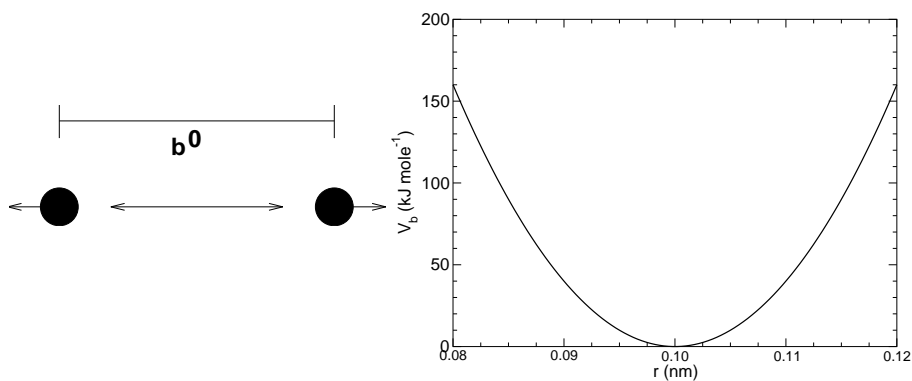


Figure 4.5: Principle of bond stretching (left), and the bond stretching potential (right).

see also Fig. 4.5, with the force

$$\mathbf{F}_i(\mathbf{r}_{ij}) = k_{ij}^b (r_{ij} - b_{ij}) \frac{\mathbf{r}_{ij}}{r_{ij}} \quad (4.34)$$

Fourth power potential

In the GROMOS-96 force field [53] the covalent bond potential is written for reasons of computational efficiency as:

$$V_b(r_{ij}) = \frac{1}{4} k_{ij}^b (r_{ij}^2 - b_{ij}^2)^2 \quad (4.35)$$

the corresponding force is:

$$\mathbf{F}_i(\mathbf{r}_{ij}) = k_{ij}^b (r_{ij}^2 - b_{ij}^2) \mathbf{r}_{ij} \quad (4.36)$$

The force constants for this form of the potential is related to the usual harmonic force constant $k^{b,harm}$ (sec. 4.2.1) as

$$2k^b b_{ij}^2 = k^{b,harm} \quad (4.37)$$

The force constants are mostly derived from the harmonic ones used in GROMOS-87 [45]. Although this form is computationally more efficient (because no square root has to be evaluated), it is conceptually more complex. One particular disadvantage is that since the form is not harmonic, the average energy of a single bond is not equal to $\frac{1}{2}kT$ as it is for the normal harmonic potential.

4.2.2 Morse potential bond stretching

For some systems that require an anharmonic bond stretching potential, the Morse potential [54] between two atoms i and j is available in GROMACS. This potential differs from the harmonic potential in having an asymmetric potential well and a zero force at infinite distance. The functional form is:

$$V_{morse}(r_{ij}) = D_{ij} [1 - \exp(-\beta_{ij}(r_{ij} - b_{ij}))]^2, \quad (4.38)$$

see also Fig. 4.6, and the corresponding force is:

$$\mathbf{F}_{morse}(\mathbf{r}_{ij}) = 2D_{ij}\beta_{ij}r_{ij} \exp(-\beta_{ij}(r_{ij} - b_{ij})) * [1 - \exp(-\beta_{ij}(r_{ij} - b_{ij}))] \frac{\mathbf{r}_{ij}}{r_{ij}}, \quad (4.39)$$

where D_{ij} is the depth of the well in kJ/mol, β_{ij} defines the steepness of the well (in nm^{-1}), and b_{ij} is the equilibrium distance in nm. The steepness parameter β_{ij} can be expressed in terms of the reduced mass of the atoms i and j , the fundamental vibration frequency ω_{ij} and the well depth D_{ij} :

$$\beta_{ij} = \omega_{ij} \sqrt{\frac{\mu_{ij}}{2D_{ij}}} \quad (4.40)$$

and because $\omega = \sqrt{k/\mu}$, one can rewrite β_{ij} in terms of the harmonic force constant k_{ij}

$$\beta_{ij} = \sqrt{\frac{k_{ij}}{2D_{ij}}} \quad (4.41)$$

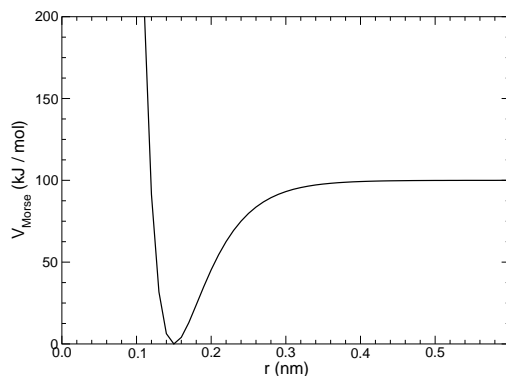


Figure 4.6: The Morse potential well, with bond length 0.15 nm.

For small deviations ($r_{ij} - b_{ij}$), one can approximate the exp-term to first-order using a Taylor expansion:

$$\exp(-x) \approx 1 - x \quad (4.42)$$

and substituting eqn. 4.41 and eqn. 4.42 in the functional from,

$$\begin{aligned} V_{morse}(r_{ij}) &= D_{ij}[1 - \exp(-\beta_{ij}(r_{ij} - b_{ij}))]^2 \\ &= D_{ij}[1 - (1 - \sqrt{\frac{k_{ij}}{2D_{ij}}}(r_{ij} - b_{ij}))]^2 \\ &= \frac{1}{2}k_{ij}(r_{ij} - b_{ij})^2, \end{aligned} \quad (4.43)$$

we recover the harmonic bond stretching potential.

4.2.3 Cubic bond stretching potential

Another anharmonic bond stretching potential that is slightly simpler than the Morse potential adds a cubic term in the distance to the simple harmonic form:

$$V_b(r_{ij}) = k_{ij}^b(r_{ij} - b_{ij})^2 + k_{ij}^b k_{ij}^{cub}(r_{ij} - b_{ij})^3 \quad (4.44)$$

A flexible water model (based on the SPC water model [55]) including a cubic bond stretching potential for the O-H bond was developed by Ferguson [56]. This model was found to yield a reasonable infrared spectrum. The Ferguson water model is available in the GROMACS library. It should be noted that the potential is asymmetric: overstretching leads to infinitely low energies. The integration timestep is therefore limited to 1 fs.

The force corresponding to this potential is:

$$\mathbf{F}_i(\mathbf{r}_{ij}) = 2k_{ij}^b(r_{ij} - b_{ij}) \frac{\mathbf{r}_{ij}}{r_{ij}} + 3k_{ij}^b k_{ij}^{cub}(r_{ij} - b_{ij})^2 \frac{\mathbf{r}_{ij}}{r_{ij}} \quad (4.45)$$

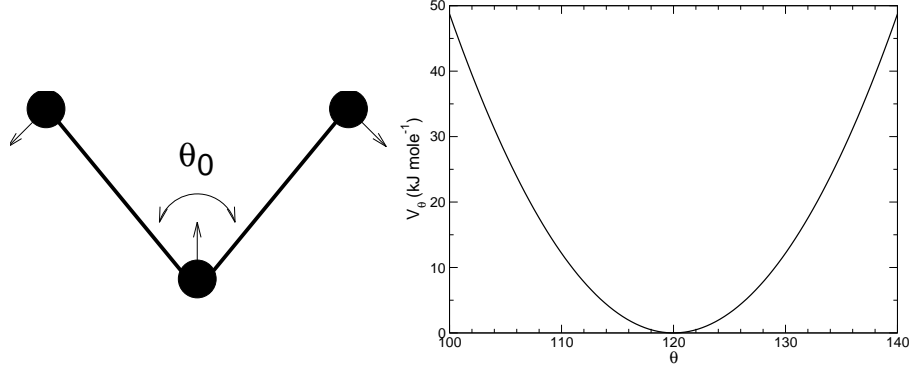


Figure 4.7: Principle of angle vibration (left) and the bond angle potential (right).

4.2.4 FENE bond stretching potential

In coarse-grained polymer simulations the beads are often connected by a FENE (finitely extensible nonlinear elastic) potential [57]:

$$V_{\text{FENE}}(r_{ij}) = -\frac{1}{2}k_{ij}^b b_{ij}^2 \log\left(1 - \frac{r_{ij}^2}{b_{ij}^2}\right) \quad (4.46)$$

The potential looks complicated, but the expression for the force is simpler:

$$\mathbf{F}_{\text{FENE}}(\mathbf{r}_{ij}) = k_{ij}^b \left(1 - \frac{r_{ij}^2}{b_{ij}^2}\right)^{-1} \mathbf{r}_{ij} \quad (4.47)$$

At short distances the potential asymptotically goes to a harmonic potential with force constant k^b , while it diverges at distance b .

4.2.5 Harmonic angle potential

The bond angle vibration between a triplet of atoms $i - j - k$ is also represented by a harmonic potential on the angle θ_{ijk}

$$V_a(\theta_{ijk}) = \frac{1}{2}k_{ijk}^\theta (\theta_{ijk} - \theta_{ijk}^0)^2 \quad (4.48)$$

As the bond-angle vibration is represented by a harmonic potential, the form is the same as the bond stretching (Fig. 4.5).

The force equations are given by the chain rule:

$$\begin{aligned} \mathbf{F}_i &= -\frac{dV_a(\theta_{ijk})}{d\mathbf{r}_i} \\ \mathbf{F}_k &= -\frac{dV_a(\theta_{ijk})}{d\mathbf{r}_k} \\ \mathbf{F}_j &= -\mathbf{F}_i - \mathbf{F}_k \end{aligned} \quad \text{where } \theta_{ijk} = \arccos\left(\frac{\mathbf{r}_{ij} \cdot \mathbf{r}_{kj}}{r_{ij}r_{kj}}\right) \quad (4.49)$$

The numbering i, j, k is in sequence of covalently bonded atoms. Atom j is in the middle; atoms i and k are at the ends (see Fig. 4.7). **Note** that in the input in topology files, angles are given in degrees and force constants in kJ/mol/rad².

4.2.6 Cosine based angle potential

In the GROMOS-96 force field a simplified function is used to represent angle vibrations:

$$V_a(\theta_{ijk}) = \frac{1}{2} k_{ijk}^{\theta} \left(\cos(\theta_{ijk}) - \cos(\theta_{ijk}^0) \right)^2 \quad (4.50)$$

where

$$\cos(\theta_{ijk}) = \frac{\mathbf{r}_{ij} \cdot \mathbf{r}_{kj}}{r_{ij} r_{kj}} \quad (4.51)$$

The corresponding force can be derived by partial differentiation with respect to the atomic positions. The force constants in this function are related to the force constants in the harmonic form $k^{\theta, harm}$ (sec. 4.2.5) by:

$$k^{\theta} \sin^2(\theta_{ijk}^0) = k^{\theta, harm} \quad (4.52)$$

In the GROMOS-96 manual there is a much more complicated conversion formula which is temperature dependent. The formulas are equivalent at 0 K and the differences at 300 K are on the order of 0.1 to 0.2%. **Note** that in the input in topology files, angles are given in degrees and force constants in kJ/mol.

4.2.7 Urey-Bradley potential

The bond Urey-Bradley angle vibration between a triplet of atoms $i - j - k$ is represented by a harmonic potential on the angle θ_{ijk} and a harmonic correction term on the distance between the atoms i and k . Although this can be easily written as a simple sum of two terms, it is convenient to have it as a single entry in the topology file and in the output as a separate energy term. It is used mainly in the CHARMM force field [58]. The energy is given by:

$$V_a(\theta_{ijk}) = \frac{1}{2} k_{ijk}^{\theta} (\theta_{ijk} - \theta_{ijk}^0)^2 + \frac{1}{2} k_{ijk}^{UB} (r_{ik} - r_{ik}^0)^2 \quad (4.53)$$

The force equations can be deduced from sections 4.2.1 and 4.2.5.

4.2.8 Bond-Bond cross term

The bond-bond cross term for three particles i, j, k forming bonds $i - j$ and $k - j$ is given by [59]:

$$V_{rr'} = k_{rr'} (|\mathbf{r}_i - \mathbf{r}_j| - r_{1e}) (|\mathbf{r}_k - \mathbf{r}_j| - r_{2e}) \quad (4.54)$$

where $k_{rr'}$ is the force constant, and r_{1e} and r_{2e} are the equilibrium bond lengths of the $i - j$ and $k - j$ bonds respectively. The force associated with this potential on particle i is:

$$\mathbf{F}_i = -k_{rr'} (|\mathbf{r}_k - \mathbf{r}_j| - r_{2e}) \frac{\mathbf{r}_i - \mathbf{r}_j}{|\mathbf{r}_i - \mathbf{r}_j|} \quad (4.55)$$

the force on atom k can be obtained by swapping i and k in the above equation. Finally the force on atom j follows from the fact that the sum of internal forces should be zero: $\mathbf{F}_j = -\mathbf{F}_i - \mathbf{F}_k$.

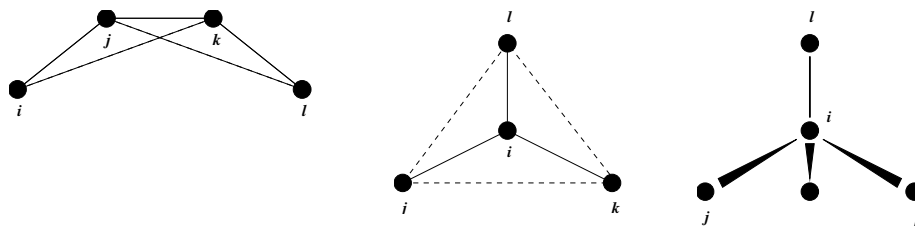


Figure 4.8: Principle of improper dihedral angles. Out of plane bending for rings (left), substituents of rings (middle), out of tetrahedral (right). The improper dihedral angle ξ is defined as the angle between planes (i,j,k) and (j,k,l) in all cases.

4.2.9 Bond-Angle cross term

The bond-angle cross term for three particles i, j, k forming bonds $i-j$ and $k-j$ is given by [59]:

$$V_{r\theta} = k_{r\theta} (|\mathbf{r}_i - \mathbf{r}_k| - r_{3e}) (|\mathbf{r}_i - \mathbf{r}_j| - r_{1e} + |\mathbf{r}_k - \mathbf{r}_j| - r_{2e}) \quad (4.56)$$

where $k_{r\theta}$ is the force constant, r_{3e} is the $i-k$ distance, and the other constants are the same as in Eqn. 4.54. The force associated with the potential on atom i is:

$$\mathbf{F}_i = -k_{r\theta} \left[(|\mathbf{r}_i - \mathbf{r}_k| - r_{3e}) \frac{\mathbf{r}_i - \mathbf{r}_j}{|\mathbf{r}_i - \mathbf{r}_j|} + (|\mathbf{r}_i - \mathbf{r}_j| - r_{1e} + |\mathbf{r}_k - \mathbf{r}_j| - r_{2e}) \frac{\mathbf{r}_i - \mathbf{r}_k}{|\mathbf{r}_i - \mathbf{r}_k|} \right] \quad (4.57)$$

4.2.10 Quartic angle potential

For special purposes there is an angle potential that uses a fourth order polynomial:

$$V_q(\theta_{ijk}) = \sum_{n=0}^5 C_n (\theta_{ijk} - \theta_{ijk}^0)^n \quad (4.58)$$

4.2.11 Improper dihedrals

Improper dihedrals are meant to keep planar groups planar (*e.g.* aromatic rings) or to prevent molecules from flipping over to their mirror images, see Fig. 4.8.

$$V_{id}(\xi_{ijkl}) = \frac{1}{2} k_\xi (\xi_{ijkl} - \xi_0)^2 \quad (4.59)$$

This is also a harmonic potential; it is plotted in Fig. 4.9. Note that, since it is harmonic, periodicity is not taken into account, so it is best to define improper dihedrals to have a ξ_0 as far away from $\pm 180^\circ$ as you can manage. **Note** that in the input in topology files, angles are given in degrees and force constants in kJ/mol/rad².

4.2.12 Proper dihedrals

For the normal dihedral interaction there is a choice of either the GROMOS periodic function or a function based on expansion in powers of $\cos \phi$ (the so-called Ryckaert-Bellemans potential). This

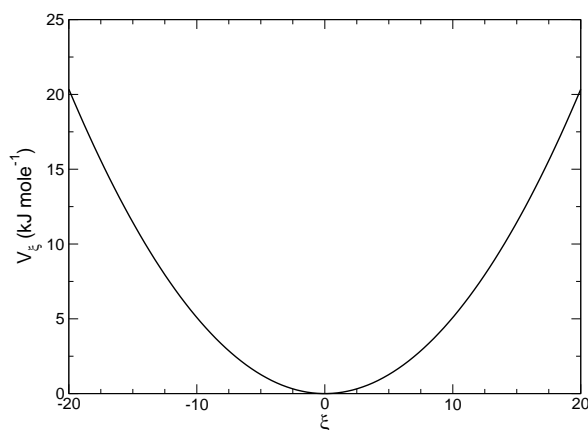
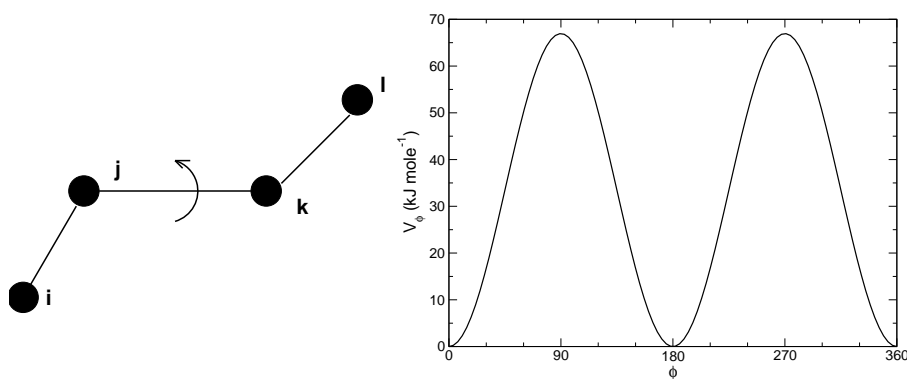


Figure 4.9: Improper dihedral potential.

Figure 4.10: Principle of proper dihedral angle (left, in *trans* form) and the dihedral angle potential (right).

choice has consequences for the inclusion of special interactions between the first and the fourth atom of the dihedral quadruple. With the periodic GROMOS potential a special 1-4 LJ-interaction must be included; with the Ryckaert-Bellemans potential *for alkanes* the 1-4 interactions must be excluded from the non-bonded list. **Note:** Ryckaert-Bellemans potentials are also used in e.g. the OPLS force field in combination with 1-4 interactions. You should therefore not modify topologies generated by `pdb2gmx` in this case.

Proper dihedrals: periodic type

Proper dihedral angles are defined according to the IUPAC/IUB convention, where ϕ is the angle between the ijk and the jkl planes, with **zero** corresponding to the *cis* configuration (i and l on the same side).

$$V_d(\phi_{ijkl}) = k_\phi(1 + \cos(n\phi - \phi_s)) \quad (4.60)$$

C_0	9.28	C_2	-13.12	C_4	26.24
C_1	12.16	C_3	-3.06	C_5	-31.5

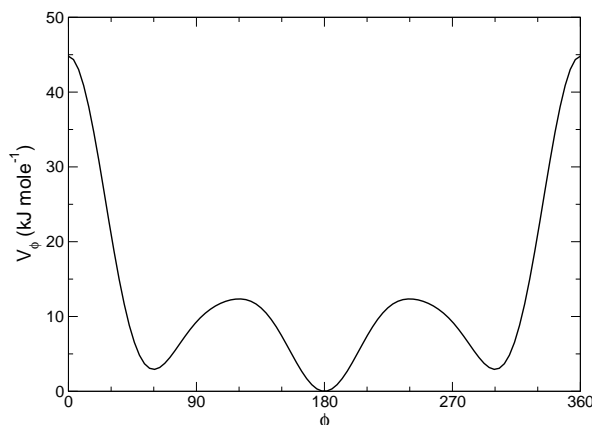
Table 4.1: Constants for Ryckaert-Bellemans potential (kJ mol^{-1}).

Figure 4.11: Ryckaert-Bellemans dihedral potential.

Proper dihedrals: Ryckaert-Bellemans function

For alkanes, the following proper dihedral potential is often used (see Fig. 4.11)

$$V_{rb}(\phi_{ijkl}) = \sum_{n=0}^5 C_n (\cos(\psi))^n, \quad (4.61)$$

where $\psi = \phi - 180^\circ$.

Note: A conversion from one convention to another can be achieved by multiplying every coefficient C_n by $(-1)^n$.

An example of constants for C is given in Table 4.1.

(Note: The use of this potential implies exclusion of LJ interactions between the first and the last atom of the dihedral, and ψ is defined according to the 'polymer convention' ($\psi_{trans} = 0$).)

The RB dihedral function can also be used to include the OPLS dihedral potential [60]. The OPLS potential function is given as the first four terms of a Fourier series:

$$V_{rb}(\phi_{ijkl}) = V_0 + \frac{1}{2}(V_1(1 + \cos(\psi)) + V_2(1 - \cos(2\psi)) + V_3(1 + \cos(3\psi))), \quad (4.62)$$

with $\psi = \phi$ (protein convention). Because of the equalities $\cos(2\phi) = 2(\cos(\phi))^2 - 1$ and $\cos(3\phi) = 4(\cos(\phi))^3 - 3\cos(\phi)$, one can translate the OPLS parameters to Ryckaert-Bellemans

parameters as follows:

$$\begin{aligned}
 C_0 &= V_0 + V_2 + \frac{1}{2}(V_1 + V_3) \\
 C_1 &= \frac{1}{2}(3V_3 - V_1) \\
 C_2 &= -V_2 \\
 C_3 &= -2V_3 \\
 C_4 &= 0 \\
 C_5 &= 0
 \end{aligned} \tag{4.63}$$

with OPLS parameters in protein convention and RB parameters in polymer convention.

Note: Mind the conversion from $kcal\ mol^{-1}$ for literature OPLS and RB parameters to $kJ\ mol^{-1}$ in GROMACS.

4.3 Restraints

Special potentials are used for imposing restraints on the motion of the system, either to avoid disastrous deviations, or to include knowledge from experimental data. In either case they are not really part of the force field and the reliability of the parameters is not important. The potential forms, as implemented in GROMACS, are mentioned just for the sake of completeness.

4.3.1 Position restraints

These are used to restrain particles to fixed reference positions \mathbf{R}_i . They can be used during equilibration in order to avoid too drastic rearrangements of critical parts (*e.g.* to restrain motion in a protein that is subjected to large solvent forces when the solvent is not yet equilibrated). Another application is the restraining of particles in a shell around a region that is simulated in detail, while the shell is only approximated because it lacks proper interaction from missing particles outside the shell. Restraining will then maintain the integrity of the inner part. For spherical shells it is a wise procedure to make the force constant depend on the radius, increasing from zero at the inner boundary to a large value at the outer boundary. This feature has not, however, been implemented in GROMACS.

The following form is used:

$$V_{pr}(\mathbf{r}_i) = \frac{1}{2}k_{pr}|\mathbf{r}_i - \mathbf{R}_i|^2 \tag{4.64}$$

The potential is plotted in Fig. 4.12.

The potential form can be rewritten without loss of generality as:

$$V_{pr}(\mathbf{r}_i) = \frac{1}{2} \left[k_{pr}^x (x_i - X_i)^2 \hat{\mathbf{x}} + k_{pr}^y (y_i - Y_i)^2 \hat{\mathbf{y}} + k_{pr}^z (z_i - Z_i)^2 \hat{\mathbf{z}} \right] \tag{4.65}$$

Now the forces are:

$$\begin{aligned}
 F_i^x &= -k_{pr}^x (x_i - X_i) \\
 F_i^y &= -k_{pr}^y (y_i - Y_i) \\
 F_i^z &= -k_{pr}^z (z_i - Z_i)
 \end{aligned} \tag{4.66}$$

Using three different force constants the position restraints can be turned on or off in each spatial dimension; this means that atoms can be harmonically restrained to a plane or a line. Position

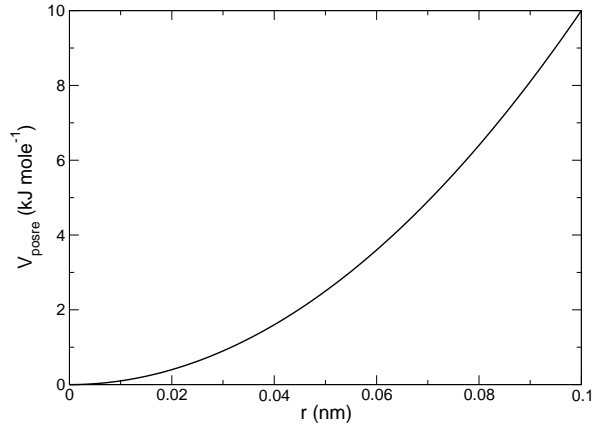


Figure 4.12: Position restraint potential.

restraints are applied to a special fixed list of atoms. Such a list is usually generated by the `pdb2gmx` program.

4.3.2 Angle restraints

These are used to restrain the angle between two pairs of particles or between one pair of particles and the Z-axis. The functional form is similar to that of a proper dihedral. For two pairs of atoms:

$$V_{ar}(\mathbf{r}_i, \mathbf{r}_j, \mathbf{r}_k, \mathbf{r}_l) = k_{ar}(1 - \cos(n(\theta - \theta_0))), \quad \text{where } \theta = \arccos\left(\frac{\mathbf{r}_j - \mathbf{r}_i}{\|\mathbf{r}_j - \mathbf{r}_i\|} \cdot \frac{\mathbf{r}_l - \mathbf{r}_k}{\|\mathbf{r}_l - \mathbf{r}_k\|}\right) \quad (4.67)$$

For one pair of atoms and the Z-axis:

$$V_{ar}(\mathbf{r}_i, \mathbf{r}_j) = k_{ar}(1 - \cos(n(\theta - \theta_0))), \quad \text{where } \theta = \arccos\left(\frac{\mathbf{r}_j - \mathbf{r}_i}{\|\mathbf{r}_j - \mathbf{r}_i\|} \cdot \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}\right) \quad (4.68)$$

A multiplicity (n) of 2 is useful when you do not want to distinguish between parallel and anti-parallel vectors.

4.3.3 Dihedral restraints

These are used to restrain the dihedral angle ϕ defined by four particles as in an improper dihedral (sec. 4.2.11) but with a slightly modified potential. Using

$$\phi' = (\phi - \phi_0) \text{ MOD } 2\pi \quad (4.69)$$

where ϕ_0 is the reference angle, the potential is defined as:

$$V_{dih}(\phi') = \begin{cases} \frac{1}{2}k_{dih}(\phi' - \phi_0 - \Delta\phi)^2 & \text{for } \phi' > \Delta\phi \\ 0 & \text{for } \phi' \leq \Delta\phi \end{cases} \quad (4.70)$$

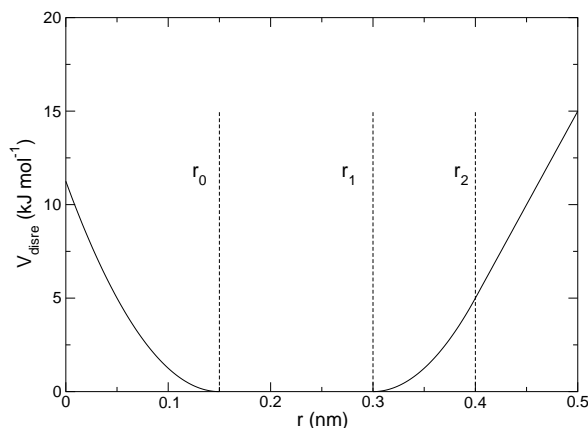


Figure 4.13: Distance Restraint potential.

where $\Delta\phi$ is a user defined angle and k_{dih} is the force constant. **Note** that in the input in topology files, angles are given in degrees and force constants in kJ/mol/rad².

4.3.4 Distance restraints

Distance restraints add a penalty to the potential when the distance between specified pairs of atoms exceeds a threshold value. They are normally used to impose experimental restraints, as from experiments in nuclear magnetic resonance (NMR), on the motion of the system. Thus MD can be used for structure refinement using NMR data. If one just wants to restrain the distance between two particles using a harmonic potential one should use [bonds] type 6 (see 5.3.5). The potential form for distance restraints is quadratic below a specified lower bound and between two specified upper bounds and linear beyond the largest bound (see Fig. 4.13).

$$V_{dr}(r_{ij}) = \begin{cases} \frac{1}{2}k_{dr}(r_{ij} - r_0)^2 & \text{for } r_{ij} < r_0 \\ 0 & \text{for } r_0 \leq r_{ij} < r_1 \\ \frac{1}{2}k_{dr}(r_{ij} - r_1)^2 & \text{for } r_1 \leq r_{ij} < r_2 \\ \frac{1}{2}k_{dr}(r_2 - r_1)(2r_{ij} - r_2 - r_1) & \text{for } r_2 \leq r_{ij} \end{cases} \quad (4.71)$$

The forces are

$$\mathbf{F}_i = \begin{cases} -k_{dr}(r_{ij} - r_0)\frac{\mathbf{r}_{ij}}{r_{ij}} & \text{for } r_{ij} < r_0 \\ 0 & \text{for } r_0 \leq r_{ij} < r_1 \\ -k_{dr}(r_{ij} - r_1)\frac{\mathbf{r}_{ij}}{r_{ij}} & \text{for } r_1 \leq r_{ij} < r_2 \\ -k_{dr}(r_2 - r_1)\frac{\mathbf{r}_{ij}}{r_{ij}} & \text{for } r_2 \leq r_{ij} \end{cases} \quad (4.72)$$

Time averaging

Distance restraints based on instantaneous distances can potentially reduce the fluctuations in a molecule significantly. This problem can be overcome by restraining to a *time averaged* distance [61]. The forces with time averaging are:

$$\mathbf{F}_i = \begin{cases} -k_{dr}^a (\bar{r}_{ij} - r_0) \frac{\mathbf{r}_{ij}}{r_{ij}} & \text{for } \bar{r}_{ij} < r_0 \\ 0 & \text{for } r_0 \leq \bar{r}_{ij} < r_1 \\ -k_{dr}^a (\bar{r}_{ij} - r_1) \frac{\mathbf{r}_{ij}}{r_{ij}} & \text{for } r_1 \leq \bar{r}_{ij} < r_2 \\ -k_{dr}^a (r_2 - r_1) \frac{\mathbf{r}_{ij}}{r_{ij}} & \text{for } r_2 \leq \bar{r}_{ij} \end{cases} \quad (4.73)$$

where \bar{r}_{ij} is given by an exponential running average with decay time τ :

$$\bar{r}_{ij} = \langle r_{ij}^{-3} \rangle^{-1/3} \quad (4.74)$$

and the force constant k_{dr}^a is switched on slowly to compensate for the lack of history at the beginning of the simulation:

$$k_{dr}^a = k_{dr} \left(1 - \exp\left(-\frac{t}{\tau}\right) \right) \quad (4.75)$$

Because of the time averaging we can no longer speak of a distance restraint potential.

This way an atom can satisfy two incompatible distance restraints *on average* by moving between two positions. An example would be an amino-acid side-chain which is rotating around its χ dihedral angle, thereby coming close to various other groups. Such a mobile side chain can give rise to multiple NOEs that can not be fulfilled by a single structure.

The computation of the time averaged distance in the `mdrun` program is done in the following fashion:

$$\begin{aligned} \overline{r_{ij}^{-3}}(0) &= r_{ij}(0)^{-3} \\ \overline{r_{ij}^{-3}}(t) &= \overline{r_{ij}^{-3}}(t - \Delta t) \exp\left(-\frac{\Delta t}{\tau}\right) + r_{ij}(t)^{-3} \left[1 - \exp\left(-\frac{\Delta t}{\tau}\right) \right] \end{aligned} \quad (4.76)$$

When a pair is within the bounds it can still feel a force, because the time averaged distance can still be beyond a bound. To prevent the protons from being pulled too close together a mixed approach can be used. In this approach the penalty is zero when the instantaneous distance is within the bounds, otherwise the violation is the square root of the product of the instantaneous violation and the time averaged violation:

$$\mathbf{F}_i = \begin{cases} k_{dr}^a \sqrt{(r_{ij} - r_0)(\bar{r}_{ij} - r_0)} \frac{\mathbf{r}_{ij}}{r_{ij}} & \text{for } r_{ij} < r_0 \quad \text{and} \quad \bar{r}_{ij} < r_0 \\ -k_{dr}^a \min\left(\sqrt{(r_{ij} - r_1)(\bar{r}_{ij} - r_1)}, r_2 - r_1\right) \frac{\mathbf{r}_{ij}}{r_{ij}} & \text{for } r_{ij} > r_1 \quad \text{and} \quad \bar{r}_{ij} > r_1 \\ 0 & \text{else} \end{cases} \quad (4.77)$$

Averaging over multiple pairs

Sometimes it is unclear from experimental data which atom pair gives rise to a single NOE, in other occasions it can be obvious that more than one pair contributes due to the symmetry of the system, *e.g.* a methyl group with three protons. For such a group it is not possible to distinguish between the protons, therefore they should all be taken into account when calculating the distance between this methyl group and another proton (or group of protons). Due to the physical nature of magnetic resonance, the intensity of the NOE signal is inversely proportional to the sixth power of the interatomic distance. Thus, when combining atom pairs, a fixed list of N restraints may be taken together, where the apparent “distance” is given by:

$$r_N(t) = \left[\sum_{n=1}^N \bar{r}_n(t)^{-6} \right]^{-1/6} \quad (4.78)$$

where we use r_{ij} or eqn. 4.74 for the \bar{r}_n . The r_N of the instantaneous and time-averaged distances can be combined to do a mixed restraining as indicated above. As more pairs of protons contribute to the same NOE signal, the intensity will increase, and the summed “distance” will be shorter than any of its components due to the reciprocal summation.

There are two options for distributing the forces over the atom pairs. In the conservative option the force is defined as the derivative of the restraint potential with respect to the coordinates. This results in a conservative potential when time averaging is not used. The force distribution over the pairs is proportional to r^{-6} . This means that a close pair feels a much larger force than a distant pair, which might lead to a ‘too rigid’ molecule. The other option is an equal force distribution. In this case each pair feels $1/N$ of the derivative of the restraint potential with respect to r_N . The advantage of this method is that more conformations might be sampled, but the non-conservative nature of the forces can lead to local heating of the protons.

It is also possible to use *ensemble averaging* using multiple (protein) molecules. In this case the bounds should be lowered as in:

$$\begin{aligned} r_1 &= r_1 * M^{-1/6} \\ r_2 &= r_2 * M^{-1/6} \end{aligned} \quad (4.79)$$

where M is the number of molecules. The GROMACS preprocessor `grompp` can do this automatically when the appropriate option is given. The resulting “distance” is then used to calculate the scalar force according to:

$$\begin{aligned} \mathbf{F}_i &= 0 & r_N < r_1 \\ &= -k_{dr}(r_N - r_1) \frac{\mathbf{r}_{ij}}{r_{ij}} & r_1 \leq r_N < r_2 \\ &= -k_{dr}(r_2 - r_1) \frac{\mathbf{r}_{ij}}{r_{ij}} & r_N \geq r_2 \end{aligned} \quad (4.80)$$

where i and j denote the atoms of all the pairs that contribute to the NOE signal.

Using distance restraints

A list of distance restraints based on NOE data can be added to a molecule definition in your topology file, like in the following example:

```
[ distance_restraints ]
; ai  aj  type  index  type'  low  up1  up2  fac
  10  16   1    0     1    0.0  0.3  0.4  1.0
  10  28   1    1     1    0.0  0.3  0.4  1.0
  10  46   1    1     1    0.0  0.3  0.4  1.0
  16  22   1    2     1    0.0  0.3  0.4  2.5
  16  34   1    3     1    0.0  0.5  0.6  1.0
```

In this example a number of features can be found. In columns `ai` and `aj` you find the atom numbers of the particles to be restrained. The `type` column should always be 1. As explained in sec. 4.3.4, multiple distances can contribute to a single NOE signal. In the topology this can be set using the `index` column. In our example, the restraints 10-28 and 10-46 both have index 1, therefore they are treated simultaneously. An extra requirement for treating restraints together, is that the restraints should be on successive lines, without any other intervening restraint. The `type'` column will usually be 1, but can be set to 2 to obtain a distance restraint which will never be time and ensemble averaged; this can be useful for restraining hydrogen bonds. The columns `low`, `up1` and `up2` hold the values of r_0 , r_1 and r_2 from eqn. 4.71. In some cases it can be useful to have different force constants for some restraints; this is controlled by the column `fac`. The force constant in the parameter file is multiplied by the value in the column `fac` for each restraint.

Some parameters for NMR refinement can be specified in the `grompp.mdp` file:

disre: type of distance restraining. The `disre` variable sets the type of distance restraint. `no/simple` turns the distance restraints off/on. When multiple proteins or peptides are present in one simulation box, ensemble averaging can be turned on by setting `disre = ensemble`. Normally one would perform ensemble averaging over multiple subsystems, each in a separate box, using `mdrun -multi; supply topol0.tpr, topol1.tpr, ...` with different coordinates and/or velocities.

disre_weighting: force-weighting in restraints with multiple pairs. By default, the force due to the distance restraint is distributed equally over all the pairs involved in the restraint. This can also be explicitly selected with `disre_weighting = equal`. If you instead set this option to `disre_weighting = conservative` you get conservative forces when `disre_tau = 0`.

disre_mixed: how to calculate the violations. `disre_mixed = no` gives normal time-averaged violations. When `disre_mixed = yes` the square root of the product of the time-averaged and the instantaneous violations is used.

disre_fc: force constant k_{dr} for distance restraints. k_{dr} (eqn. 4.71) can be set as variable `disre_fc = 1000` for a force constant of 1000 kJ mol⁻¹ nm⁻². This value is multiplied by the value in the `fac` column in the distance restraint entries in the topology file.

disre_tau: time constant for restraints. τ (eqn. 4.76) can be set as variable `disre_tau = 10` for a time constant of 10 ps. Time averaging can be turned off by setting `disre_tau` to 0.

nstdisreout: pair distance output frequency. Determines how often the time-averaged and instantaneous distances of all atom pairs involved in distance restraints are written to the energy file.

4.3.5 Orientation restraints

This section describes how orientations between vectors, as measured in certain NMR experiments, can be calculated and restrained in MD simulations. The presented refinement methodology and a comparison of results with and without time and ensemble averaging have been published [62].

Theory

In an NMR experiment orientations of vectors can be measured when a molecule does not tumble completely isotropically in the solvent. Two examples of such orientation measurements are residual dipolar couplings (between two nuclei) or chemical shift anisotropies. An observable for a vector \mathbf{r}_i can be written as follows:

$$\delta_i = \frac{2}{3} \text{tr}(\mathbf{S}\mathbf{D}_i) \quad (4.81)$$

where \mathbf{S} is the dimensionless order tensor of the molecule. The tensor \mathbf{D}_i is given by:

$$\mathbf{D}_i = \frac{c_i}{\|\mathbf{r}_i\|^\alpha} \begin{pmatrix} 3xx - 1 & 3xy & 3xz \\ 3xy & 3yy - 1 & 3yz \\ 3xz & 3yz & 3zz - 1 \end{pmatrix} \quad (4.82)$$

$$\text{with: } x = \frac{r_{i,x}}{\|\mathbf{r}_i\|}, \quad y = \frac{r_{i,y}}{\|\mathbf{r}_i\|}, \quad z = \frac{r_{i,z}}{\|\mathbf{r}_i\|} \quad (4.83)$$

For a dipolar coupling \mathbf{r}_i is the vector connecting the two nuclei, $\alpha = 3$ and the constant c_i is given by:

$$c_i = \frac{\mu_0}{4\pi} \gamma_1^i \gamma_2^i \frac{\hbar}{4\pi} \quad (4.84)$$

where γ_1^i and γ_2^i are the gyromagnetic ratios of the two nuclei.

The order tensor is symmetric and has trace zero. Using a rotation matrix \mathbf{T} it can be transformed into the following form:

$$\mathbf{T}^T \mathbf{S} \mathbf{T} = s \begin{pmatrix} -\frac{1}{2}(1 - \eta) & 0 & 0 \\ 0 & -\frac{1}{2}(1 + \eta) & 0 \\ 0 & 0 & 1 \end{pmatrix} \quad (4.85)$$

where $-1 \leq s \leq 1$ and $0 \leq \eta \leq 1$. s is called the order parameter and η the asymmetry of the order tensor \mathbf{S} . When the molecule tumbles isotropically in the solvent s is zero and no orientational effects can be observed as all δ_i are zero.

Calculating orientations in a simulation

For reasons which are explained below, the \mathbf{D} matrices are calculated which respect to a reference orientation of the molecule. The orientation is defined by a rotation matrix \mathbf{R} which is needed to least-squares fit the current coordinates of a selected set of atoms onto a reference conformation.

The reference conformation is the starting conformation of the simulation. In case of ensemble averaging, which will be treated later, the structure is taken from the first subsystem. The calculated \mathbf{D}_i^c matrix is given by:

$$\mathbf{D}_i^c(t) = \mathbf{R}(t)\mathbf{D}_i(t)\mathbf{R}^T(t) \quad (4.86)$$

The calculated orientation for vector i is given by:

$$\delta_i^c(t) = \frac{2}{3}\text{tr}(\mathbf{S}(t)\mathbf{D}_i^c(t)) \quad (4.87)$$

The order tensor $\mathbf{S}(t)$ is usually unknown. A reasonable choice for the order tensor is the tensor which minimizes the (weighted) mean square difference between the calculated and the observed orientations:

$$MSD(t) = \left(\sum_{i=1}^N w_i \right)^{-1} \sum_{i=1}^N w_i (\delta_i^c(t) - \delta_i^{exp})^2 \quad (4.88)$$

To properly combine different types of measurements the unit of w_i should be such that all terms are dimensionless. This means the unit of w_i is the unit of δ_i to the power -2 . Note that scaling all w_i with a constant factor does not influence the order tensor.

Time averaging

Since the tensors \mathbf{D}_i fluctuate rapidly in time, much faster than can be observed in experiment, they should be time averaged in the simulation. However, in a simulation the time as well as the number of copies of a molecule is limited. Usually one can not obtain a converged average of the \mathbf{D}_i tensors over all orientations of the molecule. If one assumes that the average orientations of the r_i vectors within the molecule converge much faster than the tumbling time of the molecule, the tensor can be averaged in an axis system which rotates with the molecule, as expressed by equation (4.86). The time averaged tensors are calculated using an exponentially decaying memory function:

$$\mathbf{D}_i^a(t) = \frac{\int_{u=t_0}^t \mathbf{D}_i^c(u) \exp\left(-\frac{t-u}{\tau}\right) du}{\int_{u=t_0}^t \exp\left(-\frac{t-u}{\tau}\right) du} \quad (4.89)$$

Assuming that the order tensor \mathbf{S} fluctuates slower than the \mathbf{D}_i , the time averaged orientation can be calculated as:

$$\delta_i^a(t) = \frac{2}{3}\text{tr}(\mathbf{S}(t)\mathbf{D}_i^a(t)) \quad (4.90)$$

where the order tensor $\mathbf{S}(t)$ is calculated using expression (4.88) with $\delta_i^c(t)$ replaced by $\delta_i^a(t)$.

Restraining

The simulated structure can be restrained by applying a force proportional to the difference between the calculated and the experimental orientations. When no time averaging is applied a proper potential can be defined as:

$$V = \frac{1}{2}k \sum_{i=1}^N w_i (\delta_i^c(t) - \delta_i^{exp})^2 \quad (4.91)$$

where the unit of k is the unit of energy. Thus the effective force constant for restraint i is kw_i . The forces are given by minus the gradient of V . The force \mathbf{F}_i working on vector \mathbf{r}_i is:

$$\begin{aligned}\mathbf{F}_i(t) &= -\frac{dV}{d\mathbf{r}_i} \\ &= -kw_i(\delta_i^c(t) - \delta_i^{exp})\frac{d\delta_i(t)}{d\mathbf{r}_i} \\ &= -kw_i(\delta_i^c(t) - \delta_i^{exp})\frac{2c_i}{\|\mathbf{r}\|^{2+\alpha}}\left(2\mathbf{R}^T\mathbf{S}\mathbf{R}\mathbf{r}_i - \frac{2+\alpha}{\|\mathbf{r}\|^2}\text{tr}(\mathbf{R}^T\mathbf{S}\mathbf{R}\mathbf{r}_i\mathbf{r}_i^T)\mathbf{r}_i\right)\end{aligned}$$

Ensemble averaging

Ensemble averaging can be applied by simulating a system of M subsystems which each contain an identical set of orientation restraints. The systems only interact via the orientation restraint potential which is defined as:

$$V = M\frac{1}{2}k\sum_{i=1}^N w_i\langle\delta_i^c(t) - \delta_i^{exp}\rangle^2 \quad (4.92)$$

The force on vector $\mathbf{r}_{i,m}$ in subsystem m is given by:

$$\mathbf{F}_{i,m}(t) = -\frac{dV}{d\mathbf{r}_{i,m}} = -kw_i\langle\delta_i^c(t) - \delta_i^{exp}\rangle\frac{d\delta_{i,m}^c(t)}{d\mathbf{r}_{i,m}} \quad (4.93)$$

Time averaging

When using time averaging it is not possible to define a potential. We can still define a quantity which gives a rough idea of the energy stored in the restraints:

$$V = M\frac{1}{2}k^a\sum_{i=1}^N w_i\langle\delta_i^a(t) - \delta_i^{exp}\rangle^2 \quad (4.94)$$

The force constant k_a is switched on slowly to compensate for the lack of history at times close to t_0 . It is exactly proportional to the amount of average which has been accumulated:

$$k^a = k\frac{1}{\tau}\int_{u=t_0}^t \exp\left(-\frac{t-u}{\tau}\right) du \quad (4.95)$$

What really matters is the definition of the force. It is chosen to be proportional to the square root of the product of the time averaged and the instantaneous deviation. Using only the time averaged deviation induces large oscillations. The force is given by:

$$\mathbf{F}_{i,m}(t) = \begin{cases} 0 & \text{for } ab \leq 0 \\ k^a w_i \frac{a}{|a|} \sqrt{ab} \frac{d\delta_{i,m}^c(t)}{d\mathbf{r}_{i,m}} & \text{for } ab > 0 \end{cases} \quad (4.96)$$

$$a = \langle\delta_i^a(t) - \delta_i^{exp}\rangle$$

$$b = \langle\delta_i^c(t) - \delta_i^{exp}\rangle$$

Using orientation restraints

Orientation restraints can be added to a molecule definition in the topology in the section [`orientation_restraints`]. Here we give an example section containing five N-H residual dipolar coupling restraints:

```
[ orientation_restraints ]
; ai  aj  type  exp.  label  alpha  const.  obs.  weight
;                               Hz nm3  Hz  1/Hz2
31  32   1    1    3    3    6.083  -6.73  1.0
43  44   1    1    4    3    6.083  -7.87  1.0
55  56   1    1    5    3    6.083  -7.13  1.0
65  66   1    1    6    3    6.083  -2.57  1.0
73  74   1    1    7    3    6.083  -2.10  1.0
```

The unit of the observable is Hz, but one can choose any other unit. In columns `ai` and `aj` you find the atom numbers of the particles to be restrained. The `type` column should always be 1. The `exp.` column denotes the experiment number, this starts numbering at 1. For each experiment a separate order tensor \mathbf{S} is optimized. The `label` should be a unique number larger than zero for each restraint. The `alpha` column contains the power α which is used in equation (4.82) to calculate the orientation. The `const.` column contains the constant c_i used in the same equation. The constant should have the unit of the observable times nm^α . The column `obs.` contains the observable, in any unit you like. The last column contains the weights w_i , the unit should be the inverse of the square of the unit of the observable.

Some parameters for orientation restraints can be specified in the `grompp.mdp` file, for a study of the effect of different force constants and averaging times and ensemble averaging see [62].

orire: use orientation restraining. `no/yes` turns the distance restraints off/on. Ensemble averaging can be performed using `mdrun -multi`, which simulates multiple subsystems in separate boxes; supply `topol0.tpr`, `topol1.tpr`, ... with different coordinates and/or velocities.

orire_fc: force constant k for orientation restraints. The unit of k is kJ mol^{-1} . Note that the force constant for a restraint is this force constant times the weight of the restraint. When set to zero one obtain the calculated orientation without affecting the simulation.

orire_tau: time constant τ for restraints. Set `orire_tau = 10` for a time constant of 10 ps. Time averaging can be turned off by setting `orire_tau` to 0.

orire_fitgrp: the fit group for the restraints. This group of atoms is used to determine the rotation \mathbf{R} of the system with respect to the reference orientation. The reference orientation is the starting conformation of the first subsystem. For a protein backbone should be a reasonable choice.

nstorireout: orientation output frequency. Determines how often the orientations for all restraints and the order tensor(s) \mathbf{S} are written to the energy file. When using time and/or ensemble averaging, the time and ensemble averaged orientations as well as the instantaneous non-ensemble averaged orientations are written to the energy file. These can be analyzed using `g_energy`.

4.4 Polarization

Polarization can be treated by GROMACS by attaching shell (drude) particles to atoms and/or virtual sites. The energy of the shell particle is then minimized at each time step in order to remain on the Born-Oppenheimer surface.

4.4.1 Simple polarization

This is merely a harmonic potential with equilibrium distance 0.

4.4.2 Water polarization

A special potential for water that allows anisotropic polarization of a single shell particle [29].

4.4.3 Thole polarization

Based on early work by Thole [63] Roux and coworkers have implemented potentials for molecules like ethanol [64, 65, 66]. Within such molecules there are intramolecular interaction between shell particles, however these must be screened because full Coulomb would be too strong. The potential between two shell particles i and j is:

$$V_{thole} = \frac{q_i q_j}{r_{ij}} \left[1 - \left(1 - \frac{\bar{r}_{ij}}{2} \right) \exp^{-\bar{r}_{ij}} \right] \quad (4.97)$$

where

$$\bar{r}_{ij} = a \frac{r_{ij}}{(\alpha_i \alpha_j)^{1/6}} \quad (4.98)$$

where a is a magic (dimensionless) constant, usually chosen to be 2.6 [66] and α_i, α_j are the polarizabilities of the respective shell particles.

4.5 Free energy interactions

This section describes the λ -dependence of the potentials used for free energy calculations (see sec. 3.12). All common types of potentials and constraints can be interpolated smoothly from state A ($\lambda = 0$) to state B ($\lambda = 1$) and vice versa. All bonded interactions are interpolated by linear interpolation of the interaction parameters. Non-bonded interactions can be interpolated linearly or via soft-core interactions.

Harmonic potentials

The example given here is for the bond potential, which is harmonic in GROMACS. However, these equations apply to the angle potential and the improper dihedral potential as well.

$$V_b = \frac{1}{2} ((1 - \lambda)k_b^A + \lambda k_b^B) (b - (1 - \lambda)b_0^A - \lambda b_0^B)^2 \quad (4.99)$$

$$\frac{\partial V_b}{\partial \lambda} = \frac{1}{2}(k_b^B - k_b^A) \left[b - (1 - \lambda)b_0^A + \lambda b_0^B \right]^2 + (b_0^A - b_0^B)(b - (1 - \lambda)b_0^A - \lambda b_0^B) \quad (4.100)$$

GROMOS-96 bonds and angles

Fourth power bond stretching and cosine based angle potentials are interpolated by linear interpolation of the force constant and the equilibrium position. Formulas are not given here.

Proper dihedrals

For the proper dihedrals, the equations are somewhat more complicated:

$$V_d = ((1 - \lambda)k_d^A + \lambda k_d^B)(1 + \cos(n_\phi \phi - ((1 - \lambda)\phi_s^A + \lambda\phi_s^B))) \quad (4.101)$$

$$\begin{aligned} \frac{\partial V_d}{\partial \lambda} = & (k_d^B - k_d^A) \left[1 + \cos(n_\phi \phi - [(1 - \lambda)\phi_s^A + \lambda\phi_s^B]) - \right. \\ & \left. ((1 - \lambda)k_d^A + \lambda k_d^B)(\phi_s^A - \phi_s^B) \sin(n_\phi \phi - [(1 - \lambda)\phi_s^A + \lambda\phi_s^B]) \right] \quad (4.102) \end{aligned}$$

Note: that the multiplicity n_ϕ can not be parameterized because the function should remain periodic on the interval $[0, 2\pi]$.

Coulomb interaction

The Coulomb interaction between two particles of which the charge varies with λ is:

$$V_c = \frac{f}{\epsilon_{rf} r_{ij}} \left[((1 - \lambda)q_i^A + \lambda q_i^B) \cdot ((1 - \lambda)q_j^A + \lambda q_j^B) \right] \quad (4.103)$$

$$\frac{\partial V_c}{\partial \lambda} = \frac{f}{\epsilon_{rf} r_{ij}} \left[(q_j^B - q_j^A)((1 - \lambda)q_i^A + \lambda q_i^B) + (q_i^B - q_i^A)((1 - \lambda)q_j^A + \lambda q_j^B) \right] \quad (4.104)$$

where $f = \frac{1}{4\pi\epsilon_0} = 138.935485$ (see chapter 2)

Coulomb interaction with Reaction Field

The coulomb interaction including a reaction field, between two particles of which the charge varies with λ is:

$$V_c = f \left[\frac{1}{r_{ij}} + k_{rf} r_{ij}^2 - c_{rf} \right] \left[((1 - \lambda)q_i^A + \lambda q_i^B) \cdot ((1 - \lambda)q_j^A + \lambda q_j^B) \right] \quad (4.105)$$

$$\begin{aligned} \frac{\partial V_c}{\partial \lambda} = & f \left[\frac{1}{r_{ij}} + k_{rf} r_{ij}^2 - c_{rf} \right] \cdot \\ & \left[(q_j^B - q_j^A)((1 - \lambda)q_i^A + \lambda q_i^B) + (q_i^B - q_i^A)((1 - \lambda)q_j^A + \lambda q_j^B) \right] \quad (4.106) \end{aligned}$$

Note that the constants k_{rf} and c_{rf} are defined using the dielectric constant ϵ_{rf} of the medium (see sec. 4.1.4).

Lennard-Jones interaction

For the Lennard-Jones interaction between two particles of which the *atom type* varies with λ we can write:

$$V_{LJ} = \frac{((1 - \lambda)C_{12}^A + \lambda C_{12}^B)}{r_{ij}^{12}} - \frac{(1 - \lambda)C_6^A + \lambda C_6^B}{r_{ij}^6} \quad (4.107)$$

$$\frac{\partial V_{LJ}}{\partial \lambda} = \frac{C_{12}^B - C_{12}^A}{r_{ij}^{12}} - \frac{C_6^B - C_6^A}{r_{ij}^6} \quad (4.108)$$

It should be noted that it is also possible to express a pathway from state A to state B using σ and ϵ (see eqn. 4.5). It may seem to make sense physically, to vary the forcefield parameters σ and ϵ rather than the derived parameters C_{12} and C_6 . However, the difference between the pathways in parameter space is not large, and the free energy itself does not depend on the pathway, so we use the simple formulation presented above.

Kinetic Energy

When the mass of a particle changes, there is also a contribution of the kinetic energy to the free energy (note that we can not write the momentum \mathbf{p} as $m\mathbf{v}$, since that would result in the sign of $\frac{\partial Ek}{\partial \lambda}$ being incorrect [67]):

$$Ek = \frac{1}{2} \frac{\mathbf{p}^2}{(1 - \lambda)m^A + \lambda m^B} \quad (4.109)$$

$$\frac{\partial Ek}{\partial \lambda} = -\frac{1}{2} \frac{\mathbf{p}^2(m^B - m^A)}{((1 - \lambda)m^A + \lambda m^B)^2} \quad (4.110)$$

after taking the derivative, we *can* insert $\mathbf{p} = m\mathbf{v}$, such that:

$$\frac{\partial Ek}{\partial \lambda} = -\frac{1}{2} \mathbf{v}^2(m^B - m^A) \quad (4.111)$$

Constraints

The constraints are formally part of the Hamiltonian, and therefore they give a contribution to the free energy. In GROMACS this can be calculated using the LINCS or the SHAKE algorithm. If we have a number of constraint equations g_k :

$$g_k = r_k - d_k \quad (4.112)$$

where r_k is the distance vector between two particles and d_k is the constraint distance between the two particles, we can write this using a λ -dependent distance as

$$g_k = r_k - ((1 - \lambda)d_k^A + \lambda d_k^B) \quad (4.113)$$

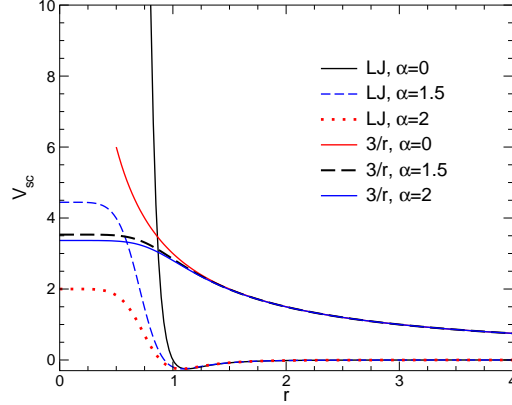


Figure 4.14: Soft-core interactions at $\lambda = 0.5$, with $p = 2$ and $C_6^A = C_{12}^A = C_6^B = C_{12}^B = 1$.

the contribution C_λ to the Hamiltonian using Lagrange multipliers λ :

$$C_\lambda = \sum_k \lambda_k g_k \quad (4.114)$$

$$\frac{\partial C_\lambda}{\partial \lambda} = \sum_k \lambda_k (d_k^B - d_k^A) \quad (4.115)$$

4.5.1 Soft-core interactions

The linear interpolation of the Lennard-Jones and Coulomb potentials gives problems when growing particles out of nothing or when making particles disappear (λ close to 0 or 1). To circumvent these problems, the singularities in the potentials need to be removed. This is done with soft-core potentials. In GROMACS the soft-core potential V_{sc} is:

$$V_{sc}(r) = (1 - \lambda)V^A(r_A) + \lambda V^B(r_B) \quad (4.116)$$

$$r_A = \left(\alpha \sigma_A^6 \lambda^p + r^6 \right)^{\frac{1}{6}} \quad (4.117)$$

$$r_B = \left(\alpha \sigma_B^6 (1 - \lambda)^p + r^6 \right)^{\frac{1}{6}} \quad (4.118)$$

where V^A and V^B are the normal “hard core” Van der Waals or electrostatic potentials in state A ($\lambda = 0$) and state B ($\lambda = 1$) respectively, α is the soft-core parameter, p is the soft-core λ power, σ is the radius of the interaction, which is $(C_{12}/C_6)^{1/6}$ or a predefined value when C_6 or C_{12} is zero. For intermediate λ , r_A and r_B alter the interactions very little when $r > \alpha^{1/6}\sigma$ and they quickly switch the soft-core interaction to an almost constant value when r becomes smaller (Fig. 4.14). The force is:

$$F_{sc}(r) = -\frac{\partial V_{sc}(r)}{\partial r} = (1 - \lambda)F^A(r_A) \left(\frac{r}{r_A} \right)^5 + \lambda F^B(r_B) \left(\frac{r}{r_B} \right)^5 \quad (4.119)$$

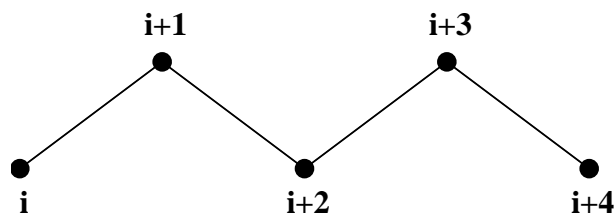


Figure 4.15: Atoms along an alkane chain.

where F^A and F^B are the 'hard core' forces. The contribution to the derivative of the free energy is:

$$\frac{\partial V_{sc}(r)}{\partial \lambda} = -V^A(r_A) + V^B(r_B) + \alpha \frac{p}{6} \left(-(1-\lambda)\lambda^{p-1} F^A(r_A) \sigma_A^6 r_A^{-5} + \lambda(1-\lambda)^{p-1} F^B(r_B) \sigma_B^6 r_B^{-5} \right) \quad (4.120)$$

The original Gromos Lennard-Jones soft-core function [68] uses $p=2$, but $p=1$ gives a smoother $\partial H/\partial \lambda$ curve. When the changes between the two states involve both the disappearing and appearing of atoms, it is important that the overlapping of atoms happens around $\lambda=0.5$. This can usually be achieved with $\alpha \approx 0.7$ for $p=1$ and $\alpha \approx 1.5$ for $p=2$. Another issue which should be considered is the soft-core effect of hydrogens without Lennard-Jones interaction. Their soft-core σ is set with `sc-sigma` in the `.mdp` file. These hydrogens produce peaks in $\partial H/\partial \lambda$ at λ is 0 and/or 1 for $p=1$ and close to 0 and/or 1 with $p=2$. Lowering `sc-sigma` will decrease this effect, but it will also increase the interactions with hydrogens relative to the other interactions in the soft-core state.

4.6 Methods

4.6.1 Exclusions and 1-4 Interactions.

Atoms within a molecule that are close by in the chain, *i.e.* atoms that are covalently bonded, or linked by one respectively two atoms are so-called *first neighbors*, *second neighbors* and *third neighbors*, (see Fig. 4.15). Since the interactions of atom **i** with atoms **i+1** and **i+2**

are mainly quantum mechanical, they can not be modeled by a Lennard-Jones potential. Instead it is assumed that these interactions are adequately modeled by a harmonic bond term or constraint (**i**, **i+1**) and a harmonic angle term (**i**, **i+2**). The first and second neighbors (atoms **i+1** and **i+2**) are therefore *excluded* from the Lennard-Jones interaction list of atom **i**; atoms **i+1** and **i+2** are called *exclusions* of atom **i**.

For third neighbors the normal Lennard-Jones repulsion is sometimes still too strong, which means that when applied to a molecule the molecule would deform or break due to the internal strain. This is especially the case for carbon-carbon interactions in a *cis*-conformation (*e.g.* *cis*-butane). Therefore for some of these interactions the Lennard-Jones repulsion has been reduced in the GROMOS force field, which is implemented by keeping a separate list of 1-4 and normal Lennard-Jones parameters. In other force fields, such as OPLS [60], the standard Lennard-Jones parameters

are reduced by a factor of two, but in that case also the dispersion (r^{-6}) and the coulomb interaction are scaled. GROMACS can use either of these methods.

4.6.2 Charge Groups.

In principle the force calculation in MD is an $O(N^2)$ problem. Therefore we apply a cutoff for non-bonded force (NBF) calculations: only the particles within a certain distance of each other are interacting. This reduces the cost to $O(N)$ (typically $100N$ to $200N$) of the NBF. It also introduces an error, which is, in most cases, acceptable, except when applying the cutoff implies the creation of charges, in which case you should consider using the lattice sum methods provided by GROMACS.

Consider a water molecule interacting with another atom. When we would apply the cutoff on an atom-atom basis we might include the atom-Oxygen interaction (with a charge of -0.82) without the compensating charge of the protons and so induce a large dipole moment over the system. Therefore we have to keep groups of atoms with total charge 0 together. These groups are called *charge groups*.

4.6.3 Treatment of Cutoffs

GROMACS is quite flexible in treating cutoffs, which implies there can be quite a number of parameters to set. These parameters are set in the input file for `grompp`. There are two sort of parameters that affect the cutoff interactions; you can select which type of interaction to use in each case, and which cutoffs should be used in the neighborsearching.

For both Coulomb and van der Waals interactions there are interaction type selectors (termed `vdwtype` and `coulombtype`) and two parameters, for a total of six nonbonded interaction parameters. See sec. 7.3.1 for a complete description of these parameters.

The neighbor searching (NS) can be performed using a single-range, or a twin-range approach. Since the former is merely a special case of the latter we will discuss the more general twin-range. In this case NS is described by two radii `rlist` and `max(rcoulomb,rvdw)`. Usually one builds the neighbor list every 10 time steps or every 20 fs (parameter `nstlist`). In the neighbor list all interaction pairs that fall within `rlist` are stored. Furthermore, the interactions between pairs that do not fall within `rlist` but do fall within `max(rcoulomb,rvdw)` are computed during NS, and the forces and energy are stored separately, and added to short-range forces at every time step between successive NS. If `rlist = max(rcoulomb,rvdw)`, no forces are evaluated during neighbor list generation. The virial is calculated from the sum of the short- and long-range forces. This means that the virial can be slightly asymmetrical at non-NS steps. In single precision the virial is almost always asymmetrical, because the off-diagonal elements are about as large as each element in the sum. In most cases this is not really a problem, since the fluctuations in de virial can be 2 orders of magnitude larger than the average.

Except for the plain cutoff, all of the interaction functions in Table 4.2 require that neighbor searching is done with a larger radius than the r_c specified for the functional form, because of the use of charge groups. The extra radius is typically of the order of 0.25 nm (roughly the largest distance between two atoms in a charge group plus the distance a charge group can diffuse within neighbor list updates).

	Type	Parameters
Coulomb	Plain cutoff	r_c, ϵ_r
	Reaction field	r_c, ϵ_{rf}
	Shift function	r_1, r_c, ϵ_r
	Switch function	r_1, r_c, ϵ_r
VdW	Plain cutoff	r_c
	Shift function	r_1, r_c
	Switch function	r_1, r_c

Table 4.2: Parameters for the different functional forms of the non-bonded interactions.

4.7 Virtual interaction-sites

Virtual interaction-sites (called dummy atoms in GROMACS versions before 3.3) can be used in GROMACS in a number of ways. We write the position of the virtual site \mathbf{r}_v as a function of the positions of other particles \mathbf{r}_i : $\mathbf{r}_d = f(\mathbf{r}_1.. \mathbf{r}_n)$. The virtual site, which may carry charge, or can be involved in other interactions can now be used in the force calculation. The force acting on the virtual site must be redistributed over the particles with mass in a consistent way. A good way to do this can be found in ref. [69]. We can write the potential energy as

$$V = V(\mathbf{r}_v, \mathbf{r}_1.. \mathbf{r}_n) = V^*(\mathbf{r}_1.. \mathbf{r}_n) \quad (4.121)$$

The force on the particle i is then

$$\mathbf{F}_i = -\frac{\partial V^*}{\partial \mathbf{r}_i} = -\frac{\partial V}{\partial \mathbf{r}_i} - \frac{\partial \mathbf{r}_v}{\partial \mathbf{r}_i} \frac{\partial V}{\partial \mathbf{r}_v} = \mathbf{F}_i^{direct} + \mathbf{F}'_i \quad (4.122)$$

the first term of which is the normal force. The second term is the force on particle i due to the virtual site, which can be written in tensor notation:

$$\mathbf{F}'_i = \begin{bmatrix} \frac{\partial x_v}{\partial x_i} & \frac{\partial y_v}{\partial x_i} & \frac{\partial z_v}{\partial x_i} \\ \frac{\partial x_v}{\partial y_i} & \frac{\partial y_v}{\partial y_i} & \frac{\partial z_v}{\partial y_i} \\ \frac{\partial x_v}{\partial z_i} & \frac{\partial y_v}{\partial z_i} & \frac{\partial z_v}{\partial z_i} \end{bmatrix} \mathbf{F}_d \quad (4.123)$$

where \mathbf{F}_d is the force on the virtual site and x_v, y_v and z_v are the coordinates of the virtual site. In this way the total force and the total torque are conserved [69].

As a further note, the computation of the virial (eqn. 3.19) is non-trivial when virtual sites are used. Since the virial involves a summation over all the atoms (rather than virtual sites) the forces must be redistributed from the virtual sites to the atoms (using eqn. 4.123) *before* computation of the virial. In some special cases where the forces on the atoms can be written as a linear combination of the forces on the virtual sites (types 2 and 3 below) there is no difference between computing the virial before and after the redistribution of forces. However, in the general case redistribution should be done first.

There are six ways to construct virtual sites from surrounding atoms in GROMACS, which we classify by the number of constructing atoms. Note that all site types mentioned can be constructed

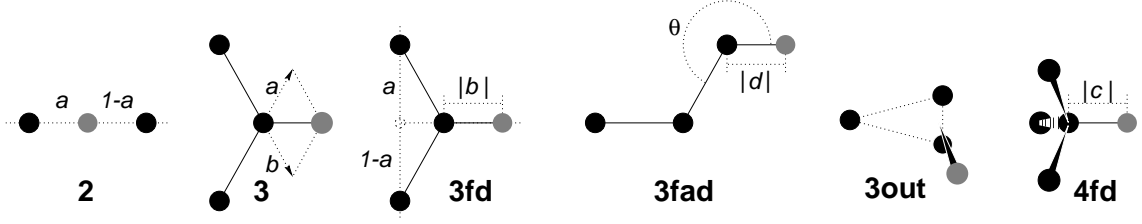


Figure 4.16: The six different types of virtual site construction in GROMACS. The constructing atoms are shown as black circles, the virtual sites in grey.

from types 3fd (normalized, in-plane) and 3out (non-normalized, out of plane). However, the amount of computation involved increases sharply along this list, so we strongly recommended using the first adequate virtual site type that will be sufficient for a certain purpose. Fig. 4.16 gives an overview of the available virtual site constructions.

2. As a linear combination of two atoms (Fig. 4.16 2):

$$\mathbf{r}_v = \mathbf{r}_i + a\mathbf{r}_{ij} \quad (4.124)$$

In this case the virtual site is on the line through atoms i and j . The force on particles i and j due to the force on the virtual site can be computed as:

$$\begin{aligned} \mathbf{F}'_i &= (1-a)\mathbf{F}_d \\ \mathbf{F}'_j &= a\mathbf{F}_d \end{aligned} \quad (4.125)$$

3. As a linear combination of three atoms (Fig. 4.16 3):

$$\mathbf{r}_v = \mathbf{r}_i + a\mathbf{r}_{ij} + b\mathbf{r}_{ik} \quad (4.126)$$

In this case the virtual site is in the plane of the other three particles. The force on particles i , j and k due to the force on the virtual site can be computed as:

$$\begin{aligned} \mathbf{F}'_i &= (1-a-b)\mathbf{F}_d \\ \mathbf{F}'_j &= a\mathbf{F}_d \\ \mathbf{F}'_k &= b\mathbf{F}_d \end{aligned} \quad (4.127)$$

- 3fd. In the plane of three atoms, with a fixed distance (Fig. 4.16 3fd):

$$\mathbf{r}_v = \mathbf{r}_i + b \frac{\mathbf{r}_{ij} + a\mathbf{r}_{jk}}{|\mathbf{r}_{ij} + a\mathbf{r}_{jk}|} \quad (4.128)$$

In this case the virtual site is in the plane of the other three particles at a distance of $|b|$ from i . The force on particles i , j and k due to the force on the virtual site can be computed as:

$$\begin{aligned} \mathbf{F}'_i &= \mathbf{F}_d - \gamma(\mathbf{F}_d - \mathbf{p}) \\ \mathbf{F}'_j &= (1-a)\gamma(\mathbf{F}_d - \mathbf{p}) \\ \mathbf{F}'_k &= a\gamma(\mathbf{F}_d - \mathbf{p}) \end{aligned} \quad \text{where} \quad \begin{aligned} \gamma &= \frac{b}{|\mathbf{r}_{ij} + a\mathbf{r}_{jk}|} \\ \mathbf{p} &= \frac{\mathbf{r}_{id} \cdot \mathbf{F}_d}{\mathbf{r}_{id} \cdot \mathbf{r}_{id}} \mathbf{r}_{id} \end{aligned} \quad (4.129)$$

3fad. In the plane of three atoms, with a fixed angle and distance (Fig. 4.16 3fad):

$$\mathbf{r}_v = \mathbf{r}_i + d \cos \theta \frac{\mathbf{r}_{ij}}{|\mathbf{r}_{ij}|} + d \sin \theta \frac{\mathbf{r}_\perp}{|\mathbf{r}_\perp|} \quad \text{where} \quad \mathbf{r}_\perp = \mathbf{r}_{jk} - \frac{\mathbf{r}_{ij} \cdot \mathbf{r}_{jk}}{\mathbf{r}_{ij} \cdot \mathbf{r}_{ij}} \mathbf{r}_{ij} \quad (4.130)$$

In this case the virtual site is in the plane of the other three particles at a distance of $|d|$ from i at an angle of α with \mathbf{r}_{ij} . Atom k defines the plane and the direction of the angle. Note that in this case b and α must be specified, instead of a and b (see also sec. 5.2.2). The force on particles i , j and k due to the force on the virtual site can be computed as (with \mathbf{r}_\perp as defined in eqn. 4.130):

$$\begin{aligned} \mathbf{F}'_i &= \mathbf{F}_d - \frac{d \cos \theta}{|\mathbf{r}_{ij}|} \mathbf{F}_1 + \frac{d \sin \theta}{|\mathbf{r}_\perp|} \left(\frac{\mathbf{r}_{ij} \cdot \mathbf{r}_{jk}}{\mathbf{r}_{ij} \cdot \mathbf{r}_{ij}} \mathbf{F}_2 + \mathbf{F}_3 \right) \\ \mathbf{F}'_j &= \frac{d \cos \theta}{|\mathbf{r}_{ij}|} \mathbf{F}_1 - \frac{d \sin \theta}{|\mathbf{r}_\perp|} \left(\mathbf{F}_2 + \frac{\mathbf{r}_{ij} \cdot \mathbf{r}_{jk}}{\mathbf{r}_{ij} \cdot \mathbf{r}_{ij}} \mathbf{F}_2 + \mathbf{F}_3 \right) \\ \mathbf{F}'_k &= \frac{d \sin \theta}{|\mathbf{r}_\perp|} \mathbf{F}_2 \end{aligned}$$

$$\text{where } \mathbf{F}_1 = \mathbf{F}_d - \frac{\mathbf{r}_{ij} \cdot \mathbf{F}_d}{\mathbf{r}_{ij} \cdot \mathbf{r}_{ij}} \mathbf{r}_{ij}, \quad \mathbf{F}_2 = \mathbf{F}_1 - \frac{\mathbf{r}_\perp \cdot \mathbf{F}_d}{\mathbf{r}_\perp \cdot \mathbf{r}_\perp} \mathbf{r}_\perp \quad \text{and} \quad \mathbf{F}_3 = \frac{\mathbf{r}_{ij} \cdot \mathbf{F}_d}{\mathbf{r}_{ij} \cdot \mathbf{r}_{ij}} \mathbf{r}_\perp \quad (4.131)$$

3out. As a non-linear combination of three atoms, out of plane (Fig. 4.16 3out):

$$\mathbf{r}_v = \mathbf{r}_i + a\mathbf{r}_{ij} + b\mathbf{r}_{ik} + c(\mathbf{r}_{ij} \times \mathbf{r}_{ik}) \quad (4.132)$$

This enables the construction of virtual sites out of the plane of the other atoms. The force on particles i , j and k due to the force on the virtual site can be computed as:

$$\begin{aligned} \mathbf{F}'_j &= \begin{bmatrix} a & -c z_{ik} & c y_{ik} \\ c z_{ik} & a & -c x_{ik} \\ -c y_{ik} & c x_{ik} & a \end{bmatrix} \mathbf{F}_d \\ \mathbf{F}'_k &= \begin{bmatrix} b & c z_{ij} & -c y_{ij} \\ -c z_{ij} & b & c x_{ij} \\ c y_{ij} & -c x_{ij} & b \end{bmatrix} \mathbf{F}_d \\ \mathbf{F}'_i &= \mathbf{F}_d - \mathbf{F}'_j - \mathbf{F}'_k \end{aligned} \quad (4.133)$$

4fd. From four atoms, with a fixed distance (Fig. 4.16 4fd):

$$\mathbf{r}_v = \mathbf{r}_i + c \frac{\mathbf{r}_{ij} + a\mathbf{r}_{jk} + b\mathbf{r}_{jl}}{|\mathbf{r}_{ij} + a\mathbf{r}_{jk} + b\mathbf{r}_{jl}|} \quad (4.134)$$

In this case the virtual site is at a distance of $|c|$ from i . The force on particles i , j , k and l due to the force on the virtual site can be computed as:

$$\begin{aligned} \mathbf{F}'_i &= \mathbf{F}_d - \gamma(\mathbf{F}_d - \mathbf{p}) \\ \mathbf{F}'_j &= (1 - a - b)\gamma(\mathbf{F}_d - \mathbf{p}) \\ \mathbf{F}'_k &= a\gamma(\mathbf{F}_d - \mathbf{p}) \\ \mathbf{F}'_l &= b\gamma(\mathbf{F}_d - \mathbf{p}) \end{aligned} \quad \text{where} \quad \gamma = \frac{c}{|\mathbf{r}_{ij} + a\mathbf{r}_{jk} + b\mathbf{r}_{jl}|} \quad (4.135)$$

$$\mathbf{p} = \frac{\mathbf{r}_{id} \cdot \mathbf{F}_d}{\mathbf{r}_{id} \cdot \mathbf{r}_{id}} \mathbf{r}_{id}$$

4.8 Dispersion correction

In this section we derive long range corrections due to the use of a cut-off for Lennard Jones or Buckingham interactions. We assume that the cut-off is so long that the repulsion term can safely be neglected, and therefore only the dispersion term is taken into account. Due to the nature of the dispersion interaction, energy and pressure corrections both are negative. While the energy correction is usually small, it may be important for free energy calculations. The pressure correction in contrast is very large and can not be neglected. Although it is in principle possible to parameterize a force field such that the pressure is close to 1 bar even without correction, such a method makes the parameterization dependent on the cut-off and is therefore undesirable. Please note that it is not consistent to use the long range correction to the dispersion without using either a reaction field method or a proper long range electrostatics method such as Ewald summation or PPPM.

4.8.1 Energy

The long range contribution of the dispersion interaction to the virial can be derived analytically, if we assume a homogeneous system beyond the cut-off distance r_c . The dispersion energy between two particles is written as:

$$V(r_{ij}) = -C_6 r_{ij}^{-6} \quad (4.136)$$

and the corresponding force is

$$\mathbf{F}_{ij} = -6 C_6 r_{ij}^{-8} \mathbf{r}_{ij} \quad (4.137)$$

In a periodic system it is not easy to calculate the full potentials, so usually a cut-off is applied, which can be abrupt or smooth. We will call the potential and force with cut-off V_c and \mathbf{F}_c . The long-range contribution to the dispersion energy in a system with N particles and particle density $\rho = N/V$ is:

$$V_{lr} = \frac{1}{2} N \rho \int_0^\infty 4\pi r^2 g(r) (V(r) - V_c(r)) dr \quad (4.138)$$

We will integrate this for the shift function, which is the most general form of Van der Waals interaction available in Gromacs. The shift function has a constant difference S from 0 to r_1 and is 0 beyond the cut-off distance r_c . We can integrate eqn. 4.138 assuming that the density in the sphere within r_1 is equal to the global density and the radial distribution function $g(r)$ is 1 beyond r_1 :

$$\begin{aligned} V_{lr} &= \frac{1}{2} N \left(\rho \int_0^{r_1} 4\pi r^2 g(r) C_6 S dr + \rho \int_{r_1}^{r_c} 4\pi r^2 (V(r) - V_c(r)) dr + \rho \int_{r_c}^\infty 4\pi r^2 V(r) dr \right) \\ &= \frac{1}{2} N \left(\left(\frac{4}{3} \pi \rho r_1^3 - 1 \right) C_6 S + \rho \int_{r_1}^{r_c} 4\pi r^2 (V(r) - V_c(r)) dr - \frac{4}{3} \pi N \rho C_6 r_c^{-3} \right) \quad (4.139) \end{aligned}$$

where the term -1 corrects for the self-interaction. For a plain cut-off we only need to assume that $g(r)$ is 1 beyond r_c and the correction reduces to [70]:

$$V_{lr} = -\frac{2}{3} \pi N \rho C_6 r_c^{-3} \quad (4.140)$$

If we consider for example a box of pure water, simulated with a cut-off of 0.9 nm and a density of 1 g cm^{-3} this correction is $-0.75 \text{ kJ mol}^{-1}$ per molecule.

For a homogeneous mixture we need to define an *average dispersion constant*:

$$\langle C_6 \rangle = \frac{2}{N(N-1)} \sum_i^N \sum_{j>i}^N C_6(i, j) \quad (4.141)$$

In GROMACS excluded pairs of atoms do not contribute to the average.

In the case of inhomogeneous simulation systems, *e.g.* a system with a lipid interface, the energy correction can be applied if $\langle C_6 \rangle$ for both components is comparable.

4.8.2 Virial and pressure

The scalar virial of the system due to the dispersion interaction between two particles i and j is given by:

$$\Xi = -\frac{1}{2} \mathbf{r}_{ij} \cdot \mathbf{F}_{ij} = 3 C_6 r_{ij}^{-6} \quad (4.142)$$

The pressure is given by:

$$P = \frac{2}{3V} (E_{kin} - \Xi) \quad (4.143)$$

The long-range correction to the virial is given by:

$$\Xi_{lr} = \frac{1}{2} N \rho \int_0^\infty 4\pi r^2 g(r) (\Xi - \Xi_c) dr \quad (4.144)$$

We can again integrate the long range contribution to the virial assuming $g(r)$ is 1 beyond r_1 :

$$\begin{aligned} \Xi_{lr} &= \frac{1}{2} N \rho \left(\int_{r_1}^{r_c} 4\pi r^2 (\Xi - \Xi_c) dr + \int_{r_c}^\infty 4\pi r^2 3 C_6 r_{ij}^{-6} dr \right) \\ &= \frac{1}{2} N \rho \left(\int_{r_1}^{r_c} 4\pi r^2 (\Xi - \Xi_c) dr + 4\pi C_6 r_c^{-3} \right) \end{aligned} \quad (4.145)$$

For a plain cut-off the correction to the pressure is [70]:

$$P_{lr} = -\frac{4}{3} \pi C_6 \rho^2 r_c^{-3} \quad (4.146)$$

Using the same example of a water box, the correction to the virial is 0.75 kJ mol^{-1} per molecule, the corresponding correction to the pressure for SPC water is approximately -280 bar.

For homogeneous mixtures we can again use the average dispersion constant $\langle C_6 \rangle$ (eqn. 4.141):

$$P_{lr} = -\frac{4}{3} \pi \langle C_6 \rangle \rho^2 r_c^{-3} \quad (4.147)$$

For inhomogeneous systems eqn. 4.147 can be applied under the same restriction as holds for the energy (see sec. 4.8.1).

4.9 Long Range Electrostatics

4.9.1 Ewald summation

The total electrostatic energy of N particles and the periodic images are given by

$$V = \frac{f}{2} \sum_{n_x} \sum_{n_y} \sum_{n_z^*} \sum_i^N \sum_j^N \frac{q_i q_j}{\mathbf{r}_{ij,\mathbf{n}}}. \quad (4.148)$$

$(n_x, n_y, n_z) = \mathbf{n}$ is the box index vector, and the star indicates that terms with $i = j$ should be omitted when $(n_x, n_y, n_z) = (0, 0, 0)$. The distance $\mathbf{r}_{ij,\mathbf{n}}$ is the real distance between the charges and not the minimum-image. This sum is conditionally convergent, but very slow.

Ewald summation was first introduced as a method to calculate long-range interactions of the periodic images in crystals [71]. The idea is to convert the single slowly-converging sum eqn. 4.148 into two quickly-converging terms and a constant term:

$$V = V_{dir} + V_{rec} + V_0 \quad (4.149)$$

$$V_{dir} = \frac{f}{2} \sum_{i,j}^N \sum_{n_x} \sum_{n_y} \sum_{n_z^*} q_i q_j \frac{\text{erfc}(\beta r_{ij,\mathbf{n}})}{r_{ij,\mathbf{n}}} \quad (4.150)$$

$$V_{rec} = \frac{f}{2\pi V} \sum_{i,j}^N q_i q_j \sum_{m_x} \sum_{m_y} \sum_{m_z^*} \frac{\exp(-(\pi \mathbf{m}/\beta)^2 + 2\pi i \mathbf{m} \cdot (\mathbf{r}_i - \mathbf{r}_j))}{\mathbf{m}^2} \quad (4.151)$$

$$V_0 = -\frac{f\beta}{\sqrt{\pi}} \sum_i^N q_i^2, \quad (4.152)$$

where β is a parameter that determines the relative weight of the direct and reciprocal sums and $\mathbf{m} = (m_x, m_y, m_z)$. In this way we can use a short cutoff (of the order of 1 nm) in the direct space sum and a short cutoff in the reciprocal space sum (*e.g.* 10 wave vectors in each direction). Unfortunately, the computational cost of the reciprocal part of the sum increases as N^2 (or $N^{3/2}$ with a slightly better algorithm) and it is therefore not realistic for use in large systems.

Using Ewald

Don't use Ewald unless you are absolutely sure this is what you want - for almost all cases the PME method below will perform much better. If you still want to employ classical Ewald summation enter this in your `.mdp` file, if the side of your box is about 3 nm:

```
coulombtype = Ewald
rvdw = 0.9
rlist = 0.9
rcoulomb = 0.9
fourierspacing = 0.6
ewaldrtol = 1e-5
```

The `fourierspacing` parameter times the box dimensions determines the highest magnitude of wave vectors m_x, m_y, m_z to use in each direction. With a 3 nm cubic box this example would use 11 wave vectors (from -5 to 5) in each direction. The `ewald_rtol` parameter is the relative strength of the electrostatic interaction at the cutoff. Decreasing this gives you a more accurate direct sum, but a less accurate reciprocal sum.

4.9.2 PME

Particle-mesh Ewald is a method proposed by Tom Darden [9, 10] to improve the performance of the reciprocal sum. Instead of directly summing wave vectors, the charges are assigned to a grid using cardinal B-spline interpolation. This grid is then Fourier transformed with a 3D FFT algorithm and the reciprocal energy term obtained by a single sum over the grid in k-space.

The potential at the grid points is calculated by inverse transformation, and by using the interpolation factors we get the forces on each atom.

The PME algorithm scales as $N \log(N)$, and is substantially faster than ordinary Ewald summation on medium to large systems. On very small systems it might still be better to use Ewald to avoid the overhead in setting up grids and transforms.

Using PME

To use Particle-mesh Ewald summation in GROMACS, specify the following lines in your `.mdp` file:

```
coulombtype = PME
rvdw = 0.9
rlist = 0.9
rcoulomb = 0.9
fourierspacing = 0.12
pme_order = 4
ewald_rtol = 1e-5
```

In this case the `fourierspacing` parameter determines the maximum spacing for the FFT grid and `pme_order` controls the interpolation order. Using 4th order (cubic) interpolation and this spacing should give electrostatic energies accurate to about $5 \cdot 10^{-3}$. Since the Lennard-Jones energies are not this accurate it might even be possible to increase this spacing slightly.

Pressure scaling works with PME, but be aware of the fact that anisotropic scaling can introduce artificial ordering in some systems.

4.9.3 PPPM

The Particle-Particle Particle-Mesh methods of Hockney & Eastwood can also be applied in GROMACS for the treatment of long range electrostatic interactions [72, 9, 73]. With this algorithm the charges of all particles are spread over a grid of dimensions (n_x, n_y, n_z) using a weighting function

called the triangle-shaped charged distribution:

$$W(\mathbf{r}) = W(x) W(y) W(z)$$

$$W(\xi) = \begin{cases} \frac{3}{4} - \left(\frac{\xi}{h}\right)^2 & |\xi| \leq \frac{h}{2} \\ \frac{1}{2} \left(\frac{3}{2} - \frac{|\xi|}{h}\right)^2 & \frac{h}{2} < |\xi| < \frac{3h}{2} \\ 0 & \frac{3h}{2} \leq |\xi| \end{cases} \quad (4.153)$$

where ξ (is x, y or z) is the distance to a grid point in the corresponding dimension. Only the 27 closest grid points need to be taken into account for each charge.

Then, this charge distribution is Fourier transformed using a 3D inverse FFT routine. In Fourier space a convolution with function \hat{G} is performed:

$$\hat{G}(k) = \frac{\hat{g}(k)}{\epsilon_0 k^2} \quad (4.154)$$

where \hat{g} is the Fourier transform of the charge spread function $g(\mathbf{r})$. This yields the long range potential $\hat{\phi}(k)$ on the mesh, which can be transformed using a forward FFT routine into the real space potential. Finally the potential and forces are retrieved using interpolation [73]. It is not easy to calculate the full long-range virial tensor with PPPM, but it is possible to obtain the trace. This means that the sum of the pressure components is correct (and therefore the isotropic pressure) but not necessarily the individual pressure components!

Using PPPM

To use the PPPM algorithm in GROMACS, specify the following lines in your `.mdp` file:

```
coulombtype = PPPM
rlist = 1.0
rcoulomb = 0.85
rcoulomb_switch = 0.0
rvdw = 1.0
fourierspacing = 0.075
```

For details on the switch parameters see the section on modified long-range interactions in this manual. When using PPPM we recommend to take at most 0.075 nm per gridpoint (*e.g.* 20 gridpoints for 1.5 nm). PPPM does not provide the same accuracy as PME but can be slightly faster in some cases. Due to the problem with the pressure tensor you shouldn't use it with pressure coupling.

We're somewhat ambivalent about PPPM, so if you use it please contact us - otherwise it might be removed from future releases so we can concentrate our efforts on PME.

4.9.4 Optimizing Fourier transforms

To get the best possible performance you should try to avoid large prime numbers for grid dimensions. The FFT code used in GROMACS is optimized for grid sizes of the form $2^a 3^b 5^c 7^d 11^e 13^f$,

where $e + f$ is 0 or 1 and the other exponents arbitrary. (See further the documentation of the FFT algorithms at www.fftw.org.)

It is also possible to optimize the transforms for the current problem by performing some calculations at the start of the run. This is not done per default since it takes a couple of minutes, but for large runs it will save time. Turn it on by specifying

```
optimize_fft = yes  
in your .mdp file.
```

When running in parallel the grid must be communicated several times and thus hurting scaling performance. With PME you can improve this by increasing grid spacing while simultaneously increasing the interpolation to *e.g.* 6th order. Since the interpolation is entirely local a this will improve the scaling in most cases.

4.10 All-hydrogen force-field

The GROMACS all-hydrogen force-field is almost identical to the normal GROMACS forcefield, since the extra hydrogens have no Lennard-Jones interaction and zero charge. The only differences are in the bond angle and improper dihedral angle terms. This forcefield is only useful when you need the exact hydrogen positions, for instance for distance restraints derived from NMR measurements.

4.11 GROMOS-96 notes

4.11.1 The GROMOS-96 force field

GROMACS supports the GROMOS-96 force fields [53]. All parameters for the 43a1, 43a2 (development, improved alkane dihedrals) and 43b1 (vacuum) force fields are included. All standard building blocks are included and topologies can be build automatically by `pdb2gmx`. The GROMOS-96 force field is a further development of the GROMOS-87 force field on which the GROMACS forcefield is based. The GROMOS-96 force field has improvements over the GROMACS force field for proteins and small molecules. It is not, however, recommended for use with long alkanes and lipids. The GROMOS-96 force field differs from the GROMACS force field in a few aspects:

- the force field parameters
- the parameters for the bonded interactions are not linked to atom types
- a fourth power bond stretching potential (sec. 4.2.1)
- an angle potential based on the cosine of the angle (sec. 4.2.5)

There are two differences in implementation between GROMACS and GROMOS-96 which can lead to slightly different results when simulating the same system with both packages:

- in GROMOS-96 neighbor searching for solvents is performed on the first atom of the solvent molecule, this is not implemented in GROMACS, but the difference with searching with centers of charge groups is very small
- the virial in GROMOS-96 is molecule-based. This is not implemented in GROMACS, which uses atomic virials

The GROMOS-96 force field was parameterized with a Lennard-Jones cutoff of 1.4 nm, so be sure to use a Lennard-Jones cutoff of at least 1.4. A larger cutoff is possible, because the Lennard-Jones potential and forces are almost zero beyond 1.4 nm.

4.11.2 GROMOS-96 files

GROMACS can read and write GROMOS-96 coordinate and trajectory files. These files should have the extension `.g96`. Such a file can be a GROMOS-96 initial/final configuration file or a coordinate trajectory file or a combination of both. The file is fixed format; all floats are written as 15.9 (files can get huge). GROMACS supports the following data blocks in the given order:

- Header block: `TITLE` (mandatory)
- Frame blocks: `TIMESTEP` (optional)
`POSITION/POSITIONRED` (mandatory)
`VELOCITY/VELOCITYRED` (optional)
`BOX` (optional)

See the GROMOS-96 manual [53] for a complete description of the blocks. Note that all GROMACS programs can read compressed (`.Z`) or gzipped (`.gz`) files.

Chapter 5

Topologies

5.1 Introduction

GROMACS must know on which atoms and combinations of atoms the various contributions to the potential functions (see chapter 4) must act. It must also know what parameters must be applied to the various functions. All this is described in the *topology* file `*.top`, which lists the *constant attributes* of each atom. There are many more atom types than elements, but only atom types present in biological systems are parameterized in the force field, plus some metals, ions and silicon. The bonded and special interactions are determined by fixed lists that are included in the topology file. Certain non-bonded interactions must be excluded (first and second neighbors), as these are already treated in bonded interactions. In addition there are *dynamic attributes* of atoms: their positions, velocities and forces, but these do not strictly belong to the molecular topology.

This Chapter describes the set up of the topology file, the `*.top` file and the database files: what the parameters stand for and how/where to change them if needed. First all file formats are explained. Section 5.7.1 describes the organization of the force-field files.

Note: if you construct your own topologies, we encourage you to upload them to our topology archive at www.gromacs.org! Just imagine how thankful you'd have been if your topology had been available there before you started. The same goes for new force field or modified versions of the standard force fields - contribute them to the force field archive!

5.2 Particle type

In GROMACS there are 5 types of particles, see Table 5.1. Only regular atoms and virtual interaction-sites are used in GROMACS; shells are necessary for polarizable models like the Shell-Water models [29].

Particle	Symbol
atoms	A
shells	S
virtual interaction-sites	V (or D)

Table 5.1: Particle types in GROMACS

5.2.1 Atom types

Depending on the force field GROMACS uses different atom types, a sample from the deprecated “gromacs” force field is listed below, with their corresponding masses (in a.m.u.). This is the same listing as in the file `ff???.atp` (.atp = **atom type parameter file**), therefore in this file you can change and/or add an atom type.

O	15.99940	; carbonyl oxygen (C=O)
OM	15.99940	; carboxyl oxygen (CO-)
OA	15.99940	; hydroxyl oxygen (OH)
OW	15.99940	; water oxygen
N	14.00670	; peptide nitrogen (N or NH)
NT	14.00670	; terminal nitrogen (NH2)
NL	14.00670	; terminal nitrogen (NH3)
NR5	14.00670	; aromatic N (5-ring,2 bonds)
NR5*	14.00670	; aromatic N (5-ring,3 bonds)
NP	14.00670	; porphyrin nitrogen
C	12.01100	; bare carbon (peptide,C=O,C-N)
CH1	13.01900	; aliphatic CH-group
CH2	14.02700	; aliphatic CH2-group
CH3	15.03500	; aliphatic CH3-group

Atomic detail is used except for hydrogen atoms bound to (aliphatic) carbon atoms, which are treated as *united atoms*. No special hydrogen-bond term is included. Note that other force field like OPLS/AA and Amber99 use all atoms.

For backward compatibility we retain here some reference to parameters present in the “gromacs” force field. The last 10 atom types were not part of the original GROMOS-87 force field [45] and when you use them you can refer to one or more of the following papers:

- F was taken from ref. [49],
- CP2 and CP3 from ref. [46] and references cited therein,
- CR5, CR6 and HCR from ref. [74]
- OWT3 from ref. [48]
- SD, OD and CD from ref. [50]

Note that we recommend against using these parameters in new projects since they are not well-tested.

Note: GROMACS makes use of the atom types as a name, *not* as a number (as *e.g.* in GROMOS).

5.2.2 Virtual sites

Some force fields use virtual interaction-sites (interaction sites that are constructed from other particle positions) on which certain interactions are located (*e.g.* on benzene rings, to reproduce the correct quadrupole). This is described in sec. 4.7.

To make virtual sites in your system, you should include a section [`virtual_sites?`] (for backward compatibility the old name [`dummies?`] can also be used) in your topology file, where the ‘?’ stands for the number constructing particles for the virtual site. This will be ‘2’ for type 2, ‘3’ for types 3, 3fd, 3fad and 3out and ‘4’ for type 4fd (the different types are explained in sec. 4.7).

Parameters for type 2 should look like this:

```
[ virtual_sites2 ]
; Site from funct a
5 1 2 1 0.7439756
```

for type 3 like this:

```
[ virtual_sites3 ]
; Site from funct a b
5 1 2 3 1 0.7439756 0.128012
```

for type 3fd like this:

```
[ virtual_sites3 ]
; Site from funct a d
5 1 2 3 2 0.5 -0.105
```

for type 3fad like this:

```
[ virtual_sites3 ]
; Site from funct theta d
5 1 2 3 3 120 0.5
```

for type 3out like this:

```
[ virtual_sites3 ]
; Site from funct a b c
5 1 2 3 4 -0.4 -0.4 6.9281
```

for type 4fd like this:

```
[ virtual_sites4 ]
; Site from funct a b d
5 1 2 3 4 1 0.33333 0.33333 -0.105
```

This will result in the construction of a virtual site, number 5 (first column ‘Site’), based on the positions of 1 and 2 or 1, 2 and 3 or 1, 2, 3 and 4 (next two, three or four columns ‘from’) following the rules determined by the function number (next column ‘funct’) with the parameters

Property	Symbol	Unit
Type	-	-
Mass	m	a.m.u.
Charge	q	electron
epsilon	ϵ	kJ/mol
sigma	σ	nm

Table 5.2: Static atom type properties in GROMACS

specified (last one, two or three columns ‘a b .’).

Note that if any constant bonded interactions defined between virtual sites and/or normal atoms will be removed by `grompp`, this happens after the exclusions have been generated. This way, exclusions will not be affected by an atom being defined as virtual site or not, but by the bonding configuration of the atom.

5.3 Parameter files

5.3.1 Atoms

A number of *static* properties are assigned to the atom types in the GROMACS force field: Type, Mass, Charge, ϵ and σ (see Table 5.2). The mass is listed in `ff???.atp` (see 5.2.1), whereas the charge is listed in `ff???.rtp` (`rtp` = residue topology parameter file, see 5.5.1). This implies that the charges are only defined in the building blocks of amino acids or user defined building blocks. When generating a topology (`*.top`) using the `pdb2gmx` program the information from these files is combined.

The following *dynamic* quantities are associated with an atom

- Position \mathbf{x}
- Velocity \mathbf{v}

These quantities are listed in the coordinate file, `*.gro` (see section File format, 5.6.6).

5.3.2 Bonded parameters

The bonded parameters (*i.e.* bonds, bond angles, improper and proper dihedrals) are listed in `ff???.bon.itp`. The term `func` is 1 for harmonic and 2 for GROMOS-96 bond and angle potentials. For the dihedral, this is explained after this listing.

```
[ bondtypes ]
; i j func b0 kb
C O 1 0.12300 502080.
C OM 1 0.12500 418400.
.....
```

```

[ angletypes ]
; i j k func th0 cth
HO OA C 1 109.500 397.480
HO OA CH1 1 109.500 397.480
.....

[ dihedraltypes ]
; i l func q0 cq
NR5* NR5 2 0.000 167.360
NR5* NR5* 2 0.000 167.360
.....

[ dihedraltypes ]
; j k func phi0 cp mult
C OA 1 180.000 16.736 2
C N 1 180.000 33.472 2
.....

[ dihedraltypes ]
;
; Ryckaert-Bellemans Dihedrals
;
; aj ak funct
CP2 CP2 3 9.2789 12.156 -13.120 -3.0597 26.240 -31.495

```

Also in this file are the Ryckaert-Bellemans [75] parameters for the CP2-CP2 dihedrals in alkanes or alkane tails with the following constants:

$$\begin{array}{rclclcl}
 & & & & & & \text{(kJ/mol)} \\
 C_0 & = & 9.28 & C_2 & = & -13.12 & C_4 & = & 26.24 \\
 C_1 & = & 12.16 & C_3 & = & -3.06 & C_5 & = & -31.5
 \end{array}$$

(Note: The use of this potential implies the exclusion of LJ interactions between the first and the last atom of the dihedral, and ψ is defined according to the 'polymer convention' ($\psi_{trans} = 0$)).

So there are three types of dihedrals in the GROMACS force field:

- proper dihedral : funct = 1, with mult = multiplicity, so the number of possible angles
- improper dihedral : funct = 2
- Ryckaert-Bellemans dihedral : funct = 3

In the file `ff???bon.itp` you can add bonded parameters. If you want to include parameters for new atom types, make sure you define this new atom type in `ff???atp` as well.

5.3.3 Non-bonded parameters

The non-bonded parameters consist of the Van der Waals parameters V (c6) and W (c12), as listed in the file `ff???nb.itp`, where `p`type is the particle type (see Table 5.1):

```
[ atomtypes ]
; name mass charge ptype c6 c12
O 15.99940 0.000 A 0.22617E-02 0.74158E-06
OM 15.99940 0.000 A 0.22617E-02 0.74158E-06
.....

[ nonbond_params ]
; i j func c6 c12
O O 1 0.22617E-02 0.74158E-06
O OA 1 0.22617E-02 0.13807E-05
.....

[ pairtypes ]
; i j func cs6 cs12 ; THESE ARE 1-4 INTERACTIONS
O O 1 0.22617E-02 0.74158E-06
O OM 1 0.22617E-02 0.74158E-06
.....
```

The parameters V and W can be defined in two different ways, depending on the combination rule that was chosen in the `[defaults]` section of the topology file (see 5.6.1):

$$\text{for combination rule 1 : } \begin{aligned} V_{ii} &= C_i^{(6)} = 4 \epsilon_i \sigma_i^6 \quad [\text{kJ mol}^{-1} \text{ nm}^6] \\ W_{ii} &= C_i^{(12)} = 4 \epsilon_i \sigma_i^{12} \quad [\text{kJ mol}^{-1} \text{ nm}^{12}] \end{aligned} \quad (5.1)$$

$$\text{for combination rules 2 and 3 : } \begin{aligned} V_{ii} &= \sigma_i \quad [\text{nm}] \\ W_{ii} &= \epsilon_i \quad [\text{kJ mol}^{-1}] \end{aligned} \quad (5.2)$$

Some or all combinations for different atom-types can be given in the `[nonbond_params]` section. Any combination that is not given will be computed according to the combination rule:

$$\text{for combination rules 1 and 3 : } \begin{aligned} C_{ij}^{(6)} &= \left(C_i^{(6)} C_j^{(6)} \right)^{\frac{1}{2}} \\ C_{ij}^{(12)} &= \left(C_i^{(12)} C_j^{(12)} \right)^{\frac{1}{2}} \end{aligned} \quad (5.3)$$

$$\text{for combination rule 2 : } \begin{aligned} \sigma_{ij} &= \frac{1}{2}(\sigma_i + \sigma_j) \\ \epsilon_{ij} &= \sqrt{\epsilon_i \epsilon_j} \end{aligned} \quad (5.4)$$

5.3.4 Pair interactions

Extra Lennard-Jones and electrostatic interactions between pairs of atoms in a molecule can be added in the `[pairs]` section of a molecule definition. The parameters for these interactions can be set independently from the non-bonded interaction parameters. In the GROMOS force

fields pairs are only used to modify the 1-4 interactions (interactions of atoms separated by three bonds). In these forcefields the 1-4 interactions are excluded from the non-bonded interactions (see 5.3.5).

The pair interaction parameters for the atom types in `ff???nb.itp` are listed in the `[pairtypes]` section. The GROMOS force fields list all these interactions explicitly, but this section might be empty for force fields like OPLS that calculate the 1-4 interactions by scaling. Pair parameters which are not present in the `[pairtypes]` section are only generated when `generate pairs` is set to `yes` in the topology (see 5.6.1). When `generate pairs` is set to `no`, `grompp` will give a warning for each pair type for which no parameters are given.

5.3.5 Exclusions

The exclusions for bonded particles are generated by `grompp` for neighboring atoms up to a certain number of bonds away, as defined in the `[moleculetype]` section in the topology file (see 5.6.1). Particles are considered bonded when they are connected by bonds (`[bonds]` types 1 to 5) or constraints (`[constraints]` type 1). `[bonds]` type 5 can be used to create a connection between two atoms without creating an interaction. There is a harmonic interaction (`[bonds]` type 6) which does not connect the atoms by a chemical bond. There is also a second constraint type (`[constraints]` type 2) which fixes the distance, but does not connect the atoms by a chemical bond. For a complete list of all these interactions see Table 5.4.

Extra exclusions within a molecule can be added manually in a `[exclusions]` section. Each line should start with one atom index, followed by one or more atom indices. All non-bonded interactions between the first atom and the other atoms will be excluded.

When all non-bonded interactions within or between groups of atoms need to be excluded, it is more convenient and much more efficient to use energy monitor group exclusions (see sec. 3.3).

5.4 Constraints

Constraints are defined in the `[constraints]` section. The format is two atom numbers followed by the function type, which can be 1 or 2 and the constraint distance. The only difference between the two types is that type 1 is used for generating exclusions and type 2 is not (see 5.3.5). The distances are constrained using the LINCS or the SHAKE algorithm, which can be selected in the `*.mdp` file. Both types of constraints can be perturbed in free-energy calculations by adding a second constraint distance (see 5.6.5). Several types of bonds and angles (see Table 5.4) can be converted automatically to constraints by `grompp`. There are several options for this in the `*.mdp` file.

We have also implemented the SETTLE algorithm [31] which is an analytical solution of SHAKE specifically for water. SETTLE can be selected in the topology file. Check for instance the SPC molecule definition:

```
[ moleculetype ]
; molname nrexcl
SOL 1
```

```
[ atoms ]
; nr at type res nr ren nm at nm cg nr charge
1 OW 1 SOL OW1 1 -0.82
2 HW 1 SOL HW2 1 0.41
3 HW 1 SOL HW3 1 0.41
```

```
[ settles ]
; OW funct doh dhh
1 1 0.1 0.16333
```

```
[ exclusions ]
1 2 3
2 1 3
3 1 2
```

The section `[settles]` defines the first atom of the watery molecule. The settle funct is always one, and the distance between O-H and H-H distances must be given. Note that the algorithm can also be used for TIP3P and TIP4P [48]. TIP3P just has another geometry. TIP4P has a virtual site, but since that is generated it does not need to be shaken (nor stirred).

5.5 Databases

5.5.1 Residue database

The file holding the residue database is `ff???.rtp`. Originally this file contained building blocks (amino acids) for proteins, and is the GROMACS interpretation of the `rt37c4.dat` file of GROMOS. So the residue file contains information (bonds, charge, charge groups and improper dihedrals) for a frequently used building block. It is better *not* to change this file because it is standard input for `pdb2gmx`, but if changes are needed make them in the `*.top` file (see 5.6.1). However, in the `ff???.rtp` file the user can define a new building block or molecule: see for example 2,2,2-trifluoroethanol (TFE) or *n*-decane (C10). But when defining new molecules (non-protein) it is preferable to create a `*.itp` file. This will be discussed in section 5.6.2. When adding a new protein residue to the database, don't forget to add the residue name to the `aminoacids.dat` file, so that `grompp`, `make_ndx` and analysis tools can recognize the residue as a protein residue (see 8.1.1).

The file `ff???.rtp` is only used by `pdb2gmx`. As mentioned before, the only extra information this program needs from `ff???.rtp` is bonds, charges of atoms, charge groups and improper dihedrals, because the rest is read from the coordinate input file (in the case of `pdb2gmx`, a `pdb` format file). Some proteins contain residues that are not standard, but are listed in the coordinate file. You have to construct a building block for this “strange” residue, otherwise you will not obtain a `*.top` file. This also holds for molecules in the coordinate file such as phosphate or sulphate ions. The residue database is constructed in the following way:

```
[ bondedtypes ] ; mandatory
; bonds angles dihedrals impropers
1 1 1 2 ; mandatory
```

```
[ GLY ] ; mandatory

[ atoms ] ; mandatory
; name type charge chargegroup
N N -0.280 0
H H 0.280 0
CA CH2 0.000 1
C C 0.380 2
O O -0.380 2

[ bonds ] ; optional
;atom1 atom2 b0 kb
N H
N CA
CA C
C O
-C N

[ exclusions ] ; optional
;atom1 atom2

[ angles ] ; optional
;atom1 atom2 atom3 th0 cth

[ dihedrals ] ; optional
;atom1 atom2 atom3 atom4 phi0 cp mult

[ impropers ] ; optional
;atom1 atom2 atom3 atom4 q0 cq
N -C CA H
-C -CA N -O

[ ZN ]

[ atoms ]
ZN ZN 2.000 0
```

The file is free format, the only restriction is that there can be at most one entry on a line. The first field in the file is the `[bondedtypes]` field, which is followed by four numbers, that indicate the interaction type for bonds, angles, dihedrals and improper dihedrals. The file contains residue entries, which consist of atoms and optionally bonds, angles dihedrals and impropers. The charge group codes denote the charge group numbers. Atoms in the same charge group should always be below each other. When using the hydrogen database with `pdb2gmx` for adding missing hydrogens, the atom names defined in the `.rtp` entry should correspond exactly to the naming

convention used in the hydrogen database, see 5.5.2. The atom names in the bonded interaction can be preceded by a minus or a plus, indicating that the atom is in the preceding or following residue respectively. Parameters can be added to bonds, angles, dihedrals and impropers, these parameters override the standard parameters in the `.itp` files. This should only be used in special cases. Instead of parameters, a string can be added for each bonded interaction, this is used in GROMOS96 `.rtp` files. These strings are copied to the topology file and can be replaced by force field parameters by the C-preprocessor in `grompp` using `#define` statements.

`pdb2gmx` automatically generates all angles. This means that for the GROMACS force field the `[angles]` field is only useful for overriding `.itp` parameters. For the GROMOS-96 force field the interaction number off all angles need to be specified.

`pdb2gmx` automatically generates one proper dihedral for every rotatable bond, preferably on heavy atoms. When the `[dihedrals]` field is used, no other dihedrals will be generated for the bonds corresponding to the specified dihedrals. It is possible to put more than one dihedral on a rotatable bond.

`pdb2gmx` sets the number of exclusions to 3, which means that interactions between atoms connected by at most 3 bonds are excluded. Pair interactions are generated for all pairs of atoms which are separated by 3 bonds (except pairs of hydrogens). When more interactions need to be excluded, or some pair interactions should not be generated, an `[exclusions]` field can be added, followed by pairs of atom names on separate lines. All non-bonded and pair interactions between these atoms will be excluded.

5.5.2 Hydrogen database

The hydrogen database is stored in `ff???.hdb`. It contains information for the `pdb2gmx` program on how to connect hydrogen atoms to existing atoms. In versions of the database before GROMACS 3.3, hydrogen atoms were named after the atom they are connected to: the first letter of the atom name was replaced by an 'H'. In the versions from 3.3 onwards, the H atom has to be listed explicitly, because the old behaviour was protein-specific and hence could not be generalized to other molecules. If more than one hydrogen atom is connected to the same atom, a number will be added to the end of the hydrogen atom name. For example, adding two hydrogen atoms to ND2 (in asparagine), the hydrogen atoms will be named HD21 and HD22. This is important since atom naming in the `.rtp` file (see 5.5.1) must be the same. The format of the hydrogen database is as follows:

```
; res # additions
# H add type H i j k
ALA 1
1 1 H N -C CA
ARG 4
1 2 H N CA C
1 1 HE NE CD CZ
2 3 HH1 NH1 CZ NE
2 3 HH2 NH2 CZ NE
```

On the first line we see the residue name (ALA or ARG) and the number of additions. After that

follows one line for each addition, on which we see:

- The number of H atoms added
- The way of adding H atoms, can be any of
 - 1 *one planar hydrogen, e.g. rings or peptide bond*
one hydrogen atom (n) is generated, lying in the plane of atoms (i,j,k) on the plane bisecting angle (j-i-k) at a distance of 0.1 nm from atom i, such that the angles (n-i-j) and (n-i-k) are $> 90^\circ$
 - 2 *one single hydrogen, e.g. hydroxyl*
one hydrogen atom (n) is generated at a distance of 0.1 nm from atom i, such that angle (n-i-j)=109.5 degrees and dihedral (n-i-j-k)=trans
 - 3 *two planar hydrogens, e.g. -NH₂*
two hydrogens (n1,n2) are generated at a distance of 0.1 nm from atom i, such that angle (n1-i-j)=(n2-i-j)=120 degrees and dihedral (n1-i-j-k)=cis and (n2-i-j-k)=trans, such that names are according to IUPAC standards [76]
 - 4 *two or three tetrahedral hydrogens, e.g. -CH₃*
three (n1,n2,n3) or two (n1,n2) hydrogens are generated at a distance of 0.1 nm from atom i, such that angle (n1-i-j)=(n2-i-j)=(n3-i-j)=109.47°, dihedral (n1-i-j-k)=trans, (n2-i-j-k)=trans+120 and (n3-i-j-k)=trans+240 degrees
 - 5 *one tetrahedral hydrogen, e.g. C₃CH*
one hydrogen atom (n') is generated at a distance of 0.1 nm from atom i in tetrahedral conformation such that angle (n'-i-j)=(n'-i-k)=(n'-i-l)=109.47°
 - 6 *two tetrahedral hydrogens, e.g. C-CH₂-C*
two hydrogen atoms (n1,n2) are generated at a distance of 0.1 nm from atom i in tetrahedral conformation on the plane bisecting angle i-j-k with angle (n-i-n2)=(n1-i-j)=(n1-i-k)=109.5
 - 7 *two water hydrogens*
two hydrogens are generated around atom i according to SPC [55] water geometry. The symmetry axis will alternate between three coordinate axes in both directions
 - 10 *three water "hydrogens"*
two hydrogens are generated around atom i according to SPC [55] water geometry. The symmetry axis will alternate between three coordinate axes in both directions. In addition an extra particle is generated on the position of the oxygen. This is for use with four-atom water models such as TIP4P [48]
 - 10 *four water "hydrogens"*
Same as above, except that two additional particles are generated on the position of the oxygen. This is for use with five-atom water models such as TIP5P [77]
- The name of the new H atom
- Three or four control atoms (i,j,k,l), where the first always is the atom to which the H atoms are connected. The other two or three depend on the code selected (for water there is only one control atom).

5.5.3 Termini database

The termini databases are stored in `ff???-n.tdb` and `ff???-c.tdb` for the N- and C-termini respectively. They contain information for the `pdb2gmx` program on how to connect new atoms to existing ones, which atoms should be removed or changed and which bonded interactions should be added. The format of the is as follows (this is an example from the `ffgmx-c.tdb`):

```
[ None ]

[ COO- ]

[ replace ]
C C C 12.011 0.27

[ add ]
2 8 O C CA N
OM 15.9994 -0.635

[ delete ]
O

[ impropers ]
C O1 O2 CA
```

The file is organized in blocks, each with a header specifying the name of the block. These blocks correspond to different types of termini that can be added to a molecule. In this example `[None]` is the first block, corresponding to a terminus that leaves the molecule as it is; `[COO-]` is the second terminus type, corresponding to changing the terminal carbon atom into a deprotonated carboxyl group. Block names cannot be any of the following: `replace`, `add`, `delete`, `bonds`, `angles`, `dihedrals`, `impropers`; this would interfere with the parameters of the block, and would probably also be very confusing to human readers.

Per block the following options are present:

- `[replace]`
replace an existing atom by one with a different atom type, atom name, charge and/or mass. For each atom to be replaced on line should be entered with the following fields:
 - name of the atom to be replaced
 - new atom name
 - new atom type
 - new mass
 - new charge
- `[add]`
add new atoms. For each (group of) added atom(s), a two-line entry is necessary. The first line contains the same fields as an entry in the hydrogen database (name of the new

atom, number of atoms, type of addition, control atoms, see 5.5.2), but the possible types of addition are extended by two more, specifically for C-terminal additions:

8 *two carboxyl oxygens, -COO⁻*

two oxygens (n1,n2) are generated according to rule 3, at a distance of 0.136 nm from atom i and an angle (n1-i-j)=(n2-i-j)=117 degrees

9 *carboxyl oxygens and hydrogen, -COOH*

two oxygens (n1,n2) are generated according to rule 3, at distances of 0.123 nm and 0.125 nm from atom i for n1 and n2 resp. and angles (n1-i-j)=121 and (n2-i-j)=115 degrees. One hydrogen (n') is generated around n2 according to rule 2, where n-i-j and n-i-j-k should be read as n'-n2-i and n'-n2-i-j resp.

After this line another line follows which specifies the details of the added atom(s), in the same way as for replacing atoms, *i.e.*:

- atom type
- mass
- charge

Like in the hydrogen database (see 5.5.1), when more than one atom is connected to an existing one, a number will be appended to the end of the atom name. Note that, like in the hydrogen database the atom name is now on the same line as the control atoms, whereas it was at the beginning of the second line prior to GROMACS version 3.3.

- [delete]
delete existing atoms. One atom name per line.
- [bonds], [angles], [dihedrals] and [impropers]
add additional bonded parameters. The format is identical to that used in the `ff???.rtp`, see 5.5.1.

5.6 File formats

5.6.1 Topology file

The topology file is built following the GROMACS specification for a molecular topology. A `*.top` file can be generated by `pdb2gmx`. All possible entries in the topology file are listed in Table 5.3 and Table 5.4. Also listed are all the units of the parameters, which interactions can be perturbed for free energy calculations, which bonded interactions are used by `grompp` for generating exclusions and which bonded interactions can be converted to constraints by `grompp`.

Description of the file layout:

- semicolon (;) and newline surround comments
- on a line ending with \ the newline character is ignored.
- directives are surrounded by [and]

Parameters

interaction type	directive	# at.	f. tp	parameters	F. E.
<i>mandatory</i>	defaults			non-bonded function type; combination rule ^(a) ; generate pairs (no/yes); fudge LJ (); fudge QQ ()	
<i>mandatory</i>	atomtypes			atom type; m (u); q (e); particle type; $V^{(a)}$; $W^{(a)}$	
	bondtypes		(see Table 5.4, directive bonds)		
	constrainttypes		(see Table 5.4, directive constraints)		
	pairtypes		(see Table 5.4, directive pairs)		
	angletypes		(see Table 5.4, directive angles)		
proper dih.	dihedraltypes	2/4 ^(b)	1	ϕ_s (deg); k_ϕ (kJ mol ⁻¹); multiplicity	ϕ, k
improper dih.	dihedraltypes	2/4 ^(c)	2	ξ_0 (deg); k_ξ (kJ mol ⁻¹ rad ⁻²)	all
RB dihedral	dihedraltypes	2/4 ^(b)	3	$C_0, C_1, C_2, C_3, C_4, C_5$ (kJ mol ⁻¹)	all
LJ	nonbond_params	2	1	$V^{(a)}$; $W^{(a)}$	
Buckingham	nonbond_params	2	2	a (kJ mol ⁻¹); b (nm ⁻¹); c_6 (kJ mol ⁻¹ nm ⁶)	

Molecule definition(s)

<i>mandatory</i>	moleculetype			molecule name; exclude neighbors # bonds away for non-bonded interactions	
<i>mandatory</i>	atoms	1		atom type; residue number; residue name; atom name; charge group number; q (e); m (u)	type q, m
intramolecular interaction definitions as described in Table 5.4 (next page)					

System

<i>mandatory</i>	system			system name	
<i>mandatory</i>	molecules			molecule name; number of molecules	

'# at' is the number of atom types

'f. tp' is function type

'F. E.' indicates which parameters can be interpolated during free energy calculations

^(a) the combination rule determines the type of LJ parameters, see 5.3.3

^(b) the inner two or all four atoms in the dihedral

^(c) the outer two or all four atoms in the dihedral

For free energy calculations, the parameters for topology 'B' ($\lambda = 1$) should be added on the same line, after the normal parameters, in the same order as the normal parameters.

Table 5.3: The topology (* .top) file.

Intramolecular interaction definitions

interaction type	directive	# at.	f. tp	parameters	F. E.
bond	bonds ^(b,c)	2	1	b_0 (nm); k_b (kJ mol ⁻¹ nm ⁻²)	all
G96 bond	bonds ^(b,c)	2	2	b_0 (nm); k_b (kJ mol ⁻¹ nm ⁻⁴)	all
morse	bonds ^(b,c)	2	3	b_0 (nm); D (kJ mol ⁻¹); β (nm ⁻¹)	
cubic bond connection	bonds ^(b,c)	2	4	b_0 (nm); $C_{i=2,3}$ (kJ mol ⁻¹ nm ⁻ⁱ);	
harmonic pot.	bonds ^(b)	2	5		
FENE bond	bonds	2	6	b_0 (nm); k_b (kJ mol ⁻¹ nm ⁻²)	all
LJ/Coul. 1-4	bonds	2	7	b_m (nm); k_b (kJ mol ⁻¹ nm ⁻²)	
LJ/C. 1-4 A	pairs	2	1	$V^{(a)}$; $W^{(a)}$	all
LJ/C. pair A	pairs	2	2	$V^{(a)}$; $W^{(a)}$	
angle	pairs	2	3		
G96 angle	angles ^(c)	3	1	θ_0 (deg); k_θ (kJ mol ⁻¹ rad ⁻²)	all
quartic angle	angles ^(c)	3	2	θ_0 (deg); k_θ (kJ mol ⁻¹)	all
proper dih.	angles ^(c)	3	6	θ_0 (deg); $C_{i=0,1,2,3,4}$ (kJ mol ⁻¹ rad ⁻ⁱ)	
improper dih.	dihedrals	4	1	ϕ_s (deg); k_ϕ (kJ mol ⁻¹); multiplicity	ϕ, k
RB dihedral	dihedrals	4	2	ξ_0 (deg); k_ξ (kJ mol ⁻¹ rad ⁻²)	all
constraint	dihedrals	4	3	$C_0, C_1, C_2, C_3, C_4, C_5$ (kJ mol ⁻¹)	all
constr. n.c.	constraints ^(b)	2	1	b_0 (nm)	all
settle	constraints	2	2	b_0 (nm)	all
vsite2	settles	3	1	d_{OH}, d_{HH} (nm)	
vsite3	virtual_sites2	3	1	a ()	
vsite3fd	virtual_sites3	4	1	a, b ()	
vsite3fad	virtual_sites3	4	2	a (); d (nm)	
vsite3out	virtual_sites3	4	3	θ (deg); d (nm)	
vsite4fd	virtual_sites3	4	4	a, b (); c (nm ⁻¹)	
position res.	virtual_sites4	5	1	a, b (); d (nm);	
distance res.	position_restraints	1	1	k_x, k_y, k_z (kJ mol ⁻¹ nm ⁻²)	all
orient. res.	distance_restraints	2	1	type; label; low, up ₁ , up ₂ (nm); weight ()	
angle res.	orientation_restraints	2	1	exp.; label; α ; c (U nm ^{α}); obs. (U); weight (U ⁻¹)	
angle res. z	angle_restraints	4	1	θ_0 (deg); k_c (kJ mol ⁻¹); multiplicity	θ, k
exclusions	angle_restraints_z	2	1	θ_0 (deg); k_c (kJ mol ⁻¹); multiplicity	θ, k
	exclusions	1		one or more atom indices	

'# at' is the number of atom indices

'f. tp' is function type

'F. E.' indicates which parameters can be interpolated during free energy calculations

^(a) the combination rule determines the type of LJ parameters, see 5.3.3

^(b) used by grompp for generating exclusions

^(c) can be converted to constraints by grompp

For free energy calculations, the parameters for topology 'B' ($\lambda = 1$) should be added on the same line, after the normal parameters, in the same order as the normal parameters.

Table 5.4: Intramolecular interaction definitions.

- the topology consists of three levels:
 - the parameter level (see Table 5.3)
 - the molecule level, which should contain one or more molecule definitions (see Table 5.4)
 - the system level: [system], [molecules]
- items should be separated by spaces or tabs, not commas
- atoms in molecules should be numbered consecutively starting at 1
- the file is parsed once only which implies that no forward references can be treated: items must be defined before they can be used
- exclusions can be generated from the bonds or overridden manually
- the bonded force types can be generated from the atom types or overridden per bond
- it is possible to apply multiple bonded interactions of the same type on the same atoms
- descriptive comment lines and empty lines are highly recommended
- starting with GROMACS version 3.1.3 all directives at the parameter level can be used multiple times and there are no restrictions on the order, except that an atom type needs to be defined before it can be used in other parameter definitions
- If parameters for a certain interaction are defined multiple times for the same combination of atom types the last definition is used; starting with GROMACS version 3.1.3 `grompp` generates a warning for parameter redefinitions with different values
- using one of the [atoms], [bonds], [pairs], [angles], etc. without having used [moleculetype] before is meaningless and generates a warning
- using [molecules] without having used [system] before is meaningless and generates a warning.
- after [system] the only allowed directive is [molecules]
- using an unknown string in [] causes all the data until the next directive to be ignored, and generates a warning

Here is an example of a topology file, `urea.top`:

```
;
; Example topology file
;
; The force field files to be included
#include "ffgmx.itp"

[ moleculetype ]
; name nrexcl
```

Urea 3

```
[ atoms ]
; nr type resnr residu atom cgnr charge
1 C 1 UREA C1 1 0.683
2 O 1 UREA O2 1 -0.683
3 NT 1 UREA N3 2 -0.622
4 H 1 UREA H4 2 0.346
5 H 1 UREA H5 2 0.276
6 NT 1 UREA N6 3 -0.622
7 H 1 UREA H7 3 0.346
8 H 1 UREA H8 3 0.276
```

```
[ bonds ]
; ai aj funct b0 kb
3 4 1 1.000000e-01 3.744680e+05
3 5 1 1.000000e-01 3.744680e+05
6 7 1 1.000000e-01 3.744680e+05
6 8 1 1.000000e-01 3.744680e+05
1 2 1 1.230000e-01 5.020800e+05
1 3 1 1.330000e-01 3.765600e+05
1 6 1 1.330000e-01 3.765600e+05
```

```
[ pairs ]
; ai aj funct c6 c12
2 4 1 0.000000e+00 0.000000e+00
2 5 1 0.000000e+00 0.000000e+00
2 7 1 0.000000e+00 0.000000e+00
2 8 1 0.000000e+00 0.000000e+00
3 7 1 0.000000e+00 0.000000e+00
3 8 1 0.000000e+00 0.000000e+00
4 6 1 0.000000e+00 0.000000e+00
5 6 1 0.000000e+00 0.000000e+00
```

```
[ angles ]
; ai aj ak funct th0 cth
1 3 4 1 1.200000e+02 2.928800e+02
1 3 5 1 1.200000e+02 2.928800e+02
4 3 5 1 1.200000e+02 3.347200e+02
1 6 7 1 1.200000e+02 2.928800e+02
1 6 8 1 1.200000e+02 2.928800e+02
7 6 8 1 1.200000e+02 3.347200e+02
2 1 3 1 1.215000e+02 5.020800e+02
2 1 6 1 1.215000e+02 5.020800e+02
3 1 6 1 1.170000e+02 5.020800e+02
```

```

[ dihedrals ]
; ai aj ak al funct phi cp mult
2 1 3 4 1 1.800000e+02 3.347200e+01 2.000000e+00
6 1 3 4 1 1.800000e+02 3.347200e+01 2.000000e+00
2 1 3 5 1 1.800000e+02 3.347200e+01 2.000000e+00
6 1 3 5 1 1.800000e+02 3.347200e+01 2.000000e+00
2 1 6 7 1 1.800000e+02 3.347200e+01 2.000000e+00
3 1 6 7 1 1.800000e+02 3.347200e+01 2.000000e+00
2 1 6 8 1 1.800000e+02 3.347200e+01 2.000000e+00
3 1 6 8 1 1.800000e+02 3.347200e+01 2.000000e+00

[ dihedrals ]
; ai aj ak al funct q0 cq
3 4 5 1 2 0.000000e+00 1.673600e+02
6 7 8 1 2 0.000000e+00 1.673600e+02
1 3 6 2 2 0.000000e+00 1.673600e+02

[ position_restraints ]
; you wouldn't normally use this for a molecule like Urea,
; but we include it here for didactic purposes
; ai funct fc
1 1 1000 1000 1000 ; Restrain to a point
2 1 1000 0 1000 ; Restrain to a line (Y-axis)
3 1 1000 0 0 ; Restrain to a plane (Y-Z-plane)

; Include SPC water topology
#include "spc.itp"

[ system ]
Urea in Water

[ molecules ]
;molecule name nr.
Urea 1
SOL 1000

```

Here follows the explanatory text.

```
[ defaults ]:
```

- non-bond type = 1 (Lennard-Jones) or 2 (Buckingham)
- combination rule =

1. For Lennard Jones: supply $C^{(6)}$ and $C^{(N)}$, $C_{ij}^M = \sqrt{C_i^M C_j^M}$ ($M = 6, N$). Default value for $N = 12$, but it can be overridden using the last parameter on this line. For

Buckingham potentials the combination rule is such that you give the A, B and C parameters. $A_{ij} = \sqrt{A_i A_j}$ and similar for C_{ij} , $B_{ij} = 2/(1/B_i + 1/B_j)$.

2. supply σ and ϵ , $\sigma_{ij} = \frac{1}{2}(\sigma_i + \sigma_j)$ and $\epsilon_{ij} = \sqrt{\epsilon_i \epsilon_j}$

3. supply σ and ϵ , $\sigma_{ij} = \sqrt{\sigma_i \sigma_j}$, $\epsilon_{ij} = \sqrt{\epsilon_i \epsilon_j}$

- generate pairs = no (the default, get 1-4 interactions from the pair list, when parameters are not present in the list give a warning and use zeros) or yes (generate 1-4 interactions which are not present in the pair list from normal Lennard-Jones parameters using FudgeLJ)
- FudgeLJ = factor to change Lennard-Jones 1-4 interactions, default 1
- FudgeQQ = factor to change electrostatic 1-4 interactions, default 1
- N = power for the repulsion term in a 6-N potential (with nonbonded-type Lennard Jones only)

note: generate pairs, FudgeLJ, FudgeQQ and N are optional, FudgeLJ is only used when generate pairs is set to 'yes'. However if you want to specify N you need to give a value for the other parameters as well.

`#include "ffgmx.itp"` : this includes the bonded and non-bonded GROMACS parameters, the `gmx` in `ffgmx` will be replaced by the name of the force field you are actually using.

`[moleculetype]` : defines the name of your molecule in this `*.top` and `nrexcl = 3` stands for excluding non-bonded interactions between atoms that are no further than 3 bonds away.

`[atoms]` : defines the molecule, where `nr` and `type` are fixed, the rest is user defined. So `atom` can be named as you like, `cgnr` made larger or smaller (if possible, the total charge of a charge group should be zero), and charges can be changed here too.

`[bonds]` : no comment.

`[pairs]` : LJ and Coulomb 1-4 interactions

`[angles]` : no comment

`[dihedrals]` : in this case there are 9 proper dihedrals (`funct = 1`), 3 improper (`funct = 2`) and no Ryckaert-Bellemans type dihedrals. If you want to include Ryckaert-Bellemans type dihedrals in a topology, do the following (in case of *e.g.* decane): `[dihedrals]`

```
; ai aj ak al funct c0 c1 c2
```

```
1 2 3 4 3
```

```
2 3 4 5 3
```

and do not forget to *erase the 1-4 interaction* in `[pairs]!!`

`[position_restraints]` : harmonically restrain the selected particles to reference positions (sec. 4.3.1). The reference positions are read from a separate coordinate file by `grompp`.

`#include "spc.itp"` : includes a topology file that was already constructed (see next section, `molecule.itp`).

`[system]` : title of your system, user defined

`[molecules]` : this defines the total number of (sub)molecules in your system that are defined in this `*.top`. In this example file it stands for 1 urea molecules dissolved in 1000 water molecules. The molecule type `SOL` is defined in the `spc.itp` file.

5.6.2 Molecule.itp file

If you construct a topology file you will use frequently (like a water molecule, `spc.itp`) it is better to make a `molecule.itp` file, which only lists the information of the molecule:

```
[ moleculetype ]
; name nrexcl
Urea 3

[ atoms ]
; nr type resnr residu atom cgnr charge
1 C 1 UREA C1 1 0.683
.....
.....
8 H 1 UREA H8 3 0.276

[ bonds ]
; ai aj funct c0 c1
3 4 1 1.000000e-01 3.744680e+05
.....
.....
1 6 1 1.330000e-01 3.765600e+05

[ pairs ]
; ai aj funct c0 c1
2 4 1 0.000000e+00 0.000000e+00
.....
.....
5 6 1 0.000000e+00 0.000000e+00

[ angles ]
; ai aj ak funct c0 c1
1 3 4 1 1.200000e+02 2.928800e+02
.....
.....
3 1 6 1 1.170000e+02 5.020800e+02

[ dihedrals ]
; ai aj ak al funct c0 c1 c2
2 1 3 4 1 1.800000e+02 3.347200e+01 2.000000e+00
.....
.....
3 1 6 8 1 1.800000e+02 3.347200e+01 2.000000e+00

[ dihedrals ]
; ai aj ak al funct c0 c1
```

```

3 4 5 1 2 0.000000e+00 1.673600e+02
6 7 8 1 2 0.000000e+00 1.673600e+02
1 3 6 2 2 0.000000e+00 1.673600e+02

```

This results in a very short `*.top` file as described in the previous section, but this time you only need to include files:

```

; The force field files to be included
#include "ffgm.itp"

; Include urea topology
#include "urea.itp"

; Include SPC water topology
#include "spc.itp"

[ system ]
Urea in Water

[ molecules ]
;molecule name number
Urea 1
SOL 1000

```

5.6.3 Ifdef option

A very powerful feature in GROMACS is the use of `#ifdef` statements in your `*.top` file. By making use of this statement, different parameters for one molecule can be used in the same `*.top` file. An example is given for TFE, where there is an option to use different charges on the atoms: charges derived by De Loof *et al.* [78] or by Van Buuren and Berendsen [49]. In fact you can use all the options of the C-Preprocessor, `cpp`, because this is used to scan the file. The way to make use of the `#ifdef` option is as follows:

- in `grompp.mdp` (the GROMACS preprocessor input parameters) use the option
`define = -DDeLoof`
or
`define =`
- put the `#ifdef` statements in your `*.top`, as shown below:

...

```

[ atoms ]
; nr type resnr residu atom cgnr charge mass
#ifdef DeLoof

```

```

; Use Charges from DeLoof
1 C 1 TFE C 1 0.74
2 F 1 TFE F 1 -0.25
3 F 1 TFE F 1 -0.25
4 F 1 TFE F 1 -0.25
5 CH2 1 TFE CH2 1 0.25
6 OA 1 TFE OA 1 -0.65
7 HO 1 TFE HO 1 0.41
#else
; Use Charges from VanBuuren
1 C 1 TFE C 1 0.59
2 F 1 TFE F 1 -0.2
3 F 1 TFE F 1 -0.2
4 F 1 TFE F 1 -0.2
5 CH2 1 TFE CH2 1 0.26
6 OA 1 TFE OA 1 -0.55
7 HO 1 TFE HO 1 0.3
#endif

[ bonds ]
; ai aj funct c0 c1
6 7 1 1.000000e-01 3.138000e+05
1 2 1 1.360000e-01 4.184000e+05
1 3 1 1.360000e-01 4.184000e+05
1 4 1 1.360000e-01 4.184000e+05
1 5 1 1.530000e-01 3.347000e+05
5 6 1 1.430000e-01 3.347000e+05

...

```

5.6.4 Free energy calculations

Free energy differences between two systems A and B can be calculated as described in sec. 3.12. The systems A and B are described by topologies consisting of the same number of molecules with the same number of atoms. Masses and non-bonded interactions can be perturbed by adding B parameters in the [atoms] field. Bonded interactions can be perturbed by adding B parameters to the bonded types or the bonded interactions. The parameters that can be perturbed are listed in Table 5.3 and Table 5.4. The λ -dependence of the interactions is described in section sec. 4.5. Below is an example of a topology which changes from 200 propanols to 200 pentanes using the GROMOS-96 force field.

```

; Include forcefield parameters
#include "ffG43a1.itp"

[ moleculetype ]

```



```
; Name nrexcl
PropPent 3

[ atoms ]
; nr type resnr residue atom cgnr charge mass typeB chargeB massB
1 H 1 PROP PH 1 0.398 1.008 CH3 0.0 15.035
2 OA 1 PROP PO 1 -0.548 15.9994 CH2 0.0 14.027
3 CH2 1 PROP PC1 1 0.150 14.027 CH2 0.0 14.027
4 CH2 1 PROP PC2 2 0.000 14.027
5 CH3 1 PROP PC3 2 0.000 15.035

[ bonds ]
; ai aj funct par_A par_B
1 2 2 gb_1 gb_26
2 3 2 gb_17 gb_26
3 4 2 gb_26 gb_26
4 5 2 gb_26

[ pairs ]
; ai aj funct
1 4 1
2 5 1

[ angles ]
; ai aj ak funct par_A par_B
1 2 3 2 ga_11 ga_14
2 3 4 2 ga_14 ga_14
3 4 5 2 ga_14 ga_14

[ dihedrals ]
; ai aj ak al funct par_A par_B
1 2 3 4 1 gd_12 gd_17
2 3 4 5 1 gd_17 gd_17

[ system ]
; Name
Propanol to Pentane
4
[ molecules ]
; Compound #mols
PropPent 200
```

Atoms that are not perturbed, PC2 and PC3, do not need B parameter specifications, the B parameters will be copied from the A parameters. Bonded interactions between atoms that are not perturbed do not need B parameter specifications, here this is the case for the last bond. Topolo-

gies using the GROMACS force field need no bonded parameters at all, since both the A and B parameters are determined by the atom types. Non-bonded interactions involving one or two perturbed atoms use the free-energy perturbation functional forms. Non-bonded interaction between two non-perturbed atoms use the normal functional forms. This means that when, for instance, only the charge of a particle is perturbed, its Lennard-Jones interactions will also be affected when λ is not equal to zero or one.

Note that this topology uses the GROMOS-96 force field, in which the bonded interactions are not determined by the atom types. The bonded interaction strings are converted by the C-preprocessor. The force field parameter files contain lines like:

```
#define gb_26 0.1530 7.1500e+06
#define gd_17 0.000 5.86 3
```

5.6.5 Constraint force

The constraint force between two atoms in one molecule can be calculated with the free energy perturbation code by adding a constraint between the two atoms, with a different length in the A and B topology. When the B length is 1 nanometer longer than the A length and λ is kept constant at zero, the derivative of the Hamiltonian with respect to λ is the constraint force. For constraints between molecules the pull code can be used, see sec. 6.2. Below is an example for calculating the constraint force at 0.7 nanometer between two methanes in water, by combining the two methanes into one molecule. The added constraint is of function type 2, which means that it is not used for generating exclusions (see 5.3.5).

```
; Include forcefield parameters
#include "ffG43a1.itp"

[ moleculetype ]
; Name nrexcl
Methanes 1

[ atoms ]
; nr type resnr residu atom cgnr charge mass
1 CH4 1 CH4 C1 1 0 16.043
2 CH4 1 CH4 C2 2 0 16.043

[ constraints ]
; ai aj funct length_A length_B
1 2 2 0.7 1.7

#include "spc.itp"

[ system ]
; Name
```

Methanes in Water

```
[ molecules ]
; Compound #mols
Methanes 1
SOL 2002
```

5.6.6 Coordinate file

Files with the `.gro` file extension contain a molecular structure in GROMOS87 format. A sample piece is included below:

```
MD of 2 waters, reformat step, PA aug-91
6
1WATER OW1 1 0.126 1.624 1.679 0.1227 -0.0580 0.0434
1WATER HW2 2 0.190 1.661 1.747 0.8085 0.3191 -0.7791
1WATER HW3 3 0.177 1.568 1.613 -0.9045 -2.6469 1.3180
2WATER OW1 4 1.275 0.053 0.622 0.2519 0.3140 -0.1734
2WATER HW2 5 1.337 0.002 0.680 -1.0641 -1.1349 0.0257
2WATER HW3 6 1.326 0.120 0.568 1.9427 -0.8216 -0.0244
1.82060 1.82060 1.82060
```

This format is fixed, *i.e.* all columns are in a fixed position. If you want to read such a file in your own program without using the GROMACS libraries you can use the following formats:

C-format: "%5i%5s%5s%5i%8.3f%8.3f%8.3f%8.4f%8.4f%8.4f"

Or to be more precise, with title *etc.* it looks like this:

```
"for (i=0; (i<natoms); i++)
"residuenr, residuename, atomname, atomnr, x, y, z, vx, vy, vz

"box[X] [X], box[Y] [Y], box[Z] [Z],
box[X] [Y], box[X] [Z], box[Y] [X], box[Y] [Z], box[Z] [X], box[Z] [Y]
```

Fortran format: (i5, 2a5, i5, 3f8.3, 3f8.4)

So `conf.in.gro` is the GROMACS coordinate file and is almost the same as the GROMOS-87 file (for GROMOS users: when used with `ntx=7`). The only difference is the box for which GROMACS uses a tensor, not a vector.

5.7 Force-field organization

5.7.1 Force-field files

GROMACS 3.3 includes five forcefields. They are listed the file `FF.dat`:

```

5
ffgmX Gromacs Forcefield (see manual)
ffgmX2 Gromacs Forcefield with all hydrogens (proteins only)
ffG43a1 GROMOS96 43a1 Forcefield (official distribution)
ffG43b1 GROMOS96 43b1 Vacuum Forcefield (official distribution)
ffG43a2 GROMOS96 43a2 Forcefield (development) (improved ...)

```

All files for each force field have names beginning with the `ff???` string in the `FF.dat` file. A force field is included at the beginning of a topology file with an `#include` statement followed by `ff???.itp`. This statement includes the force-field file, which in turn may include other forcefield files. The five force fields are organized in the same way. As an example we show the `ffgmX.itp` force-field file:

```

#define _FF_GROMACS
#define _FF_GROMACS1

[ defaults ]
; nbfunc comb-rule gen-pairs fudgeLJ fudgeQQ
1 1 no 1.0 1.0

#include "ffgmXnb.itp"
#include "ffgmXbon.itp"

```

The first `#define` can be used in topologies to parse data which is specific for all GROMACS force-fields, the second `#define` to parse data which is specific for this force field. The `defaults` section is explained in 5.6.1. The included file `ffgmXnb.itp` contains all atom types and non-bonded parameters. The included file `ffgmXbon.itp` contains all bonded parameters.

For each force field there are five files which are only used by `pdb2gmX`. These are: the residue database (`.rtp`, see 5.5.1) the hydrogen database (`.hdb`, see 5.5.2), two termini databases (`.tdb`, see 5.5.3) and the atom type database (`.atp`) which contains only the masses.

5.7.2 Changing force-field parameters

If one wants to change the parameters of few bonded interactions in a molecule, this is most easily accomplished by typing the parameters behind the definition of the bonded interaction in the `[moleculetype]` section (see 5.6.1 for the format and units). If one wants to change the parameters for all instances of a certain interaction one can change them in the force-field file or add a new `[???types]` section after including the force field. When parameters for a certain interaction are defined multiple times the last definition is used. As of GROMACS version 3.1.3 a warning is generated when parameters are redefined with a different value. Changing the Lennard-Jones parameters of an atom type is not recommended, because in the GROMACS and GROMOS force-fields the Lennard-Jones parameters for several combinations of atom types are not according to the standard combination rules. Such combinations (and possibly also combinations that do follow the combination rules) are defined in the `[nonbonded_params]` section and changing the Lennard-Jones parameters of an atom type has no effect on these combinations.

5.7.3 Adding atom types

As of GROMACS version 3.1.3 atom types can be added in an extra [`atomtypes`] section after the the including of the normal forcefield. After the definition of the new atom type(s), additional non-bonded and pair parameters can be defined. In pre 3.1.3 versions of GROMACS the new atom types needed to be added in the [`atomtypes`] section of the forcefield files, because all non-bonded parameters above the last [`atomtypes`] section would be overwritten using the standard combination rules.

Chapter 6

Special Topics

6.1 Potential of mean force

A potential of mean force (pmf) is a potential which is obtained by integrating the mean force from an ensemble of configurations. In GROMACS there are several different methods to calculate the mean force. Each method has its limitations, which are listed below.

- **pull code:** between the centers of mass of molecules or groups of molecules.
- **free-energy code with harmonic bonds or constraints:** between single atoms.
- **free-energy code with position restraints:** changing the conformation of a relatively immobile group of atoms.
- **not possible currently:** between groups of atoms that are part of a larger molecule for which the bonds are constrained with SHAKE or LINCS.

The pull and free-energy code are described in more detail in the following two sections.

6.1.1 Entropic effects

When a distance between two atoms or the centers of mass of two groups is constrained or restrained, there will be a purely entropic contribution to the pmf due to the rotation of the two groups. For a system of two non-interacting masses the potential of mean force is:

$$V_{pmf}(r) = -(n_c - 1)k_B T \log(r) \quad (6.1)$$

where n_c is the number of dimensions in which the constraint works (i.e. $n_c = 3$ for a normal constraint and $n_c = 1$ when only the z -direction is constrained). Whether one needs to correct for this contribution depends on what the pmf should represent. When one wants to pull a substrate into a protein, this entropic term indeed contributes to the work to get the substrate into the protein. But when calculating a pmf between two solutes in a solvent, for the purpose of simulating without solvent, the entropic contribution should be removed. Note that this term can be significant; when at 300K the distance is halved the contribution is 3.5 kJ mol^{-1} .

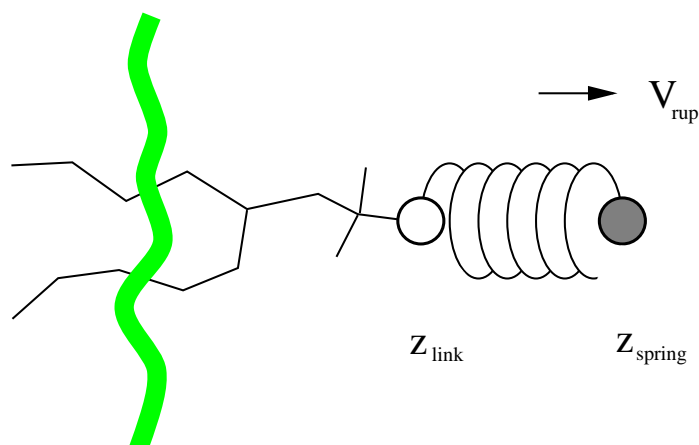


Figure 6.1: Schematic picture of pulling a lipid out of a lipid bilayer with AFM pulling. V_{rup} is the velocity at which the spring is retracted, Z_{link} is the atom to which the spring is attached and Z_{spring} is the location of the spring.

6.2 Calculating the PMF using the pull code

There are a number of options to calculate potentials of mean force and related topics. In the current version of GROMACS this is implemented through some extra files for `mdrun`.

6.2.1 Overview

Three different types of calculation are supported:

1. **AFM pulling** A spring is connected to an atom and slowly retracted. This has the effect of pulling an atom or group of atoms away from its initial location. The rate constant and spring constant for the spring can be varied to study e.g. the unbinding of a protein and a ligand (see figure 6.1).
2. **Constraint forces** The distance between the centers of mass of two groups of atoms can be constrained and the constraint force monitored. The distance can be in 1, 2, or 3 dimensions. This method uses the SHAKE algorithm but only needs 1 iteration to be exact if only two groups are constrained.
3. **Umbrella sampling** A simple umbrella sampling with an harmonic umbrella potential that acts on the center of mass of a group of atoms.

In the calculations, there has to be 1 reference group and 1 to 20 other groups of atoms. For constrained runs, the distance between the reference group and the other groups is kept constant at the distance they have in the input coordinate file (`.tpr`) file.

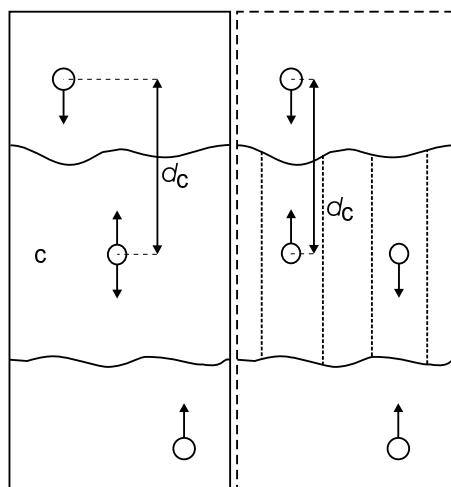


Figure 6.2: Overview of the different reference group possibilities, applied to interface systems. C is the reference group. The circles represent the center of mass of 2 groups plus the reference group, and d_c is the reference distance.

6.2.2 Usage

Input files

The `mdrun` programs needs 4 additional files: 2 input files and 2 output files.

`-pi pull.ppa`

If this file is specified the pull code will be used. It contains the parameters that control what type of calculation is done. A full explanation of all the options is given below.

`-pn index.ndx`

This file defines the different groups for use in all pull calculations. The groups are referred to by name, so the index file can contain other groups that are not used as well.

`-po pullout.ppa`

A formatted copy of the input parameter file with the parameters that were actually used in the run.

`-pdo pull.pdo`

The data file with the calculated forces (AFM pulling, constraint force) or positions (umbrella sampling).

Definition of groups

The way the reference groups and different reference types work is summarized in figure 6.2. There are four different possibilities for the reference group.

com

The center of mass of the group given under `reference_group`, calculated each step from the current coordinates.

`com_t0`

The center of mass of the group given under `reference_group`, calculated each step from the current coordinates, but corrected for atoms that have crossed the box. If the reference group consists of all the water molecules in the system, and a single water molecule moves across the box and enters from the other side, the c.o.m. will show a slight jump. This is simply due to the periodic boundary conditions, and shows that the center of mass in a simulation in periodic boundary conditions is ill defined if the group used to calculate it is *e.g.* a slab of liquid. If the 'real' positions are used instead of the coordinates that have been reset to be inside the box, the center of mass of the **whole** system is conserved.

`dynamic`

In a phospholipid bilayer system it may be of interest to calculate the pmf of a lipid as function of its distance from the whole bilayer. The whole bilayer can be taken as reference group in that case, but it might also be of interest to define the reaction coordinate for the pmf more locally. `dynamic` does not use all the atoms of the `reference_group`, but instead only those within a cylinder with radius `r` below the main group. This only works for distances defined in 1 dimension, and the cylinder is oriented with its long axis along this 1 dimension. A second cylinder can be defined with `rc`, with a linear switch function that weighs the contribution of atoms between `r` and `rc` with distance. This smoothes the effects of atoms moving in and out of the cylinder (which causes jumps in the constraint forces).

`dynamic_t0`

The same as `dynamic`, but the coordinates are corrected for boxcrossings like in `com_t0`. Note that strictly speaking this is not correct if the reference group is not the whole system, including the groups defined with `group_1, ...`

To further smooth rapidly fluctuating distances between the reference group and the other groups, the average distance can be constrained instead of the instantaneous distance. This is defined by setting `reflag` to the number of steps to average over. However, using this option is not strictly correct for calculating potentials of mean force from the average constraint force.

When relative weights w_i are used during the calculations, either due to a dynamic reference group or by supplying weights in the input, the weights need to be scaled to conserve momentum:

$$w'_i = w_i \frac{\sum_{j=1}^N w_j m_j}{\sum_{j=1}^N w_j^2 m_j} \quad (6.2)$$

where m_j is the mass of atom j of the group. The mass of the group, required for calculating the constraint force, is:

$$M = \sum_{i=1}^N w'_i m_i \quad (6.3)$$

The definition of the weighted center of mass is:

$$\mathbf{r}_{com} = \frac{\sum_{i=1}^N w'_i m_i \mathbf{r}_i}{M} \quad (6.4)$$

From the centers of mass the AFM, constraint or umbrella force F_{com} on each group can be calculated. The force on the center of mass of a group is redistributed to the atoms as follows:

$$\mathbf{F}_i = \frac{w'_i m_i}{M} \mathbf{F}_{com} \quad (6.5)$$

6.2.3 The parameter file

General

`verbose = no`

If this is set to `yes`, a large amount of detailed information is sent to `stderr`, which is only useful for diagnostic purposes. The `.pdo` file also becomes more detailed, which is not necessary for normal use.

`runtype = constraint`

Options are `afm`, `constraint`, `umbrella`. This selects the type of calculation: AFM pulling, constraint force calculation or umbrella sampling.

`ngroups = 2`

`group_1 = MB21_1`

`group_2 = MB21_2`

The groups with the atoms to act on. `ngroups` should be between 1 and 20.

`reference_group = OCTA`

The reference group. Distances are calculated between `group_1` (and `group_2` if specified) and this group. If *e.g.* the constraint force between two ions is needed, you would specify `group_1` as a group with 1 ion, and `reference_group` as the other ion. If the reference group is left blank, then AFM and umbrella sampling can be performed in absolute coordinates.

`weights_1 =`

`weights_2 =`

`reference_weights =`

The relative weights for all atoms in each group. The center of mass and the force calculation are performed using the masses of the atoms multiplied by the weights. The weights are scaled before use, thus only the ratios in the input matter. When no weights are given all weights are set to one. Weights can not be used with dynamic reference groups.

`reftype = com`

The type of reference group. Options are `com`, `com_t0`, `dynamic`, `dynamic_t0` as explained above.

`reflag = 1`

The position of the reference group can be taken as average over a number of steps, specified by `reflag` (see above).

`pulldim = N N Y`

Selects the dimensions used for the pulling. In this example the distance is calculated using only the z-component and pulling is performed only in the z-direction. If nothing is given all three dimensions are used.

Dynamic reference group options

`r = 0`

If dynamic reference groups are selected (`dynamic`, `dynamic_t0`), `r` is the radius of the cylinder used to define which atoms are part of the reference group (see above).

`rc = 0`

With dynamic reference groups, the cylinder can be smoothly switched so that atoms that fall between `r` and `rc` are weighted linearly from 1 to 0 going from `r` to `rc`. As reasonable initial values we suggest `r = 1.0` and `rc = 1.5`, but this will depend strongly on the exact system of interest.

`update = 1`

The frequency with which the dynamic reference groups are recalculated. Usually there is no reason to use anything other than 1.

Constraint run options

`constraint_direction =`

The direction for constraining. The default is 0.0 0.0 0.0, which means no direction and thus the distance between groups is constrained. The direction is normalized before use.

`constraint_distance1 =`

`constraint_distance2 =`

The initial constraint distance in nm, is 0 by default, which means that the distance is determined by the starting coordinates.

`constraint_rate1 =`

`constraint_rate2 =`

The rate of the change of the constraint distance in nm/ps, is 0 by default.

`constraint_tolerance =`

The tolerance for the constraint distance, 1E-6 nm is the default.

AFM options

AFM pulling attaches a spring to the center of mass of the pulled group. This spring moves in relation to the center of mass of the reference group. If the reference group moves, then the spring does as well. If the `reference_group` is left blank, then pulling will be in absolute coordinates, with the reference group position taken to be [0 0 0].

afm_rate1 =

afm_rate2 =

The rate at which the spring moves in nm/ps for each group.

afm_k1 =

afm_k2 =

The force constant of the spring for each pulled group in $\text{kJ mol}^{-1} \text{nm}^{-2}$.

afm_dir1 =

afm_dir2 =

Unit vector describing the direction of pulling.

afm_init1 =

afm_init2 =

Vector describing the initial position of the spring relative to the reference group. To start a simulation with zero initial force on the pulled group, the initial position should be set to the position of the pulled group relative to the reference group.

Umbrella Sampling options

Umbrella sampling applies a simple harmonic force maintain the position of the pulled group relative to the reference group. If the reference group moves, then the pulled group will move as well. If the `reference_group` is left blank, then pulling will be in absolute coordinates, with the reference group position taken to be [0 0 0].

k1 =

k2 =

This option specifies the force constant for each pulled group in $\text{kJ mol}^{-1} \text{nm}^{-2}$.

pos1 =

pos2 =

This option specifies the position that the pulled group will be restrained to, relative to the reference group.

6.2.4 Output

Header

The pull code begins with a header file. All lines in the header file are prepended with a # symbol.

```
# CONSTRAINT 3.0
```

```
# AFM 3.0
```

```
# UMBRELLA 3.0
```

The first line indicates the type of run and the version of the output file. Future versions of gromacs may change the output format and the version number will be updated to reflect this.

```
# Component selection: 0.0 0.0 1.0
```

The second line shows the value of `pull_dim` used in the simulation.

```
# Skip 1
```

The third line shows the output frequency.

```
# Ref. Group 'ref'
```

The fourth line shows the name of the reference group

```
# Nr. of pull groups 2
```

The fifth line shows the number of pulled groups.

```
# Group 1 grp_1 0.0 0.0 1.0 0.01 500
```

For AFM, the sixth and further lines show the number of each pulled group, its name, the x, y, and z components the direction vector, the pull rate, and the force constant.

```
# Group 1 grp_1 1.2 500
```

For umbrella sampling, the sixth and further lines show the number of each pulled group, its name, the position relative the pulled group, and the force constant. Note that only the dimensions for which `pull_dim` is Y are output.

```
# Group 1 grp_1 1.0
```

For constraint runs, the sixth and further lines show the number of each pulled group, its name, and the initial position of the pulled group. Note that only the dimensions for which `pull_dim` is Y are output.

Constraint

```
10.5 -81.2
```

The output file consists of the time and the force the constraint experiences for each group. A negative number means that the reference group and the pulled group attract each other.

AFM

```
10.5 3.5 3.63 3.6
```

The output file consists of the time, the position of the reference group, and the position of the pulled group, and the position of the spring. The position of the pulled group and spring are in absolute coordinates. Subtract the reference coordinate to get the position relative to the pulled group. Note that only the dimensions for which `pull_dim` is Y are output.

Umbrella Sampling

10.5 0.002

The output file consists of the time, and the deviation of each pulled group from its restrained position. Note that only the dimensions for which `pulldim` is `Y` are output.

6.2.5 Limitations

Apart from obvious limitations that are simply not implemented (*e.g.* a better umbrella sampling and analysis scheme), there is one important limitation: constraint forces can **only** be calculated between groups that are not connected by constraints to the rest of the system. If a group contains part of a molecule of which the bondlengths are constrained, SHAKE or LINCS and the constraint force calculation here will interfere with each other, making the results unreliable. This means that for simulations with `constraints = all-bonds` in the `.mdp` file pulling is limited to whole molecules or groups of molecules. In some cases this limitation can be avoided by using the free energy code, see sec. 6.3.

Another limitation is connected with pressure scaling. When pressure scaling is used, the shortest distance between the pull groups and the reference group should be within in the central unit-cell, not to periodic images of the pull groups. Otherwise the pressure scaling will affect the reference distances and thus also the forces. The exception is constraint pulling with `constraint_distance? > 0`, this will work correctly with pressure scaling under all conditions.

6.2.6 Implementation

The code for the options described above can be found in the files `pull.c`, `pullinit.c`, `pullio.c`, `pullutil.c` and the headerfiles `pull.h` and `pulls.h`. This last file defines a few datatypes, `pull.h` explains the main functions.

6.2.7 Future development

There are several additional features that would be useful, including more advanced umbrella sampling, an analysis tool to analyse the output of the pull code, incorporation of the input parameters and index file into the `grompp` program input files, extension to more groups, more flexible definition of a reaction coordinate, extension to groups that are parts of molecules that use SHAKE or LINCS, and a combination of the starting structure calculation with constraints for faster convergence of starting structures.

6.3 Calculating the PMF using the free-energy code

The free-energy coupling-parameter approach (see sec. 3.12) provides several ways to calculate potentials of mean force. A potential of mean force between two atoms can be calculated by connecting them with a harmonic potential or a constraint (for this purpose there a special potentials

that avoid the generation of extra exclusions, see 5.3.5). When the position of the minimum or the constraint length is 1 nm more in state B than in state A, the restraint or constraint force is given by $\partial H/\partial\lambda$. The distance between the atoms can be changed as a function of λ and time by setting `delta-lambda` in the `.mdp` file. The results should be identical (although not numerically due to the different implementations) to the results of the pull code with umbrella sampling and constraint pulling. Unlike the pull code, the free energy code can also handle atoms that are connected by constraints.

Potentials of mean force can also be calculated using position restraints. With position restraints atoms can be linked to a position in space with a harmonic potential (see sec. 4.3.1). These positions can be made a function of the coupling parameter λ . The positions for the A and the B state are supplied to `grompp` with the `-r` and `-rb` option, respectively. One could use this approach to do targeted MD; note that we do not encourage the use of targeted MD for proteins. A protein can be forced from one conformation to another by using these conformations as position restraint coordinates for state A and B. One can then slowly change λ from 0 to 1. The main drawback of this approach is that the conformational freedom of the protein is severely limited by the position restraints, independent of the change from state A to B. Also the protein is forced from state A to B in an almost straight line, whereas the real pathway might be very different. An example of a more fruitful application is a solid system or a liquid confined between walls where one wants to measure the force required to change the separation between the boundaries or walls. Because the boundaries or walls already need to be fixed, the position restraints do not limit the system in its sampling.

6.4 Removing fastest degrees of freedom

The maximum time step in MD simulations is limited by the smallest oscillation period that can be found in the simulated system. Bond-stretching vibrations are in their quantum-mechanical ground state and are therefore better represented by a constraint than by a harmonic potential.

For the remaining degrees of freedom, the shortest oscillation period as measured from a simulation is 13 fs for bond-angle vibrations involving hydrogen atoms. Taking as a guideline that with a Verlet (leap-frog) integration scheme a minimum of 5 numerical integration steps should be performed per period of a harmonic oscillation in order to integrate it with reasonable accuracy, the maximum time step will be about 3 fs. Disregarding these very fast oscillations of period 13 fs the next shortest periods are around 20 fs, which will allow a maximum time step of about 4 fs

Removing the bond-angle degrees of freedom from hydrogen atoms can best be done by defining them as virtual interaction-sites instead of normal atoms. Where a normal atom is connected to the molecule with bonds, angles and dihedrals, a virtual site's position is calculated from the position of three nearby heavy atoms in a predefined manner (see also sec. 4.7). For the hydrogens in water and in hydroxyl, sulfhydryl or amine groups, no degrees of freedom can be removed, because rotational freedom should be preserved. The only other option available to slow down these motions, is to increase the mass of the hydrogen atoms at the expense of the mass of the connected heavy atom. This will increase the moment of inertia of the water molecules and the hydroxyl, sulfhydryl or amine groups, without affecting the equilibrium properties of the system and without affecting the dynamical properties too much. These constructions will shortly be described in sec. 6.4.1 and have previously been described in full detail [79].

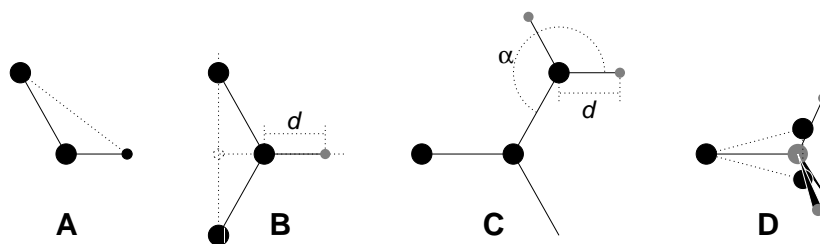


Figure 6.3: The different types of virtual site constructions used for hydrogen atoms. The atoms used in the construction of the virtual site(s) are depicted as black circles, virtual sites as grey ones. Hydrogens are smaller than heavy atoms. **A**: fixed bond angle, note that here the hydrogen is not a virtual site; **B**: in the plane of three atoms, with fixed distance; **C**: in the plane of three atoms, with fixed angle and distance; **D**: construction for amine groups ($-\text{NH}_2$ or $-\text{NH}_3^+$), see text for details.

Using both virtual sites and modified masses, the next bottleneck is likely to be formed by the improper dihedrals (which are used to preserve planarity or chirality of molecular groups) and the peptide dihedrals. The peptide dihedral cannot be changed without affecting the physical behavior of the protein. The improper dihedrals that preserve planarity, mostly deal with aromatic residues. Bonds, angles and dihedrals in these residues can also be replaced with somewhat elaborate virtual site constructions.

All modifications described in this section can be performed using the GROMACS topology building tool `pdb2gmx`. Separate options exist to increase hydrogen masses, virtualize all hydrogen atoms or also virtualize all aromatic residues. Note that when all hydrogen atoms are virtualized, also those inside the aromatic residues will be virtualized, *i.e.* hydrogens in the aromatic residues are treated differently depending on the treatment of the aromatic residues.

Parameters for the virtual site constructions for the hydrogen atoms are inferred from the forcefield parameters (*vis.* bond lengths and angles) directly by `grompp` while processing the topology file. The constructions for the aromatic residues are based on the bond lengths and angles for the geometry as described in the forcefields, but these parameters are hard-coded into `pdb2gmx` due to the complex nature of the construction needed for a whole aromatic group.

6.4.1 Hydrogen bond-angle vibrations

Construction of virtual sites

The goal of defining hydrogen atoms as virtual sites is to remove all high-frequency degrees of freedom from them. In some cases not all degrees of freedom of a hydrogen atom should be removed, *e.g.* in the case of hydroxyl or amine groups the rotational freedom of the hydrogen atom(s) should be preserved. Care should be taken that no unwanted correlations are introduced by the construction of virtual sites, *e.g.* bond-angle vibration between the constructing atoms could translate into hydrogen bond-length vibration. Additionally, since virtual sites are by definition massless, in order to preserve total system mass, the mass of each hydrogen atom that is treated as virtual site should be added to the bonded heavy atom.

Taking into account these considerations, the hydrogen atoms in a protein naturally fall into several

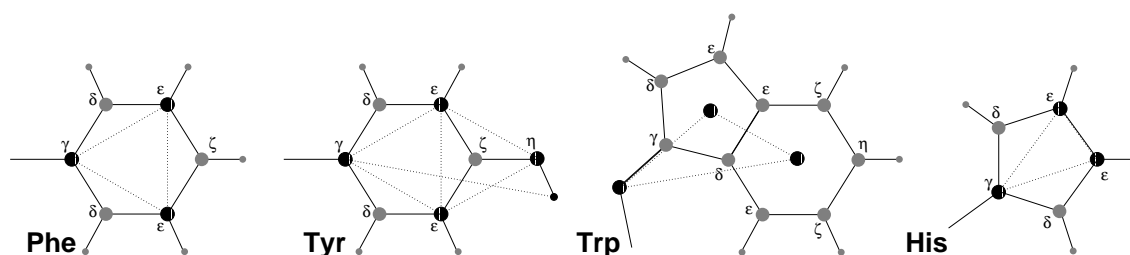


Figure 6.4: The different types of virtual site constructions used for aromatic residues. The atoms used in the construction of the virtual site(s) are depicted as black circles, virtual sites as grey ones. Hydrogens are smaller than heavy atoms. A: phenylalanine; B: tyrosine (note that the hydroxyl hydrogen is *not* a virtual site); C: tryptophane; D: histidine.

categories, each requiring a different approach (see also Fig. 6.3).

- *hydroxyl (-OH) or sulfhydryl (-SH) hydrogen*: The only internal degree of freedom in a hydroxyl group that can be constrained is the bending of the C-O-H angle. This angle is fixed by defining an additional bond of appropriate length, see Fig. 6.3A. This removes the high frequency angle bending, but leaves the dihedral rotational freedom. The same goes for a sulfhydryl group. Note that in these cases the hydrogen is not treated as a virtual site.
- *single amine or amide (-NH-) and aromatic hydrogens (-CH-)*: The position of these hydrogens cannot be constructed from a linear combination of bond vectors, because of the flexibility of the angle between the heavy atoms. Instead, the hydrogen atom is positioned at a fixed distance from the bonded heavy atom on a line going through the bonded heavy atom and a point on the line through both second bonded atoms, see Fig. 6.3B.
- *planar amine (-NH₂) hydrogens*: The method used for the single amide hydrogen is not well suited for planar amine groups, because no suitable two heavy atoms can be found to define the direction of the hydrogen atoms. Instead, the hydrogen is constructed at a fixed distance from the nitrogen atom, with a fixed angle to the carbon atom, in the plane defined by one of the other heavy atoms, see Fig. 6.3C.
- *amine group (umbrella -NH₂ or -NH₃⁺) hydrogens*: Amine hydrogens with rotational freedom cannot be constructed as virtual sites from the heavy atoms they are connected to, since this would result in loss of the rotational freedom of the amine group. To preserve the rotational freedom while removing the hydrogen bond-angle degrees of freedom, two “dummy masses” are constructed with the same total mass, moment of inertia (for rotation around the C-N bond) and center of mass as the amine group. These dummy masses have no interaction with any other atom, except for the fact that they are connected to the carbon and to each other, resulting in a rigid triangle. From these three particles the positions of the nitrogen and hydrogen atoms are constructed as linear combinations of the two carbon-mass vectors and their outer product, resulting in an amine group with rotational freedom intact, but without other internal degrees of freedom. See Fig. 6.3D.

6.4.2 Out-of-plane vibrations in aromatic groups

The planar arrangements in the side chains of the aromatic residues lends itself perfectly to a virtual-site construction, giving a perfectly planar group without the inherently instable constraints that are necessary to keep normal atoms in a plane. The basic approach is to define three atoms or dummy masses with constraints between them to fix the geometry and create the rest of the atoms as simple virtual sites type (see sec. 4.7) from these three. Each of the aromatic residues require a different approach:

- *Phenylalanine*: C_γ , $C_{\epsilon 1}$ and $C_{\epsilon 2}$ are kept as normal atoms, but with each a mass of one third the total mass of the phenyl group. See Fig. 6.3A.
- *Tyrosine*: The ring is treated identical to the phenylalanine ring. Additionally, constraints are defined between $C_{\epsilon 1}$ and $C_{\epsilon 2}$ and O_η . The original improper dihedral angles will keep both triangles (one for the ring and one with O_η) in a plane, but due to the larger moments of inertia this construction will be much more stable. The bond angle in the hydroxyl group will be constrained by a constraint between C_γ and H_η , note that the hydrogen is not treated as a virtual site. See Fig. 6.3B.
- *Tryptophane*: C_β is kept as a normal atom and two dummy masses are created at the center of mass of each of the rings, each with a mass equal to the total mass of the respective ring ($C_{\delta 2}$ and $C_{\epsilon 2}$ are each counted half for each ring). This keeps the overall center of mass and the moment of inertia almost (but not quite) equal to what it was. See Fig. 6.3C.
- *Histidine*: C_γ , $C_{\epsilon 1}$ and $N_{\epsilon 2}$ are kept as normal atoms, but with masses redistributed such that the center of mass of the ring is preserved. See Fig. 6.3D.

6.5 Viscosity calculation

The shear viscosity is a property of liquid which can be determined easily by experiment. It is useful for parameterizing the forcefield, because it is a kinetic property, while most other properties which are used for parameterization are thermodynamic. The viscosity is also an important property, since it influences the rates of conformational changes of molecules solvated in the liquid.

The viscosity can be calculated from an equilibrium simulation using an Einstein relation:

$$\eta = \frac{1}{2} \frac{V}{k_B T} \lim_{t \rightarrow \infty} \frac{d}{dt} \left\langle \left(\int_{t_0}^{t_0+t} P_{xz}(t') dt' \right)^2 \right\rangle_{t_0} \quad (6.6)$$

This can be done with `g_energy`. This method converges very slowly [80]. A nanosecond simulation might not be long enough for an accurate determination of the viscosity. The result is very dependent on the treatment of the electrostatics. Using a (short) cut-off results in large noise on the off-diagonal pressure elements, which can increase the calculated viscosity by an order of magnitude.

GROMACS also has a non-equilibrium method for determining the viscosity [80]. This makes use of the fact that energy, which is fed into system by external forces, is dissipated through viscous

friction. The generated heat is removed by coupling to a heat bath. For a Newtonian liquid adding a small force will result in a velocity gradient according to the following equation:

$$a_x(z) + \frac{\eta}{\rho} \frac{\partial^2 v_x(z)}{\partial z^2} = 0 \quad (6.7)$$

here we have applied an acceleration $a_x(z)$ in the x -direction, which is a function of the z -coordinate. In GROMACS the acceleration profile is:

$$a_x(z) = A \cos\left(\frac{2\pi z}{l_z}\right) \quad (6.8)$$

where l_z is the height of the box. The generated velocity profile is:

$$v_x(z) = V \cos\left(\frac{2\pi z}{l_z}\right) \quad (6.9)$$

$$V = A \frac{\rho}{\eta} \left(\frac{l_z}{2\pi}\right)^2 \quad (6.10)$$

The viscosity can be calculated from A and V :

$$\eta = \frac{A}{V} \rho \left(\frac{l_z}{2\pi}\right)^2 \quad (6.11)$$

In the simulation V is defined as:

$$V = \frac{\sum_{i=1}^N m_i v_{i,x} 2 \cos\left(\frac{2\pi z}{l_z}\right)}{\sum_{i=1}^N m_i} \quad (6.12)$$

The generated velocity profile is not coupled to the heat bath, moreover the velocity profile is excluded from the kinetic energy. One would like V to be as large as possible to get good statistics. However the shear rate should not be so high that the system gets too far from equilibrium. The maximum shear rate occurs where the cosine is zero, the rate being:

$$\text{sh}_{\max} = \max_z \left| \frac{\partial v_x(z)}{\partial z} \right| = A \frac{\rho}{\eta} \frac{l_z}{2\pi} \quad (6.13)$$

For a simulation with: $\eta = 10^{-3}$ [kg m⁻¹ s⁻¹], $\rho = 10^3$ [kg m⁻³] and $l_z = 2\pi$ [nm], $\text{sh}_{\max} = 1$ [ps nm⁻¹] A . This shear rate should be smaller than one over the longest correlation time in the system. For most liquids this will be the rotation correlation time, which is around 10 picoseconds. In this case A should be smaller than 0.1 [nm ps⁻²]. When the shear rate is too high, the observed viscosity will be too low. Because V is proportional to the square of the box height, the optimal box is elongated in the z -direction. In general a simulation length of 100 picoseconds is enough to obtain an accurate value for the viscosity.

The heat generated by the viscous friction is removed by coupling to a heat bath. Because this coupling is not instantaneous the real temperature of the liquid will be slightly lower than the

observed temperature. Berendsen derived this temperature shift[24], which can be written in terms of the shear rate as:

$$T_s = \frac{\eta \tau}{2\rho C_v} \text{sh}_{\text{max}}^2 \quad (6.14)$$

where τ is the coupling time for the Berendsen thermostat and C_v is the heat capacity. Using the values of the example above, $\tau = 10^{-13}$ [s] and $C_v = 2 \cdot 10^3$ [J kg⁻¹ K⁻¹], we get: $T_s = 25$ [K ps⁻²] sh_{max}^2 . When we want the shear rate to be smaller than 1/10 [ps⁻¹], T_s is smaller than 0.25 [K], which is negligible.

Note that the system has to build up the velocity profile when starting from an equilibrium state. This build-up time is of the order of the correlation time of the liquid.

Two quantities are written to the energy file, along with their averages and fluctuations: V and $1/\eta$ as obtained from (6.11).

6.6 Tabulated functions

6.6.1 Cubic splines for potentials

In some of the inner loops of GROMACS lookup tables are used for computation of potential and forces. The tables are interpolated using a cubic spline algorithm. There are separate tables for electrostatic, dispersion and repulsion interactions, but for the sake of caching performance these have been combined into a single array. The cubic spline interpolation looks like this:

$$y(x) = \eta y_i + \epsilon y_{i+1} + \frac{h^2}{6} [(\eta^3 - \eta)y_i'' + (\epsilon^3 - \epsilon)y_{i+1}''] \quad (6.15)$$

where $\epsilon = 1 - \eta$, and y_i and y_i'' are the tabulated values of a function $y(x)$ and its second derivative respectively. Furthermore,

$$h = x_{i+1} - x_i \quad (6.16)$$

$$\epsilon = (x - x_i)/h \quad (6.17)$$

so that $0 \leq \epsilon < 1$. eqn. 6.15 can be rewritten as

$$y(x) = y_i + \epsilon \left(y_{i+1} - y_i - \frac{h^2}{6} (2y_i'' + y_{i+1}'') \right) + \epsilon^2 \left(\frac{h^2}{2} y_i'' \right) + \epsilon^3 \frac{h^2}{6} (y_{i+1}'' - y_i'') \quad (6.18)$$

Note that the x-dependence is completely in ϵ . This can be abbreviated to

$$y(x) = y_i + \epsilon F_i + \epsilon^2 G_i + \epsilon^3 H_i \quad (6.19)$$

From this we can calculate the derivative in order to determine the forces:

$$\frac{dy(x)}{dx} = \frac{dy(x)}{d\epsilon} \frac{d\epsilon}{dx} = (F_i + 2\epsilon G_i + 3\epsilon^2 H_i)/h \quad (6.20)$$

If we store in the table y_i , F_i , G_i and H_i we need a table of length $4n$. The number of points per nanometer is 500 and 2000 for single and double precision compiled versions of GROMACS, respectively. The relative error is less than 10^{-4} when $n = 500$ points/nm. The force routines get a scaling factor s as a parameter that is equal to the number of points per nm. (Note that h is s^{-1}).

The algorithm goes a little something like this:

1. Calculate distance vector (\mathbf{r}_{ij}) and distance r_{ij}
2. Multiply r_{ij} by s and truncate to an integer value n_0 to get a table index
3. Calculate fractional component ($\epsilon = sr_{ij} - n_0$) and ϵ^2
4. Do the interpolation to calculate the potential V and the the scalar force f
5. Calculate the vector force \mathbf{F} by multiplying f with \mathbf{r}_{ij}

The tables are stored internally as y_i , F_i , G_i , H_i in the order coulomb, dispersion, repulsion. In total there are 12 values in each table entry. However, as evident from eqn. 6.15, only two values (y and y'') are necessary to determine each of the repulsion, dispersion and Coulomb potentials; from these y_i , F_i , G_i and H_i can be determined. Hence the GROMACS user only has to specify y and y'' , for each of these, leading to a total of six functions to be specified in the input to the `mdrun` program.

Note that table lookup is significantly *slower* than computation of the most simple Lennard-Jones and Coulomb interaction. However, it is much faster than the shifted coulomb function used in conjunction with the PPPM method. Finally it is much easier to modify a table for the potential (and get a graphical representation of it) than to modify the inner loops of the MD program.

6.6.2 User specified potential functions

You can also use your own potential functions without editing the GROMACS code. The potential function should be according to the following equation

$$V(r_{ij}) = \frac{q_i q_j}{4\pi\epsilon_0} f(r_{ij}) + C_6 g(r_{ij}) + C_{12} h(r_{ij}) \quad (6.21)$$

with f, g, h user defined functions. Note that if $g(r)$ represents a normal dispersion interaction, $g(r)$ should be < 0 . C_6 , C_{12} and the charges are read from the topology. Also note that combination rules are only supported for Lennard Jones and Buckingham, and that your tables should match the parameters in the binary topology.

When you add the following lines in your `.mdp` file:

```
rlist = 1.0
coulombtype = User
rcoulomb = 1.0
vdwtype = User
rvdw = 1.0
```

the MD program will read a single file (the name can be changed with option `-table`) with seven columns of table lookup data in the order: x , $f(x)$, $f''(x)$, $g(x)$, $g''(x)$, $h(x)$, $h''(x)$. The x should run from 0 to $r_c + 1$ (the value table extension can be changed in the `.mdp` file). You can choose the spacing you like; for the standard tables GROMACS using a spacing of 0.002 and 0.0005 nm when you run in single and double precision, respectively. In this context r_c denotes the maximum of the two cut-offs `rvdw` and `rcoulomb` (see above). These variables need not be the same (and need not be 1.0 either). Some functions used for potentials contain a singularity at $x = 0$, but since atoms are normally not closer to each other than 0.1 nm, the function value at $x =$

0 is not important. Finally, it is also possible to combine a standard Coulomb with a modified LJ potential (or vice versa). One then specifies e.g. `coulombtype = Cut-off` or `coulombtype = PME`, combined with `vdwtype = User`. The table file must always contain the 7 columns however, and meaningful data (i.e. not zeroes) must be entered in all columns. A number of pre-built table files can be found in the GMXLIB directory, for 6-8, 6-9, 6-10, 6-11, 6-12 Lennard Jones potentials combined with a normal Coulomb.

6.7 Mixed Quantum-Classical simulation techniques

In a molecular mechanics (MM) forcefield, the influence of electrons is expressed by empirical parameters that are assigned on the basis of experimental data, or on the basis of results from high-level quantum chemistry calculations. These are valid for the ground state of a given covalent structure, and the MM approximation is usually sufficiently accurate for ground-state processes in which the overall connectivity between the atoms in the system remains unchanged. However, for processes in which the connectivity does change, such as chemical reactions, or processes that involve multiple electronic states, such as photochemical conversions, electrons can no longer be ignored, and a quantum mechanical description is required for at least those parts of the system in which the reaction takes place.

One approach to the simulation of chemical reactions in solution, or in enzymes, is to use a combination of quantum mechanics (QM) and molecular mechanics (MM). The reacting parts of the system are treated quantum mechanically, with the remainder being modelled using the forcefield. The current version of Gromacs provides interfaces to several popular Quantum Chemistry packages (Mopac[81], Gamess-UK[82], Gaussian[83] and CPMD[84]).

Gromacs interactions between the two subsystems are either handled as described by Field *et al.*[85] or within the ONIOM approach by Morokuma and coworkers[86, 87].

6.7.1 Overview

Two approaches for describing the interactions between the QM and MM subsystems are supported in this version:

1. **Electronic Embedding** The electrostatic interactions between the electrons of the QM region and the MM atoms and between the QM nuclei and the MM atoms, are included in the Hamiltonian for the QM subsystem:

$$H^{QM/MM} = H_e^{QM} - \sum_i^n \sum_J^M \frac{e^2 Q_J}{4\pi\epsilon_0 r_{iJ}} + \sum_A^N \sum_J^M \frac{e^2 Z_A Q_J}{e\pi\epsilon_0 R_{AJ}}, \quad (6.22)$$

where n and N are the number of electrons and nuclei in the QM region, respectively, and M is the number of charged MM atoms. The first term on the right hand side is the original electronic Hamiltonian of an isolated QM system. The first of the double sums is the total electrostatic interaction between the QM electrons and the MM atoms. The total electrostatic interaction of the QM nuclei with the MM atoms is given by the second double sum. Bonded interactions between QM and MM atoms are described at the MM level by the

appropriate forcefield terms. Chemical bonds that connect the two subsystems are capped by a hydrogen atom to complete the valence of the QM region. The force on this atom, which is present in the QM region only, is distributed over the two atoms of the bond. The cap atom is usually referred to as a link atom.

2. **ONIOM** In the ONIOM approach, the energy and gradients are first evaluated for the isolated QM subsystem at the desired level of *ab initio* theory. Subsequently, the energy and gradients of the total system, including the QM region, are computed using the molecular mechanics forcefield and added to the energy and gradients calculated for the isolated QM subsystem. Finally in order to correct for counting the interactions inside the QM region twice, a molecular mechanics calculation is performed on the isolated QM subsystem and the energy and gradients are subtracted. This leads to the following expression for the total QM/MM energy (and gradients likewise):

$$E_{tot} = E_I^{QM} + E_{I+II}^{MM} - E_I^{MM}, \quad (6.23)$$

where the subscripts I and II refer to the QM and MM subsystems, respectively. The superscripts indicate at what level of theory the energies are computed. The ONIOM scheme has the advantage that it is not restricted to a two layer QM/MM description, but can easily handle more than two layers, with each layer described at a different level of theory.

6.7.2 Usage

To make use of the QM/MM functionality in Gromacs, one needs to:

1. introduce link atoms at the QM/MM boundary, if needed;
2. specify which atoms are to be treated at a QM level;
3. specify the QM level, basisset, type of QM/MM interface and so on.

Adding link atoms

At the bond that connects the QM and MM subsystems a link atom is introduced. In Gromacs the link atom has special atomtype, called LA. This atomtype is treated as a hydrogen atom in the QM calculation, and as a dummy atom in the forcefield calculation. The link atoms, if any, are part of the system, but have no interaction with any other atom, except that the QM force working on it is distributed over the two atoms of the bond. In the topology the link atom (LA), therefore, is defined as a virtual site atom:

```
[ virtual_sites2 ]
LA QMatom MMatom 1 0.65
```

See the dummy atoms section for more details on how dummies are treated. The link atom is replaced at every step of the simulation.

In addition, the bond itself is replaced by a constraint:

```
[ constraints ]
QMAtom MMAtom 2 0.153
```

Note that, because in our system the QM/MM bond is a carbon-carbon bond (0.153 nm), we use a constraint length of 0.153 nm, and dummy position of 0.65. The latter is the ratio between the ideal C-H bondlength and the ideal C-C bond length. With this ratio, the link atom is always 0.1 nm away from the QMAtom, consistent with the carbon-hydrogen bondlength. If the QM and MM subsystems are connected by a different kind of bond, a different constraint and a different dummy position, appropriate for that bond type, are required.

Specifying the QM atoms

Atoms that should be treated at a QM level of theory, including the link atoms, are added to the index file. In addition, the chemical bonds between the atoms in the QM region are to be defined as connect bonds (bond type 5) in the topology file:

```
[ bonds ]
QMAtom1 QMAtom2 5
QMAtom2 QMAtom3 5
```

Specifying the QM/MM simulation parameters

In the mdp file, the following parameters control a QM/MM simulation.

QMMM = no

If this is set to `yes`, a QM/MM simulation is requested. Several groups of atoms can be described at different QM levels separately. These are specified in the `QMMM-grps` field separated by spaces. The level of *ab initio* theory at which the groups are described is specified by `QMmethod` and `QMbasis` fields. Describing the groups at different levels of theory is only possible with the ONIOM QM/MM scheme, specified by `QMMMscheme`.

QMMM-grps =
groups to be described at the QM level

QMMMscheme = normal

Options are `normal` and `ONIOM`. This selects the QM/MM interface. `normal` implies that the QM subsystem is electronically embedded in the MM subsystem. There can only be one `QMMM-grps` that is modelled at the `QMmethod` and `QMbasis` level of *ab initio* theory. The rest of the system is described at the MM level. The QM and MM subsystems interact as follows: MM point charges are included in the QM one-electron hamiltonian and all Lennard-Jones interactions are described at the MM level. If `ONIOM` is selected, the interaction between the subsystem is described using the ONIOM method by Morokuma and co-workers. There can be more than one `QMMM-grps` each modelled at a different level of QM theory (`QMmethod` and `QMbasis`).

`QMmethod =`

Method used to compute the energy and gradients on the QM atoms. Available methods are AM1, PM3, RHF, UHF, DFT, B3LYP, MP2, CASSCF, MMVB and CPMD. For CASSCF, the number of electrons and orbitals included in the active space is specified by `CASelectrons` and `CASorbitals`. For CPMD, the planewave cut-off is specified by the `planewavecutoff` keyword.

`QMbasis =`

Gaussian basisset used to expand the electronic wavefunction. Only gaussian basissets are currently available, i.e. STO-3G, 3-21G, 3-21G*, 3-21+G*, 6-21G, 6-31G, 6-31G*, 6-31+G*, and 6-311G. For CPMD, which uses plane wave expansion rather than atom-centered basisfunctions, the `planewavecutoff` keyword controls the plane wave expansion.

`QMcharge =`

The total charge in e of the `QMMM-grps`. In case there are more than one `QMMM-grps`, the total charge of each ONIOM layer needs to be specified separately.

`QMmult =`

The multiplicity of the `QMMM-grps`. In case there are more than one `QMMM-grps`, the multiplicity of each ONIOM layer needs to be specified separately.

`CASorbitals =`

The number of orbitals to be included in the active space when doing a CASSCF computation.

`CASelectrons =`

The number of electrons to be included in the active space when doing a CASSCF computation.

`SH = no`

If this is set to yes, a QM/MM MD simulation on the excited state-potential energy surface and enforce a diabatic hop to the ground-state when the system hits the conical intersection hyperline in the course the simulation. This option only works in combination with the CASSCF method.

6.7.3 Output

The energies and gradients computed in the QM calculation are added to those computed by `gro-macs`. In the `.edr` file there is a section for the total QM energy.

6.7.4 Future developments

Several features are currently under development that increase the accuracy of the QM/MM interface. One useful feature is the use of delocalized MM charges in the QM computations. The most important benefit of using such smeared-out charges is that the coulombic potential has a finite value at inter atomic distances. In the point charge representation, the partially charged MM

atoms close to the QM region, tend to 'overpolarize' the QM system, which leads to artefacts in the calculation.

What is needed as well is a transition state optimizer.

Chapter 7

Run parameters and Programs

7.1 Online and html manuals

All the information in this chapter can also be found in HTML format in your GROMACS data directory. The path depends on where your files are installed, but the default location is

```
/usr/local/gromacs/share/html/online.html
```

Or, if you installed from Linux packages it can be found as

```
/usr/local/share/gromacs/html/online.html
```

You can also use the online from our web site,

www.gromacs.org/documentation/reference_3.0/online.html

In addition, we install standard UNIX manuals for all the programs. If you have sourced the GMXRC script in the GROMACS binary directory for your host they should already be present in your \$MANPATH, and you should be able to type e.g. `man grompp`.

The program manual pages can also be found in Appendix D in this manual.

7.2 File types

Table 7.1 lists the file types used by GROMACS along with a short description, and you can find a more detail description for each file in your HTML reference, or in our online version.

GROMACS files written in xdr format can be read on any architecture with GROMACS version 1.6 or later if the configuration script found the XDR libraries on your system. They should always be present on UNIX since they are necessary for NFS support.

Default Name	Ext.	Type	Default Option	Description
atomtp.atp		Asc		Atomtype file used by pdb2gmx
eiwit.brk		Asc	-f	Brookhaven data bank file
nnnice.dat		Asc		Generic data file
user.dlg		Asc		Dialog Box data for ngmx
sam.edi		Asc		ED sampling input
sam.edo		Asc		ED sampling output
ener.edr				Generic energy: edr ene
ener.edr		xdr		Energy file in portable xdr format
ener.ene		Bin		Energy file
eiwit.ent		Asc	-f	Entry in the protein date bank
plot.eps		Asc		Encapsulated PostScript (tm) file
gtraj.g87		Asc		Gromos-87 ASCII trajectory format
conf.g96		Asc	-c	Coordinate file in Gromos-96 format
conf.gro		Asc	-c	Coordinate file in Gromos-87 format
conf.gro			-c	Generic structure: gro g96 pdb tpr tpb tpa
out.gro			-o	Generic structure: gro g96 pdb
polar.hdb		Asc		Hydrogen data base
topinc.itp		Asc		Include file for topology
run.log		Asc	-l	Log file
ps.m2p		Asc		Input file for mat2ps
ss.map		Asc		File that maps matrix data to colors
ss.mat		Asc		Matrix Data file
grompp.mdp		Asc	-f	grompp input file with MD parameters
hessian.mtx		Bin	-m	Hessian matrix
index.ndx		Asc	-n	Index file
hello.out		Asc	-o	Generic output file
eiwit.pdb		Asc	-f	Protein data bank file
pull.pdo		Asc		Pull data output
pull.ppa		Asc		Pull parameters
residue.rtp		Asc		Residue Type file used by pdb2gmx
doc.tex		Asc	-o	LaTeX file
topol.top		Asc	-p	Topology file
topol.tpa		Asc	-s	Ascii run input file
topol.tpb		Bin	-s	Binary run input file
topol.tpr			-s	Generic run input: tpr tpb tpa
topol.tpr			-s	Structure+mass(db): tpr tpb tpa gro g96 pdb
topol.tpr		xdr	-s	Portable xdr run input file
traj.trj		Bin		Trajectory file (architecture specific)
traj.trr				Full precision trajectory: trr trj
traj.trr		xdr		Trajectory in portable xdr format
root.xpm		Asc		X PixMap compatible matrix file
traj.xtc			-f	Generic trajectory: xtc trr trj gro g96 pdb
traj.xtc		xdr		Compressed trajectory (portable xdr format)
graph.xvg		Asc	-o	xvgr/xmgr file

Table 7.1: The GROMACS file types.

7.3 Run Parameters

7.3.1 General

Default values are given in parentheses. The first option is always the default option. Units are given in square brackets. The difference between a dash and an underscore is ignored. A sample `.mdp` file is available. This should be appropriate to start a normal simulation. Edit it to suit your specific needs and desires.

7.3.2 Preprocessing

title:

this is redundant, so you can type anything you want

cpp: (cpp)

your preprocessor

include:

directories to include in your topology. Format:
`-I/home/john/my_lib -I../more_lib`

define: ()

defines to pass to the preprocessor, default is no defines. You can use any defines to control options in your customized topology files. Options that are already available by default are:

-DFLEXIBLE

Will tell grompp to include flexible water in stead of rigid water into your topology, this is necessary to make **conjugate gradients** or **l-bfgs** minimization work and will allow **steepest descent** to minimize further.

-DPOSRES

Will tell grompp to include `posre.itp` into your topology, used for position restraints.

7.3.3 Run control

integrator:

md

A leap-frog algorithm for integrating Newton's equations of motion.

sd

A leap-frog stochastic dynamics integrator. The temperature for one or more groups of atoms (**tc_grps**) is set with **ref_t** [K], the inverse friction constant for each group is set with **tau_t** [ps]. The parameter **tcoupl** is ignored. The random generator is initialized with **ld_seed**. **NOTE:** temperature deviations decay twice as fast as with a Berendsen thermostat with the same **tau_t**.

bd

An Euler integrator for Brownian or position Langevin dynamics, the velocity is the force divided by a friction coefficient (**bd_fric** [amu ps⁻¹]) plus random thermal noise (**ref_t**). When **bd_fric**=0, the friction coefficient for each particle is calculated as mass/tau_t , as for the integrator **sd**. The random generator is initialized with **ld_seed**.

The following algorithms are not integrators, but selected using

the integrator tag anyway

steep

A steepest descent algorithm for energy minimization. The maximum step size is **emstep** [nm], the tolerance is **emtol** [kJ mol⁻¹ nm⁻¹].

cg

A conjugate gradient algorithm for energy minimization, the tolerance is **emtol** [kJ mol⁻¹ nm⁻¹]. CG is more efficient when a steepest descent step is done every once in a while, this is determined by **nstcgsteep**. For a minimization prior to a normal mode analysis, which requires a very high accuracy, GROMACS should be compiled in double precision.

l-bfgs

A quasi-Newtonian algorithm for energy minimization according to the low-memory Broyden-Fletcher-Goldfarb-Shanno approach. In practice this seems to converge faster than Conjugate Gradients, but due to the correction steps necessary it is not (yet) parallelized.

nm

Normal mode analysis is performed on the structure in the **tpr** file. GROMACS should be compiled in double precision.

tpi

Test particle insertion. The last atom in the topology is the test particle. A trajectory should be provided with the **-rerun** option of **mdrun**. This trajectory should contain one atom less than the topology. Insertions are performed **nsteps** times in each frame. When **nstlist** is larger than one, after a first 'normal' insertion **nstlist-1** insertions are performed in a sphere with radius $0.05 * \text{rlist}$ around the 'first' location using the same neighborlist (and the same long-range energy when **rvdw** or **rcoulomb** > **rlist**). Since neighborlist construction is expensive one can perform several extra insertions with the same list almost for free. The random seed is set with **ld_seed**. The temperature for the Boltzmann weighting is set with **ref_t**, this should match the temperature of the simulation of the original trajectory. Dispersion correction is implemented correctly for **tpi**. All relevant quantities are written using the **-tpi** option of **mdrun**. No trajectory or energy file is written.

tinit: (0) [ps]

starting time for your run (only makes sense for integrators **md**, **sd** and **bd**)

dt: (0.001) [ps]

time step for integration (only makes sense for integrators **md**, **sd** and **bd**)

nsteps: (0)

maximum number of steps to integrate

init_step: (0)

The starting step. The time at a step i in a run is calculated as: $t = t_{\text{init}} + dt * (\text{init_step} + i)$. The free-energy λ is calculated as: $\lambda = \text{init_lambda} + \text{delta_lambda} * (\text{init_step} + i)$. Also non-equilibrium MD parameters can depend on the step number. Thus for exact restarts or redoing part of a run it might be necessary to set `init_step` to the step number of the restart frame. `tpbconv` does this automatically.

comm_mode:**Linear**

Remove center of mass translation

Angular

Remove center of mass translation and rotation around the center of mass

No

No restriction on the center of mass motion

nstcomm: (1) [steps]

frequency for center of mass motion removal

comm_grps:

group(s) for center of mass motion removal, default is the whole system

7.3.4 Langevin dynamics**bd_fric: (0) [amu ps⁻¹]**

Brownian dynamics friction coefficient. When `bd_fric=0`, the friction coefficient for each particle is calculated as mass/tau_t .

ld_seed: (1993) [integer]

used to initialize random generator for thermal noise for stochastic and Brownian dynamics. When `ld_seed` is set to -1, the seed is calculated as $(\text{time}() + \text{getpid}()) \% 1000000$. When running on multiple processors, each processor uses a seed equal to `ld_seed` plus the processor number.

7.3.5 Energy minimization**emtol: (10.0) [kJ mol⁻¹ nm⁻¹]**

the minimization is converged when the maximum force is smaller than this value

emstep: (0.01) [nm]

initial step-size

nstcgsteep: (1000) [steps]

frequency of performing 1 steepest descent step while doing conjugate gradient energy minimization.

nbgscorr: (10)

Number of correction steps to use for L-BFGS minimization. A higher number is (at least theoretically) more accurate, but slower.

7.3.6 Shell Molecular Dynamics

When shells or flexible constraints are present in the system the positions of the shells and the lengths of the flexible constraints are optimized at every time step until either the RMS force on the shells and constraints is less than `emtol`, or a maximum number of iterations (`niter`) has been reached

emtol: (10.0) [kJ mol⁻¹ nm⁻¹]

the minimization is converged when the maximum force is smaller than this value. For shell MD this value should be 1.0 at most, but since the variable is used for energy minimization as well the default is 10.0.

niter: (20)

maximum number of iterations for optimizing the shell positions and the flexible constraints.

fcstep: (0) [ps²]

the step size for optimizing the flexible constraints. Should be chosen as $\mu/(d^2V/dq^2)$ where μ is the reduced mass of two particles in a flexible constraint and d^2V/dq^2 is the second derivative of the potential in the constraint direction. Hopefully this number does not differ too much between the flexible constraints, as the number of iterations and thus the runtime is very sensitive to `fcstep`. Try several values!

7.3.7 Output control

nstxout: (100) [steps]

frequency to write coordinates to output trajectory file, the last coordinates are always written

nstvout: (100) [steps]

frequency to write velocities to output trajectory, the last velocities are always written

nstfout: (0) [steps]

frequency to write forces to output trajectory.

nstlog: (100) [steps]

frequency to write energies to log file, the last energies are always written

nstenergy: (100) [steps]

frequency to write energies to energy file, the last energies are always written

nstxtcout: (0) [steps]

frequency to write coordinates to xtc trajectory

xtc_precision: (1000) [real]

precision to write to xtc trajectory

xtc_grps:

group(s) to write to xtc trajectory, default the whole system is written (if **nstxtcout** is larger than zero)

energygrps:

group(s) to write to energy file

7.3.8 Neighbor searching

nstlist: (10) [steps]

Frequency to update the neighbor list (and the long-range forces, when using twin-range cut-off's). When this is 0, the neighbor list is made only once.

ns_type:**grid**

Make a grid in the box and only check atoms in neighboring grid cells when constructing a new neighbor list every **nstlist** steps. In large systems grid search is much faster than simple search.

simple

Check every atom in the box when constructing a new neighbor list every **nstlist** steps.

pbc:**xyz**

Use periodic boundary conditions in all directions.

no

Use no periodic boundary conditions, ignore the box. To simulate without cut-offs, set all cut-offs to 0 and **nstlist**=0.

rlist: (1) [nm]

cut-off distance for the short-range neighbor list

7.3.9 Electrostatics

coulombtype:**Cut-off**

Twin range cut-off's with neighborlist cut-off **rlist** and Coulomb cut-off **rcoulomb**, where **rcoulomb** > **rlist**.

Ewald

Classical Ewald sum electrostatics. The real-space cut-off **rcoulomb** should be equal to **rlist**. Use *e.g.* **rlist**=0.9, **rcoulomb**=0.9. The highest magnitude of wave vectors

used in reciprocal space is controlled by **fourierspacing**. The relative accuracy of direct/reciprocal space is controlled by **ewald_rtol**.

NOTE: Ewald scales as $O(N^{3/2})$ and is thus extremely slow for large systems. It is included mainly for reference - in most cases PME will perform much better.

PME

Fast Particle-Mesh Ewald electrostatics. Direct space is similar to the Ewald sum, while the reciprocal part is performed with FFTs. Grid dimensions are controlled with **fourierspacing** and the interpolation order with **pme_order**. With a grid spacing of 0.1 nm and cubic interpolation the electrostatic forces have an accuracy of $2\text{-}3\text{e-}4$. Since the error from the vdw-cutoff is larger than this you might try 0.15 nm. When running in parallel the interpolation parallelizes better than the FFT, so try decreasing grid dimensions while increasing interpolation.

PPPM

Particle-Particle Particle-Mesh algorithm for long range electrostatic interactions. Use for example **rlist**=0.9, **rcoulomb**=0.9. The grid dimensions are controlled by **fourierspacing**. Reasonable grid spacing for PPPM is 0.05-0.1 nm. See `Shift` for the details of the particle-particle potential.

NOTE: the pressure is incorrect when using PPPM.

Reaction-Field

Reaction field with Coulomb cut-off **rcoulomb**, where **rcoulomb** > **rlist**. The dielectric constant beyond the cut-off is **epsilon_rf**. The dielectric constant can be set to infinity by setting **epsilon_rf**=0.

Generalized-Reaction-Field

Generalized reaction field with Coulomb cut-off **rcoulomb**, where **rcoulomb** > **rlist**. The dielectric constant beyond the cut-off is **epsilon_rf**. The ionic strength is computed from the number of charged (*i.e.* with non zero charge) charge groups. The temperature for the GRF potential is set with **ref.t** [K].

Reaction-Field-nec

The same as **Reaction-Field**, but implemented as in GROMACS versions before 3.3. No reaction-field correction is applied to excluded atom pairs and self pairs. The 1-4 interactions are calculated using a reaction-field. The missing correction due to the excluded pairs that do not have a 1-4 interaction is up to a few percent of the total electrostatic energy and causes a minor difference in the forces and the pressure.

Shift

The Coulomb potential is decreased over the whole range and the forces decay smoothly to zero between **rcoulomb_switch** and **rcoulomb**. The neighbor search cut-off **rlist** should be 0.1 to 0.3 nm larger than **rcoulomb** to accommodate for the size of charge groups and diffusion between neighbor list updates.

Encad-Shift

The Coulomb potential is decreased over the whole range, using the definition from the Encad simulation package.

Switch

The Coulomb potential is normal out to **rcoulomb_switch**, after which it is switched off to reach zero at **rcoulomb**. Both the potential and force functions are continuously

smooth, but be aware that all switch functions will give rise to a bulge (increase) in the force (since we are switching the potential). The neighbor search cut-off **rlist** should be 0.1 to 0.3 nm larger than **rcoulomb** to accommodate for the size of charge groups and diffusion between neighbor list updates.

User

`mdrun` will now expect to find a file `table.xvg` with user-defined potential functions for repulsion, dispersion and Coulomb. This file should contain 7 columns: the x value, $f(x)$, $-f^{(2)}(x)$, $g(x)$, $-g^{(2)}(x)$, $h(x)$, $-h^{(2)}(x)$, where $f^{(2)}(x)$ denotes the 2nd derivative of function $f(x)$ with respect to x . $f(x)$ is the Coulomb function, $g(x)$ the dispersion function and $h(x)$ the repulsion function. The x values should run from 0 to the largest cut-off distance + **table-extension** and should be uniformly spaced. The optimal spacing, which is used for non-user tables, is 0.002 [nm] when you run in single precision or 0.0005 [nm] when you run in double precision. The function value at $x=0$ is not important. More information is in the printed manual.

PME-User

A combination of PME and user tables (see above). The PME mesh contribution is subtracted from the user table by `mdrun`. Because of this subtraction the user tables should contain about 10 decimal places.

rcoulomb_switch: (0) [nm]

where to start switching the Coulomb potential

rcoulomb: (1) [nm]

distance for the Coulomb cut-off

epsilon_r: (1)

The relative dielectric constant. A value of 0 means infinity.

epsilon_rf: (1)

The relative dielectric constant of the reaction field. This is only used with reaction-field electrostatics. A value of 0 means infinity.

7.3.10 VdW

vdwtype:

Cut-off

Twin range cut-off's with neighbor list cut-off **rlist** and VdW cut-off **rvdw**, where **rvdw** > **rlist**.

Shift

The LJ (not Buckingham) potential is decreased over the whole range and the forces decay smoothly to zero between **rvdw_switch** and **rvdw**. The neighbor search cut-off **rlist** should be 0.1 to 0.3 nm larger than **rvdw** to accommodate for the size of charge groups and diffusion between neighbor list updates.

Switch

The LJ (not Buckingham) potential is normal out to **rvdw_switch**, after which it is switched off to reach zero at **rvdw**. Both the potential and force functions are continuously smooth, but be aware that all switch functions will give rise to a bulge (increase) in the force (since we are switching the potential). The neighbor search cut-off **rlist** should be 0.1 to 0.3 nm larger than **rvdw** to accommodate for the size of charge groups and diffusion between neighbor list updates.

Encad-Shift

The LJ (not Buckingham) potential is decreased over the whole range, using the definition from the Encad simulation package.

User

See above The function value at $x=0$ is not important. When you want to use LJ correction, make sure that **rvdw** corresponds to the cut-off in the user-defined function.

rvdw_switch: (0) [nm]

where to start switching the LJ potential

rvdw: (1) [nm]

distance for the LJ or Buckingham cut-off

DispCorr:**no**

don't apply any correction

EnerPres

apply long range dispersion corrections for Energy and Pressure

Ener

apply long range dispersion corrections for Energy only

7.3.11 Tables**table-extension: (1) [nm]**

Extension of the non-bonded potential lookup tables beyond the largest cut-off distance. The value should be large enough to account for charge group sizes and the diffusion between neighbor-list updates. This value also specifies the length of the lookup tables for the 1-4 interactions, which are always tabulated irrespective of the use of tables for the non-bonded interactions.

energygrp_table:

When user tables are used for electrostatics and/or VdW, here one can give pairs of energy groups for which separate user tables should be used. The two energy groups will be appended to the table file name, in order of their definition in **energygrps**, separated by underscores. For example, if `energygrps = Na Cl Sol` and `energygrp_table = Na Na Na Cl`, `mdrun` will read `table_Na_Na.xvg` and `table_Na_Cl.xvg` in addition to the normal `table.xvg` which will be used for all other energy group pairs.

7.3.12 Ewald

fourierspacing: (0.12) [nm]

The maximum grid spacing for the FFT grid when using PPPM or PME. For ordinary Ewald the spacing times the box dimensions determines the highest magnitude to use in each direction. In all cases each direction can be overridden by entering a non-zero value for **fourier_n***.

fourier_nx (0) ; fourier_ny (0) ; fourier_nz: (0)

Highest magnitude of wave vectors in reciprocal space when using Ewald. Grid size when using PPPM or PME. These values override **fourierspacing** per direction. The best choice is powers of 2, 3, 5 and 7. Avoid large primes.

pme_order (4)

Interpolation order for PME. 4 equals cubic interpolation. You might try 6/8/10 when running in parallel and simultaneously decrease grid dimension.

ewald_rtol (1e-5)

The relative strength of the Ewald-shifted direct potential at the cutoff is given by **ewald_rtol**. Decreasing this will give a more accurate direct sum, but then you need more wave vectors for the reciprocal sum.

ewald_geometry: (3d)

The geometry to use for Ewald summations. **3d** means the sum is performed in all three dimensions. If your system has a slab geometry in the x-y plane you can try to increase box z dimension and use the **3dc** geometry. The reciprocal sum is still performed in 3d, but a force and potential correction applied in the z dimension to produce a pseudo-2d summation. In the future there might also be a true **2d** option, but this is not working yet.

epsilon_surface: (0)

This controls the dipole correction to the Ewald summation in 3d. The default value of zero means it is turned off. Turn it on by setting it to the value of the relative permittivity of the imaginary surface around your infinite system. Be careful - you shouldn't use this if you have free mobile charges in your system. This value does not affect the slab 3DC variant of the long range corrections.

optimize_fft:**no**

Don't calculate the optimal FFT plan for the grid at startup.

yes

Calculate the optimal FFT plan for the grid at startup. This saves a few percent for long simulations, but takes a couple of minutes at start.

7.3.13 Temperature coupling

tcoupl:

no

No temperature coupling.

berendsen

Temperature coupling with a Berendsen-thermostat to a bath with temperature **ref.t** [K], with time constant **tau.t** [ps]. Several groups can be coupled separately, these are specified in the **tc_grps** field separated by spaces.

nose-hoover

Temperature coupling with a by using a Nose-Hoover extended ensemble. The reference temperature and coupling groups are selected as above, but in this case **tau.t** [ps] controls the period of the temperature fluctuations at equilibrium, which is slightly different from a relaxation time.

tc_grps:

groups to couple separately to temperature bath

tau.t: [ps]

time constant for coupling (one for each group in **tc_grps**), 0 means no temperature coupling

ref.t: [K]

reference temperature for coupling (one for each group in **tc_grps**)

7.3.14 Pressure coupling**pcoupl:****no**

No pressure coupling. This means a fixed box size.

berendsen

Exponential relaxation pressure coupling with time constant **tau.p** [ps]. The box is scaled every timestep. It has been argued that this does not yield a correct thermodynamic ensemble, but it is the most efficient way to scale a box at the beginning of a run.

Parrinello-Rahman

Extended-ensemble pressure coupling where the box vectors are subject to an equation of motion. The equation of motion for the atoms is coupled to this. No instantaneous scaling takes place. As for Nose-Hoover temperature coupling the time constant **tau.p** [ps] is the period of pressure fluctuations at equilibrium. This is probably a better method when you want to apply pressure scaling during data collection, but beware that you can get very large oscillations if you are starting from a different pressure.

pcoupltype:**isotropic**

Isotropic pressure coupling with time constant **tau.p** [ps]. The compressibility and reference pressure are set with **compressibility** [bar^{-1}] and **ref.p** [bar], one value is needed.

semiisotropic

Pressure coupling which is isotropic in the x and y direction, but different in the z direction. This can be useful for membrane simulations. 2 values are needed for x/y and z directions respectively.

anisotropic

Idem, but 6 values are needed for xx, yy, zz, xy/yx, xz/zx and yz/zy components respectively. When the off-diagonal compressibilities are set to zero, a rectangular box will stay rectangular. Beware that anisotropic scaling can lead to extreme deformation of the simulation box.

surface-tension

Surface tension coupling for surfaces parallel to the xy-plane. Uses normal pressure coupling for the z-direction, while the surface tension is coupled to the x/y dimensions of the box. The first **ref_p** value is the reference surface tension times the number of surfaces [bar nm], the second value is the reference z-pressure [bar]. The two **compressibility** [bar⁻¹] values are the compressibility in the x/y and z direction respectively. The value for the z-compressibility should be reasonably accurate since it influences the convergence of the surface-tension, it can also be set to zero to have a box with constant height.

tau_p: (1) [ps]

time constant for coupling

compressibility: [bar⁻¹]

compressibility (NOTE: this is now really in bar⁻¹) For water at 1 atm and 300 K the compressibility is 4.5e-5 [bar⁻¹].

ref_p: [bar]

reference pressure for coupling

7.3.15 Simulated annealing

Simulated annealing is controlled separately for each temperature group in GROMACS. The reference temperature is a piecewise linear function, but you can use an arbitrary number of points for each group, and choose either a single sequence or a periodic behaviour for each group. The actual annealing is performed by dynamically changing the reference temperature used in the thermostat algorithm selected, so remember that the system will usually not instantaneously reach the reference temperature!

annealing:

Type of annealing for each temperature group

no

No simulated annealing - just couple to reference temperature value.

single

A single sequence of annealing points. If your simulation is longer than the time of the last point, the temperature will be coupled to this constant value after the annealing sequence has reached the last time point.

periodic

The annealing will start over at the first reference point once the last reference time is reached. This is repeated until the simulation ends.

annealing_npoints:

A list with the number of annealing reference/control points used for each temperature group. Use 0 for groups that are not annealed. The number of entries should equal the number of temperature groups.

annealing_time:

List of times at the annealing reference/control points for each group. If you are using periodic annealing, the times will be used modulo the last value, *i.e.* if the values are 0, 5, 10, and 15, the coupling will restart at the 0ps value after 15ps, 30ps, 45ps, etc. The number of entries should equal the sum of the numbers given in annealing_npoints.

annealing_temp:

List of temperatures at the annealing reference/control points for each group. The number of entries should equal the sum of the numbers given in annealing_npoints.

Confused? OK, let's use an example. Assume you have two temperature groups, set the group selections to `annealing = single periodic`, the number of points of each group to `annealing_npoints = 3 4`, the times to `annealing_time = 0 3 6 0 2 4 6` and finally temperatures to `annealing_temp = 298 280 270 298 320 320 298`. The first group will be coupled to 298K at 0ps, but the reference temperature will drop linearly to reach 280K at 3ps, and then linearly between 280K and 270K from 3ps to 6ps. After this it stays constant, at 270K. The second group is coupled to 298K at 0ps, it increases linearly to 320K at 2ps, where it stays constant until 4ps. Between 4ps and 6ps it decreases to 298K, and then it starts over with the same pattern again, *i.e.* rising linearly from 298K to 320K between 6ps and 8ps. Check the summary printed by `grompp` if you are unsure!

7.3.16 Velocity generation**gen_vel:****no**

Do not generate velocities at startup. The velocities are set to zero when there are no velocities in the input structure file.

yes

Generate velocities according to a Maxwell distribution at temperature **gen_temp** [K], with random seed **gen_seed**. This is only meaningful with integrator **md**.

gen_temp: (300) [K]

temperature for Maxwell distribution

gen_seed: (173529) [integer]

used to initialize random generator for random velocities, when **gen_seed** is set to -1, the seed is calculated as `(time() + getpid()) % 1000000`

7.3.17 Bonds

constraints:**none**

No constraints except for those defined explicitly in the topology, *i.e.* bonds are represented by a harmonic (or other) potential or a Morse potential (depending on the setting of **morse**) and angles by a harmonic (or other) potential.

hbonds

Convert the bonds with H-atoms to constraints.

all-bonds

Convert all bonds to constraints.

h-angles

Convert all bonds and additionally the angles that involve H-atoms to bond-constraints.

all-angles

Convert all bonds and angles to bond-constraints.

constraint_algorithm:**lincs**

LINear Constraint Solver. The accuracy is set with **lincs_order**, which sets the number of matrices in the expansion for the matrix inversion, 4 is enough for a "normal" MD simulation, 8 is needed for BD with large time-steps. After the matrix inversion correction the algorithm does an iterative correction to compensate for lengthening due to rotation. The number of such iterations can be controlled with **lincs_iter**. The accuracy of the constraints is printed to the log file every **nstlog** steps. If a bond rotates more than **lincs_warnangle** [degrees] in one step, a warning will be printed both to the log file and to `stderr`. Lincs should not be used with coupled angle constraints.

shake

Shake is slower and less stable than Lincs, but does work with angle constraints. The relative tolerance is set with **shake_tol**, 0.0001 is a good value for "normal" MD.

unconstrained_start:**no**

apply constraints to the start configuration and reset shells

yes

do not apply constraints to the start configuration and do not reset shells, useful for exact continuation and reruns

shake_tol: (0.0001)

relative tolerance for shake

lincs_order: (4)

Highest order in the expansion of the constraint coupling matrix.

lincs_iter: (1)

Number of iterations to correct for rotational lengthening in Lincs. For normal runs a single step is sufficient, but for NVE runs where you want to conserve energy accurately you might want to increase it to 2, and for energy minimization with constraints it is sometimes necessary to use 4 or even 8 (double precision) iteration steps.

lincs_warnangle: (30) [degrees]

maximum angle that a bond can rotate before Lincs will complain

morse:**no**

bonds are represented by a harmonic potential

yes

bonds are represented by a Morse potential

7.3.18 Energy group exclusions

energygrp_excl:

Pairs of energy groups for which all non-bonded interactions are excluded. An example: if you have two energy groups `Protein` and `SOL`, specifying

```
energygrp_excl = Protein Protein SOL SOL
```

would give only the non-bonded interactions between the protein and the solvent. This is especially useful for speeding up energy calculations with `mdrun -rerun` and for excluding interactions within frozen groups.

7.3.19 NMR refinement

disre:**no**

no distance restraints (ignore distance restraint information in topology file)

simple

simple (per-molecule) distance restraints, ensemble averaging can be performed with `mdrun -multi`

ensemble

distance restraints over an ensemble of molecules in one simulation box, should only be used for special cases, such as dimers

disre_weighting:**conservative**

the forces are the derivative of the restraint potential, this results in an r^{-7} weighting of the atom pairs

equal

divide the restraint force equally over all atom pairs in the restraint

disre_mixed:**no**

the violation used in the calculation of the restraint force is the time averaged violation

yes

the violation used in the calculation of the restraint force is the square root of the time averaged violation times the instantaneous violation

disre_fc: (1000) [kJ mol⁻¹ nm⁻²]

force constant for distance restraints, which is multiplied by a (possibly) different factor for each restraint

disre_tau: (0) [ps]

time constant for distance restraints running average

nstdisreout: (100) [steps]

frequency to write the running time averaged and instantaneous distances of all atom pairs involved in restraints to the energy file (can make the energy file very large)

orire:**no**

no orientation restraints (ignore orientation restraint information in topology file)

yes

use orientation restraints, ensemble averaging can be performed with `mdrun -multi`

orire_fc: (0) [kJ mol]

force constant for orientation restraints, which is multiplied by a (possibly) different factor for each restraint, can be set to zero to obtain the orientations from a free simulation

orire_tau: (0) [ps]

time constant for orientation restraints running average

orire_fitgrp:

fit group for orientation restraining, for a protein backbone is a good choice

nstorireout: (100) [steps]

frequency to write the running time averaged and instantaneous orientations for all restraints and the molecular order tensor to the energy file (can make the energy file very large)

7.3.20 Free Energy Perturbation

free_energy:**no**

Only use topology A.

yes

Interpolate between topology A (lambda=0) to topology B (lambda=1) and write the derivative of the Hamiltonian with respect to lambda to the energy file and to `dgd1.xvg`. The potentials, bond-lengths and angles are interpolated linearly as described in the manual. When **sc_alpha** is larger than zero, soft-core potentials are used for the LJ and Coulomb interactions.

init_lambda: (0)

starting value for lambda

delta_lambda: (0)

increment per time step for lambda

sc_alpha: (0)

the soft-core parameter, a value of 0 results in linear interpolation of the LJ and Coulomb interactions

sc_power: (0)

the power for lambda in the soft-core function, only the values 1 and 2 are supported

sc_sigma: (0.3) [nm]

the soft-core sigma for particles which have a C6 or C12 parameter equal to zero

7.3.21 Non-equilibrium MD

acc_grps:

groups for constant acceleration (*e.g.*: `Protein Sol`) all atoms in groups `Protein` and `Sol` will experience constant acceleration as specified in the **accelerate** line

accelerate: (0) [nm ps⁻²]

acceleration for **acc_grps**; x, y and z for each group (*e.g.* `0.1 0.0 0.0 -0.1 0.0 0.0` means that first group has constant acceleration of 0.1 nm ps⁻² in X direction, second group the opposite).

freezegrps:

Groups that are to be frozen (*i.e.* their X, Y, and/or Z position will not be updated; *e.g.* `Lipid SOL`). **frezedim** specifies for which dimension the freezing applies. To avoid spurious contributions to the virial and pressure due to large forces between completely frozen atoms you need to use energy group exclusions, this also saves computing time. Note that frozen coordinates are not subject to pressure scaling.

frezedim:

dimensions for which groups in **freezegrps** should be frozen, specify Y or N for X, Y and Z and for each group (*e.g.* `Y Y N N N N` means that particles in the first group can move only in Z direction. The particles in the second group can move in any direction).

cos_acceleration: (0) [nm ps⁻²]

the amplitude of the acceleration profile for calculating the viscosity. The acceleration is in the X-direction and the magnitude is **cos_acceleration** $\cos(2 \pi z/\text{boxheight})$. Two terms are added to the energy file: the amplitude of the velocity profile and 1/viscosity.

deform: (0 0 0 0 0) [nm ps⁻¹]

The velocities of deformation for the box elements: a(x) b(y) c(z) b(x) c(x) c(y). Each step the box elements for which **deform** is non-zero are calculated as: $\text{box}(ts) + (t-ts) \cdot \text{deform}$, off-diagonal elements are corrected for periodicity. The coordinates are transformed accordingly. Frozen degrees of freedom are (purposely) also transformed. The time ts is set to t at the first step and at steps at which x and v are written to trajectory to ensure exact restarts. Deformation can be used together with semiisotropic or anisotropic pressure coupling when the appropriate compressibilities are set to zero. The diagonal elements can be used to strain a solid. The off-diagonal elements can be used to shear a solid or a liquid.

7.3.22 Electric fields**E_x ; E_y ; E_z:**

If you want to use an electric field in a direction, enter 3 numbers after the appropriate E_* , the first number: the number of cosines, only 1 is implemented (with frequency 0) so enter 1, the second number: the strength of the electric field in V nm^{-1} , the third number: the phase of the cosine, you can enter any number here since a cosine of frequency zero has no phase.

E_xt; E_yt; E_zt:

not implemented yet

7.3.23 Mixed quantum/classical molecular dynamics**QMMM:**

no

No QM/MM.

yes

Do a QM/MM simulation. Several groups can be described at different QM levels separately. These are specified in the **QMMM-grps** field separated by spaces. The level of i_i ab initio/ i_i theory at which the groups are described is specified by **QM-method** and **QMbasis** Fields. Describing the groups at different levels of theory is only possible with the ONIOM QM/MM scheme, specified by **QMMMscheme**.

QMMM-grps:

groups to be described at the QM level

QMMMscheme:

normal

normal QM/MM. There can only be one **QMMM-grps** that is modelled at the **QM-method** and **QMbasis** level of i_i ab initio/ i_i theory. The rest of the system is described at the MM level. The QM and MM subsystems interact as follows: MM point charges are included in the QM one-electron hamiltonian and all Lennard-Jones interactions are described at the MM level.

ONIOM

The interaction between the subsystem is described using the ONIOM method by Morokuma and co-workers. There can be more than one **QMMM-grps** each modeled at a different level of QM theory (**QMmethod** and **QMbasis**).

QMmethod: (RHF)

Method used to compute the energy and gradients on the QM atoms. Available methods are AM1, PM3, RHF, UHF, DFT, B3LYP, MP2, CASSCF, and MMVB. For CASSCF, the number of electrons and orbitals included in the active space is specified by **CASelectrons** and **CASorbitals**.

QMbasis: (STO-3G)

Basisset used to expand the electronic wavefunction. Only gaussian basissets are currently available, i_l i.e. i_l STO-3G, 3-21G, 3-21G*, 3-21+G*, 6-21G, 6-31G, 6-31G*, 6-31+G*, and 6-311G.

QMcharge: (0) [integer]

The total charge in i_l of the **QMMM-grps**. In case there are more than one **QMMM-grps**, the total charge of each ONIOM layer needs to be specified separately.

QMmult: (1) [integer]

The multiplicity of the **QMMM-grps**. In case there are more than one **QMMM-grps**, the multiplicity of each ONIOM layer needs to be specified separately.

CASorbitals: (0) [integer]

The number of orbitals to be included in the active space when doing a CASSCF computation.

CASelectrons: (0) [integer]

The number of electrons to be included in the active space when doing a CASSCF computation.

SH:**no**

No surface hopping. The system is always in the electronic ground-state.

yes

Do a QM/MM MD simulation on the excited state-potential energy surface and enforce a i_l diabatic i_l hop to the ground-state when the system hits the conical intersection hyperline in the course the simulation. This option only works in combination with the CASSCF method.

7.3.24 User defined thingies**user1_grps; user2_grps:****userint1 (0); userint2 (0); userint3 (0); userint4 (0)**

userreal1 (0); userreal2 (0); userreal3 (0); userreal4 (0)

These you can use if you modify code. You can pass integers and reals to your subroutine. Check the inputrec definition in `src/include/types/inputrec.h`

7.4 Programs by topic

Generating topologies and coordinates

pdb2gmx	converts pdb files to topology and coordinate files
x2top	generates a primitive topology from coordinates
editconf	edits the box and writes subgroups
genbox	solvates a system
genion	generates mono atomic ions on energetically favorable positions
genconf	multiplies a conformation in 'random' orientations
genpr	generates position restraints for index groups
protonate	protonates structures

Running a simulation

grompp	makes a run input file
tpbconv	makes a run input file for restarting a crashed run
mdrun	performs a simulation

Viewing trajectories

ngmx	displays a trajectory
trjconv	converts trajectories to <i>e.g.</i> pdb which can be viewed with <i>e.g.</i> rasmol

Processing energies

g_energy	writes energies to xvg files and displays averages
g_enemat	extracts an energy matrix from an energy file
mdrun	with <code>-rerun (re)</code> calculates energies for trajectory frames

Converting files

editconf	converts and manipulates structure files
trjconv	converts and manipulates trajectory files
trjcat	concatenates trajectory files
eneconv	converts energy files
xmp2ps	converts XPM matrices to encapsulated postscript (or XPM)

Tools

make_ndx	makes index files
mk_angndx	generates index files for g_angle
gmxcheck	checks and compares files
gmxdump	makes binary files human readable
g_traj	plots x, v and f of selected atoms/groups (and more) from a trajectory
g_analyze	analyzes data sets
trjorder	orders molecules according to their distance to a group
g_filter	frequency filters trajectories, useful for making smooth movies
g_lie	free energy estimate from linear combinations
g_dyndom	interpolate and extrapolate structure rotations
g_morph	linear interpolation of conformations
g_wham	weighted histogram analysis after umbrella sampling
ffscan	scan and modify force field data for a single point energy calculation
xpm2ps	convert XPM (XPixelMap) file to postscript

Distances between structures

g_rms	calculates rmsd's with a reference structure and rmsd matrices
g_confrms	fits two structures and calculates the rmsd
g_cluster	clusters structures
g_rmsf	calculates atomic fluctuations
disco	distance geometry calculation with the CONCOORD algorithm
cdist	create input for disco

Distances in structures over time

g_mindist	calculates the minimum distance between two groups
g_dist	calculates the distances between the centers of mass of two groups
g_bond	calculates distances between atoms
g_mdmat	calculates residue contact maps
g_rmsdist	calculates atom pair distances averaged with power 2, -3 or -6

Mass distribution properties over time

g_traj	plots x, v, f, box, temperature and rotational energy
g_gyrate	calculates the radius of gyration
g_msd	calculates mean square displacements
g_rotacf	calculates the rotational correlation function for molecules
g_rdf	calculates radial distribution functions

Analyzing bonded interactions

g_bond	calculates bond length distributions
mk_angndx	generates index files for g_angle
g_angle	calculates distributions and correlations for angles and dihedrals
g_dih	analyzes dihedral transitions

Structural properties

g_hbond	computes and analyzes hydrogen bonds
g_saltbr	computes salt bridges
g_sas	computes solvent accessible surface area
g_order	computes the order parameter per atom for carbon tails
g_sgangle	computes the angle and distance between two groups
g_sorient	analyzes solvent orientation around solutes
g_bundle	analyzes bundles of axes, <i>e.g.</i> helices
g_disre	analyzes distance restraints
g_clustsize	calculate size distributions of atomic clusters
disco	distance geometry calculation with the CONCOORD algorithm
cdist	create input for disco
anadock	cluster structures from Autodock runs

Kinetic properties

g_traj	plots x, v, f, box, temperature and rotational energy
g_velacc	calculates velocity autocorrelation functions
g_tcaf	calculates viscosities of liquids
g_kinetics	derives information about kinetic processes from you trajectories

Electrostatic properties

genion	generates mono atomic ions on energetically favorable positions
g_potential	calculates the electrostatic potential across the box
g_dipoles	computes the total dipole plus fluctuations
g_dielectric	calculates frequency dependent dielectric constants

Protein specific analysis

do_dssp	assigns secondary structure and calculates solvent accessible surface area
g_chi	calculates everything you want to know about chi and other dihedrals
g_helix	calculates everything you want to know about helices
g_rama	computes Ramachandran plots
xrama	shows animated Ramachandran plots
wheel	plots helical wheels

Interfaces

g_potential	calculates the electrostatic potential across the box
g_density	calculates the density of the system
g_densmap	calculates 2D planar or axial-radial density maps
g_order	computes the order parameter per atom for carbon tails
g_h2order	computes the orientation of water molecules
g_bundle	analyzes bundles of axes, <i>e.g.</i> transmembrane helices

Covariance analysis

g_covar calculates and diagonalizes the covariance matrix
g_anaeig analyzes the eigenvectors
make.edi generate input files for essential dynamics sampling

Normal modes

grompp makes a run input file
mdrun finds a potential energy minimum
mdrun calculates the Hessian
g_nmeig diagonalizes the Hessian
g_nmtraj generate oscillating trajectory of an eigenmode
g_anaeig analyzes the normal modes
g_nmens generates an ensemble of structures from the normal modes

Chapter 8

Analysis

In this chapter different ways of analyzing your trajectory are described. The names of the corresponding analysis programs are given. Specific info on the in- and output of these programs can be found in the on-line manual at www.gromacs.org. The output files are often produced as finished Grace/Xmgr graphs.

First in sec. 8.1 the group concept in analysis is explained. Then the different analysis tools are presented.

8.1 Groups in Analysis.

`make_ndx`, `mk_angndx`

In chapter 3 it was explained how *groups of atoms* can be used in the MD-program. In most analysis programs groups of atoms are needed to work on. Most programs can generate several default index groups, but groups can always be read from an index file. Let's consider a simulation of a binary mixture of components A and B. When we want to calculate the radial distribution function (rdf) $g_{AB}(r)$ of A with respect to B, we have to calculate

$$4\pi r^2 g_{AB}(r) = V \sum_{i \in A} \sum_{j \in B} P(r) \quad (8.1)$$

where V is the volume and $P(r)$ is the probability to find a B atom at a distance r from an A atom.

By having the user define the *atom numbers* for groups A and B in a simple file we can calculate this g_{AB} in the most general way, without having to make any assumptions in the rdf-program about the type of particles.

Groups can therefore consist of a series of *atom numbers*, but in some cases also of *molecule numbers*. It is also possible to specify a series of angles by *triples* of *atom numbers*, dihedrals by *quadruples* of *atom numbers* and bonds or vectors (in a molecule) by *pairs* of *atom numbers*. When appropriate the type of index file will be specified for the following analysis programs. To help creating such index files (`index.ndx`), there are a couple of programs to generate them,

using either your input configuration or the topology. To generate an index file consisting of a series of *atom numbers* (as in the example of g_{AB}) use `make_ndx`. To generate an index file with angles or dihedrals, use `mk_angndx`. Of course you can also make them by hand. The general format is presented here:

```
[ Oxygen ]  
1 4 7  
[ Hydrogen ]  
2 3 5 6  
8 9
```

First the group name is written between square brackets. The following atom numbers may be spread out over as many lines as you like. The atom numbering starts at 1.

8.1.1 Default Groups

When no index file is supplied to analysis tools or `grompp`, a number of default groups are generated to choose from:

```
System  
    all atoms in the system  
  
Protein  
    all protein atoms  
  
Protein-H  
    protein atoms excluding hydrogens  
  
C-alpha  
    C $_{\alpha}$  atoms  
  
Backbone  
    protein backbone atoms; N, C $_{\alpha}$  and C  
  
MainChain  
    protein main chain atoms: N, C $_{\alpha}$ , C and O, including oxygens in C-terminus  
  
MainChain+Cb  
    protein main chain atoms including C $_{\beta}$   
  
MainChain+H  
    protein main chain atoms including backbone amide hydrogen and hydrogens on the N-terminus  
  
SideChain  
    protein side chain atoms; that is all atoms except N, C $_{\alpha}$ , C, O, backbone amide hydrogen, oxygens in C-terminus and hydrogens on the N-terminus  
  
SideChain-H  
    protein side chain atoms excluding all hydrogens
```

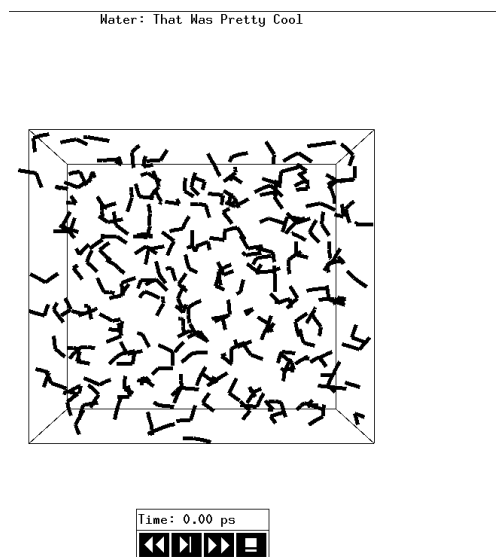


Figure 8.1: The window of ngmx showing a box of water.

Prot-Masses

protein atoms excluding dummy masses (as used in virtual site constructions of NH_3 groups and Tryptophane sidechains), see also sec. 5.2.2; this group is only included when it differs from the 'Protein' group

Non-Protein

all non-protein atoms

DNA

all DNA atoms

molecule_name

for all residues/molecules which are not recognized as protein or DNA, one group per residue/molecule name is generated

Other

all atoms which are neither protein nor DNA.

Empty groups will not be generated. Most of the groups only contain protein atoms. An atom is considered a protein atom if its residue name is listed in the `aminoacids.dat` file.

8.2 Looking at your trajectory

ngmx

Before analyzing your trajectory it is often informative to look at your trajectory first. Gromacs comes with a simple trajectory viewer `ngmx`; the advantage with this one is that it does not require OpenGL, which usually isn't present e.g. on supercomputers. It is also possible to generate a hard-copy in Encapsulated Postscript format, see Fig. 8.1. If you want a faster and more fancy viewer there are several programs that can read the GROMACS trajectory formats – have a look at our homepage www.gromacs.org for updated links.

8.3 General properties

`g_energy`, `g_traj`

To analyze some or all *energies* and other properties, such as *total pressure*, *pressure tensor*, *density*, *box-volume* and *box-sizes*, use the program `g_energy`. A choice can be made from a list a set of energies, like potential, kinetic or total energy, or individual contributions, like Lennard-Jones or dihedral energies.

The *center-of-mass velocity*, defined as

$$\mathbf{v}_{com} = \frac{1}{M} \sum_{i=1}^N m_i \mathbf{v}_i \quad (8.2)$$

with $M = \sum_{i=1}^N m_i$ the total mass of the system, can be monitored in time by the program `g_com`. It is however recommended to remove the center-of-mass velocity every step (see chapter 3)!

8.4 Radial distribution functions

`g_rdf`

The *radial distribution function* (rdf) or pair correlation function $g_{AB}(r)$ between particles of type A and B is defined in the following way:

$$\begin{aligned} g_{AB}(r) &= \frac{\langle \rho_B(r) \rangle}{\langle \rho_B \rangle_{local}} \\ &= \frac{1}{\langle \rho_B \rangle_{local}} \frac{1}{N_A} \sum_{i \in A} \sum_{j \in B} \frac{\delta(r_{ij} - r)}{4\pi r^2} \end{aligned} \quad (8.3)$$

with $\langle \rho_B(r) \rangle$ the particle density of type B at a distance r around particles A , and $\langle \rho_B \rangle_{local}$ the particle density of type B averaged over all spheres around particles A with radius r_{max} (see Fig. 8.2C).

Usually the value of r_{max} is half of the box length. The averaging is also performed in time. In practice the analysis program `g_rdf` divides the system into spherical slices (from r to $r + dr$, see Fig. 8.2A) and makes a histogram in stead of the δ -function. An example of the rdf of Oxygen-Oxygen in SPC-water [55] is given in Fig. 8.3.

With `g_rdf` it is also possible to calculate an angle dependent rdf $g_{AB}(r, \theta)$, where the angle θ is defined with respect to a certain laboratory axis \mathbf{e} , see Fig. 8.2B.

$$g_{AB}(r, \theta) = \frac{1}{\langle \rho_B \rangle_{local, \theta}} \frac{1}{N_A} \sum_{i \in A} \sum_{j \in B} \frac{\delta(r_{ij} - r) \delta(\theta_{ij} - \theta)}{2\pi r^2 \sin(\theta)} \quad (8.4)$$

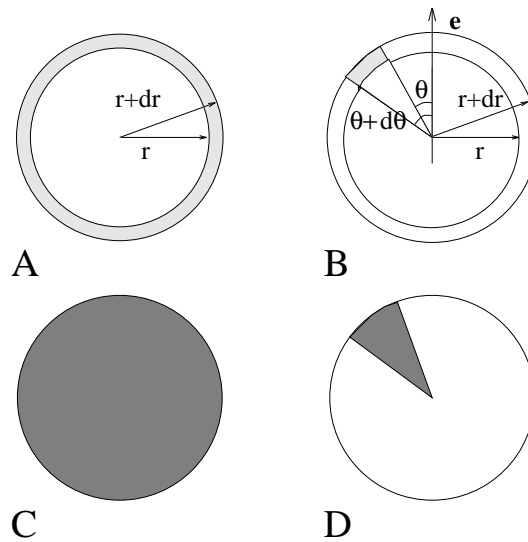


Figure 8.2: Definition of slices in g -rdf: A. $g_{AB}(r)$. B. $g_{AB}(r, \theta)$. The slices are colored grey. C. Normalization $\langle \rho_B \rangle_{local}$. D. Normalization $\langle \rho_B \rangle_{local, \theta}$. Normalization volumes are colored grey.

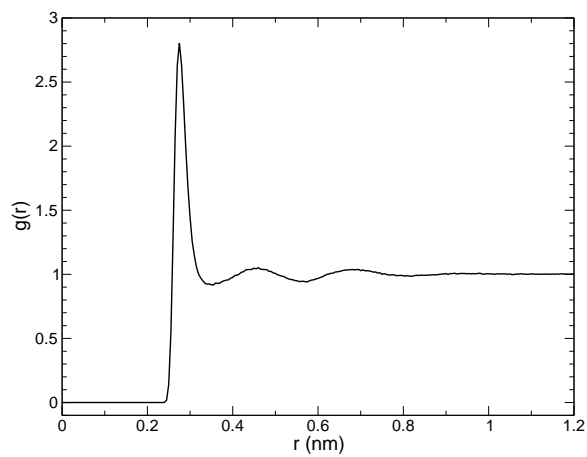


Figure 8.3: $g_{OO}(r)$ for Oxygen-Oxygen of SPC-water.

$$\cos(\theta_{ij}) = \frac{\mathbf{r}_{ij} \cdot \mathbf{e}}{\|\mathbf{r}_{ij}\| \|\mathbf{e}\|} \quad (8.5)$$

This $g_{AB}(r, \theta)$ is useful for analyzing anisotropic systems. Note that in this case the normalization $\langle \rho_B \rangle_{local, \theta}$ is the average density in all angle slices from θ to $\theta + d\theta$ up to r_{max} , so angle dependent, see Fig. 8.2D.

8.5 Correlation functions

8.5.1 Theory of correlation functions

The theory of correlation functions is well established [70]. However we want to describe here the implementation of the various correlation function flavors in the GROMACS code. The definition of the autocorrelation function (ACF) $C_f(t)$ for a property $f(t)$ is

$$C_f(t) = \langle f(\xi)f(\xi + t) \rangle_\xi \quad (8.6)$$

where the notation on the right hand side means averaging over ξ , *i.e.* over time origins. It is also possible to compute cross-correlation function from two properties $f(t)$ and $g(t)$:

$$C_{fg}(t) = \langle f(\xi)g(\xi + t) \rangle_\xi \quad (8.7)$$

however, in GROMACS there is no standard mechanism to do this (**note:** you can use the `xmgr` program to compute cross correlations). The integral of the correlation function over time is the correlation time τ_f :

$$\tau_f = \int_0^\infty C_f(t) dt \quad (8.8)$$

In practice correlation functions are calculated based on data points with discrete time intervals Δt , so that the ACF from an MD simulation is:

$$C_f(j\Delta t) = \frac{1}{N-j} \sum_{i=0}^{N-1-j} f(i\Delta t)f((i+j)\Delta t) \quad (8.9)$$

where N is the number of available time frames for the calculation. The resulting ACF is obviously only available at time points with the same interval Δt . Since for many applications it is necessary to know the short time behavior of the ACF (*e.g.* the first 10 ps) this often means that we have to save the atomic coordinates with short intervals. Another implication of eqn. 8.9 is that in principle we can not compute all points of the ACF with the same accuracy, since we have $N-1$ data points for $C_f(\Delta t)$ but only 1 for $C_f((N-1)\Delta t)$. However, if we decide to compute only an ACF of length $M\Delta t$, where $M \leq N/2$ we can compute all points with the same statistical accuracy:

$$C_f(j\Delta t) = \frac{1}{M} \sum_{i=0}^{N-1-M} f(i\Delta t)f((i+j)\Delta t) \quad (8.10)$$

here of course $j < M$. M is sometimes referred to as the time lag of the correlation function. When we decide to do this, we intentionally do not use all the available points for very short time intervals ($j \ll M$), but it makes it easier to interpret the results. Another aspect that may not be

neglected when computing ACFs from simulation, is that usually the time origins ξ (eqn. 8.6) are not statistically independent, which may introduce a bias in the results. This can be tested using a block-averaging procedure, where only time origins with a spacing at least the length of the time lag are included, *e.g.* using k time origins with spacing of $M\Delta t$ (where $kM \leq N$):

$$C_f(j\Delta t) = \frac{1}{k} \sum_{i=0}^{k-1} f(iM\Delta t) f((iM+j)\Delta t) \quad (8.11)$$

However, one needs very long simulations to get good accuracy this way, because there are many fewer points that contribute to the ACF.

8.5.2 Using FFT for computation of the ACF

The computational cost for calculating an ACF according to eqn. 8.9 is proportional to N^2 , which is considerable. However, this can be improved by using fast Fourier transforms to do the convolution [70].

8.5.3 Special forms of the ACF

There are some important varieties on the ACF, *e.g.* the ACF of a vector \mathbf{p} :

$$C_{\mathbf{p}}(t) = \int_0^\infty P_n(\cos \angle(\mathbf{p}(\xi), \mathbf{p}(\xi+t))) d\xi \quad (8.12)$$

where $P_n(x)$ is the n^{th} order Legendre polynomial¹. Such correlation times can actually be obtained experimentally using *e.g.* NMR or other relaxation experiments. GROMACS can compute correlations using the 1st and 2nd order Legendre polynomial (eqn. 8.12). This can a.o. be used for rotational autocorrelation (`g_rotacf`), dipole autocorrelation (`g_dipoles`).

In order to study torsion angle dynamics we define a dihedral autocorrelation function as [88]:

$$C(t) = \langle \cos(\theta(\tau) - \theta(\tau+t)) \rangle_\tau \quad (8.13)$$

Note that this is not a product of two functions as is generally used for correlation functions, but it may be rewritten as the sum of two products:

$$C(t) = \langle \cos(\theta(\tau)) \cos(\theta(\tau+t)) + \sin(\theta(\tau)) \sin(\theta(\tau+t)) \rangle_\tau \quad (8.14)$$

8.5.4 Some Applications

The program `g_velacc` calculates this *Velocity Auto Correlation Function*.

$$C_{\mathbf{v}}(\tau) = \langle \mathbf{v}_i(\tau) \cdot \mathbf{v}_i(0) \rangle_{i \in A} \quad (8.15)$$

The self diffusion coefficient can be calculated using the Green-Kubo relation [70]

$$D_A = \frac{1}{3} \int_0^\infty \langle \mathbf{v}_i(t) \cdot \mathbf{v}_i(0) \rangle_{i \in A} dt \quad (8.16)$$

¹ $P_0(x) = 1, P_1(x) = x, P_2(x) = (3x^2 - 1)/2$

which is just the integral of the velocity autocorrelation function. There is a widely held belief that the velocity ACF converges faster than the mean square displacement (sec. 8.6), which can also be used for the computation of diffusion constants. However, Allen & Tildesly [70] warn us that the long time contribution to the velocity ACF can not be ignored, so care must be taken.

Another important quantity is the dipole correlation time. The *dipole correlation function* for particles A is calculated as follows by `g_dipoles`:

$$C_{\mu}(\tau) = \langle \mu_i(\tau) \cdot \mu_i(0) \rangle_{i \in A} \quad (8.17)$$

with $\mu_i = \sum_{j \in i} \mathbf{r}_j q_j$. The dipole correlation time can be computed using eqn. 8.8. For some applications see [89].

The viscosity of a liquid can be related to the correlation time of the Pressure tensor P [90, 91]. `g_energy` can compute the viscosity, but this is not very accurate [80] (actually the values do not converge...).

8.6 Mean Square Displacement

`g_msd`

To determine the self diffusion coefficient D_A of particles A one can use the Einstein relation [70]

$$\lim_{t \rightarrow \infty} \langle \|\mathbf{r}_i(t) - \mathbf{r}_i(0)\|^2 \rangle_{i \in A} = 6D_A t \quad (8.18)$$

This *Mean Square Displacement* and D_A are calculated by the program `g_msd`. Normally an index file containing atom numbers is used and the MSD is averaged over atoms. For molecules consisting of more than one atom, \mathbf{r}_i can be taken as the center of mass positions of the molecules. In that case you should use an index file with molecule numbers. The results will be nearly identical to averaging over atoms, however. The `g_msd` program can also be used for calculating diffusion in one or two dimensions. This is useful for studying lateral diffusion on interfaces.

An example of the mean square displacement of SPC-water is given in Fig. 8.4.

8.7 Bonds, angles and dihedrals

`g_bond`, `g_angle`, `g_sgangle`

To monitor specific *bonds* in your molecules during time, the program `g_bond` calculates the distribution of the bond length in time. The index file consists of pairs of atom numbers, for example

```
[ bonds_1 ]
1 2
3 4
9 10
[ bonds_2 ]
12 13
```

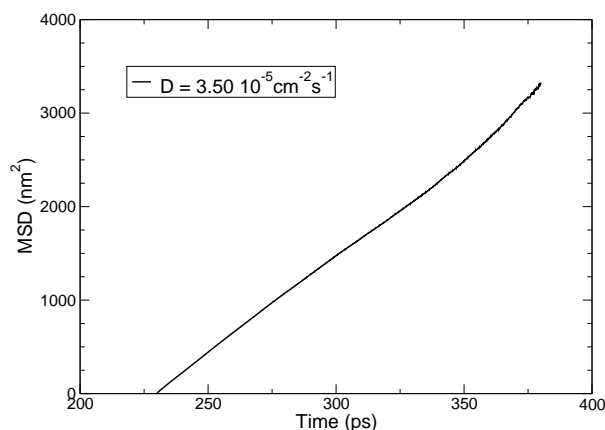


Figure 8.4: Mean Square Displacement of SPC-water.

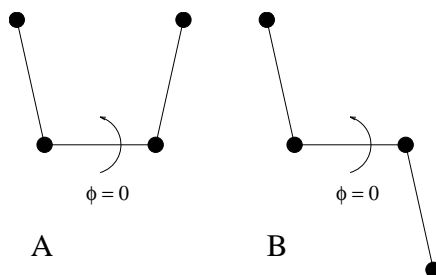


Figure 8.5: Dihedral conventions: A. “Biochemical convention”. B. “Polymer convention”.

The program `g_angle` calculates the distribution of *angles* and *dihedrals* in time. It also gives the average angle or dihedral. The index file consists of triplets or quadruples of atom numbers:

```
[ angles ]
1 2 3
2 3 4
3 4 5
[ dihedrals ]
1 2 3 4
2 3 5 5
```

For the dihedral angles you can use either the “biochemical convention” ($\phi = 0 \equiv cis$) or “polymer convention” ($\phi = 0 \equiv trans$), see Fig. 8.5.

To follow specific *angles* in time between two vectors, a vector and a plane or two planes (defined by 2, resp. 3 atoms inside your molecule, see Fig. 8.6A, B, C), use the program `g_sgangle`.

For planes it uses the normal vector perpendicular to the plane. It can also calculate the *distance* d between the geometrical center of two planes (see Fig. 8.6D), and the distances d_1 and d_2 between 2 atoms (of a vector) and the center of a plane defined by 3 atoms (see Fig. 8.6D). It further calculates the distance d between the center of the plane and the middle of this vector. Depending

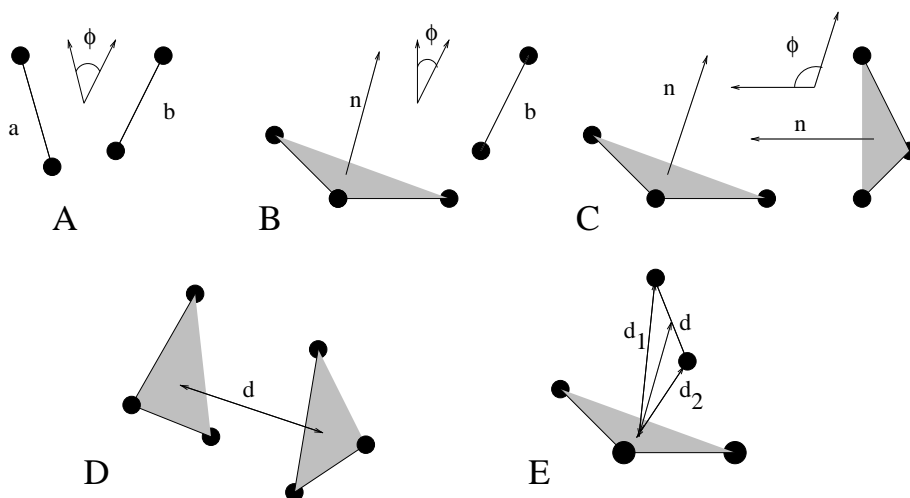


Figure 8.6: Options of `g_sgangle`: A. Angle between 2 vectors. B. Angle between a vector and the normal of a plane. C. Angle between two planes. D. Distance between the geometrical centers of 2 planes. E. Distances between a vector and the center of a plane.

on the input groups (*i.e.* groups of 2 or 3 atom numbers), the program decides what angles and distances to calculate. For example, the index-file could look like this:

```
[ a_plane ]
1 2 3
[ a_vector ]
3 4 5
```

8.8 Radius of gyration and distances

`g_gyrate`, `g_sgangle`, `g_mindist`, `g_mdmat`, `xpm2ps`

To have a rough measure for the compactness of a structure, you can calculate the *radius of gyration* with the program `g_gyrate` as follows:

$$R_g = \left(\frac{\sum_i \|\mathbf{r}_i\|^2 m_i}{\sum_i m_i} \right)^{\frac{1}{2}} \quad (8.19)$$

where m_i is the mass of atom i and \mathbf{r}_i the position of atom i with respect to the center of mass of the molecule. It is especially useful to characterize polymer solutions and proteins.

Sometimes it is interesting to plot the *distance* between two atoms, or the *minimum* distance between two groups of atoms (*e.g.*: protein side-chains in a salt bridge). To calculate these distances between certain groups there are several possibilities:

- The *distance between the geometrical centers* of two groups can be calculated with the program `g_sgangle`, as explained in sec. 8.7.

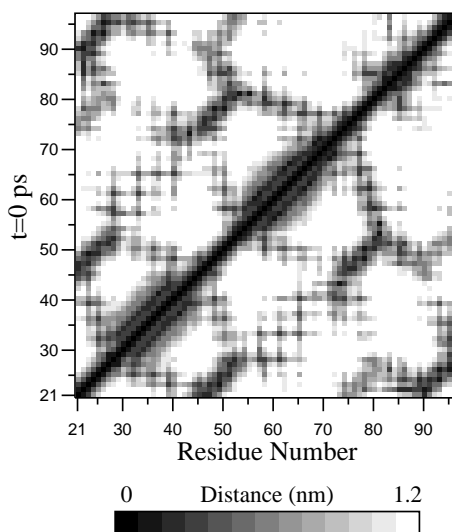


Figure 8.7: A minimum distance matrix for a peptide [92].

- The *minimum distance* between two groups of atoms during time can be calculated with the program `g_mindist`. It also calculates the *number of contacts* between these groups within a certain radius r_{max} .
- To monitor the *minimum distances between amino-acid residues* within a (protein) molecule, you can use the program `g_mdmat`. This minimum distance between two residues A_i and A_j is defined as the smallest distance between any pair of atoms ($i \in A_i, j \in A_j$). The output is a symmetrical matrix of smallest distances between all residues. To visualize this matrix, you can use a program such as `xv`. If you want to view the axes and legend or if you want to print the matrix, you can convert it with `xpm2ps` into a Postscript picture, see Fig. 8.7.
Plotting these matrices for different time-frames, one can analyze changes in the structure, and *e.g.* forming of salt bridges.

8.9 Root mean square deviations in structure

`g_rms`, `g_rmsdist`

The *root mean square deviation (RMSD)* of certain atoms in a molecule with respect to a reference structure can be calculated with the program `g_rms` by least-square fitting the structure to the reference structure ($t_2 = 0$) and subsequently calculating the *RMSD* (eqn. 8.20).

$$RMSD(t_1, t_2) = \left[\frac{1}{M} \sum_{i=1}^N m_i \|\mathbf{r}_i(t_1) - \mathbf{r}_i(t_2)\|^2 \right]^{\frac{1}{2}} \quad (8.20)$$

where $M = \sum_{i=1}^N m_i$ and $\mathbf{r}_i(t)$ is the position of atom i at time t . **NOTE** that fitting does not have to use the same atoms as the calculation of the *RMSD*; *e.g.*: a protein is usually fitted on the backbone atoms (N,C $_{\alpha}$,C), but the *RMSD* can be computed of the backbone or of the whole protein.

Instead of comparing the structures to the initial structure at time $t = 0$ (so for example a crystal structure), one can also calculate eqn. 8.20 with a structure at time $t_2 = t_1 - \tau$. This gives some insight in the mobility as a function of τ . Also a matrix can be made with the *RMSD* as a function of t_1 and t_2 , this gives a nice graphical impression of a trajectory. If there are transitions in a trajectory, they will clearly show up in such a matrix.

Alternatively the *RMSD* can be computed using a fit-free method with the program `g_rmsdist`:

$$RMSD(t) = \left[\frac{1}{N^2} \sum_{i=1}^N \sum_{j=1}^N \|\mathbf{r}_{ij}(t) - \mathbf{r}_{ij}(0)\|^2 \right]^{\frac{1}{2}} \quad (8.21)$$

where the *distance* \mathbf{r}_{ij} between atoms at time t is compared with the distance between the same atoms at time 0.

8.10 Covariance analysis

Covariance analysis, also called principal component analysis or essential dynamics [93], can find correlated motions. It uses the covariance matrix C of the atomic coordinates:

$$C_{ij} = \left\langle M_{ii}^{\frac{1}{2}} (x_i - \langle x_i \rangle) M_{jj}^{\frac{1}{2}} (x_j - \langle x_j \rangle) \right\rangle \quad (8.22)$$

where M is a diagonal matrix containing the masses of the atoms (mass-weighted analysis) or the unit matrix (non-mass weighted analysis). C is a symmetric $3N \times 3N$ matrix, which can be diagonalized with an orthonormal transformation matrix R :

$$R^T C R = \text{diag}(\lambda_1, \lambda_2, \dots, \lambda_{3N}) \quad \text{where } \lambda_1 \geq \lambda_2 \geq \dots \geq \lambda_{3N} \quad (8.23)$$

The columns of R are the eigenvectors, also called principal or essential modes. R defines a transformation to a new coordinate system. The trajectory can be projected on the principal modes to give the principal components $p_i(t)$:

$$\mathbf{p}(t) = R^T M^{\frac{1}{2}} (\mathbf{x}(t) - \langle \mathbf{x} \rangle) \quad (8.24)$$

The eigenvalue λ_i is the mean square fluctuation of principal component i . The first few principal modes often describe collective, global motions in the system. The trajectory can be filtered along one (or more) principal modes. For one principal mode i this goes as follows:

$$\mathbf{x}^f(t) = \langle \mathbf{x} \rangle + M^{-\frac{1}{2}} R_{*i} p_i(t) \quad (8.25)$$

When the analysis is performed on a macromolecule, one often wants to remove the overall rotation and translation to look at the internal motion only. This can be achieved by least square fitting to a reference structure. Care has to be taken that the reference structure is representative for the ensemble, since the choice of reference structure influences the covariance matrix.

One should always check if the principal modes are well defined. If the first principal component resembles a half cosine and the second resembles a full cosine, you might be filtering noise (see below). A good way to check the relevance of the first few principal modes is to calculate the

overlap of the sampling between the first and second half of the simulation. Note that this can only be done when the same reference structure is used for the two halves.

A good measure for the overlap has been defined in [94]. The elements of the covariance matrix are proportional to the square of the displacement, so we need to take the square root of the matrix to examine the extent of sampling. The square root can be calculated from the eigenvalues λ_i and the eigenvectors, which are the columns of the rotation matrix R . For a symmetric and diagonally-dominant matrix A of size $3N \times 3N$ the square root can be calculated as:

$$A^{\frac{1}{2}} = R \text{diag}(\lambda_1^{\frac{1}{2}}, \lambda_2^{\frac{1}{2}}, \dots, \lambda_{3N}^{\frac{1}{2}}) R^T \quad (8.26)$$

It can be verified easily that the product of this matrix with itself gives A . Now we can define a difference d between covariance matrices A and B as follows:

$$d(A, B) = \sqrt{\text{tr} \left((A^{\frac{1}{2}} - B^{\frac{1}{2}})^2 \right)} \quad (8.27)$$

$$= \sqrt{\text{tr} \left(A + B - 2A^{\frac{1}{2}}B^{\frac{1}{2}} \right)} \quad (8.28)$$

$$= \left(\sum_{i=1}^N (\lambda_i^A + \lambda_i^B) - 2 \sum_{i=1}^N \sum_{j=1}^N \sqrt{\lambda_i^A \lambda_j^B} (R_i^A \cdot R_j^B)^2 \right)^{\frac{1}{2}} \quad (8.29)$$

where tr is the trace of a matrix. We can now define the overlap s as:

$$s(A, B) = 1 - \frac{d(A, B)}{\sqrt{\text{tr}A + \text{tr}B}} \quad (8.30)$$

The overlap is 1 if and only if matrices A and B are identical. It is 0 when the sampled subspaces are completely orthogonal.

A commonly used measure is the subspace overlap of the first few eigenvectors of covariance matrices. The overlap of the subspace spanned by m orthonormal vectors $\mathbf{w}_1, \dots, \mathbf{w}_m$ with a reference subspace spanned by n orthonormal vectors $\mathbf{v}_1, \dots, \mathbf{v}_n$ can be quantified as follows:

$$\text{overlap}(\mathbf{v}, \mathbf{w}) = \frac{1}{n} \sum_{i=1}^n \sum_{j=1}^m (\mathbf{v}_i \cdot \mathbf{w}_j)^2 \quad (8.31)$$

The overlap will increase with increasing m and will be 1 when set \mathbf{v} is a subspace of set \mathbf{w} . The disadvantage of this method is that it does not take the eigenvalues into account. All eigenvectors are weighted equally and when degenerate subspaces are present (equal eigenvalues) the calculated overlap will be too low.

Another useful check is the cosine content. It has been proven the the principal components of random diffusion are cosines with the number of periods equal to half the principal component index[95, 94]. The eigenvalues are proportional to the index to the power -2 . The cosine content is defined as:

$$\frac{2}{T} \left(\int_0^T \cos(i\pi t) p_i(t) dt \right)^2 \left(\int_0^T p_i^2(t) dt \right)^{-1} \quad (8.32)$$

When the cosine content of the first few principal components is close to 1, the largest fluctuations are not connected with the potential, but with random diffusion.

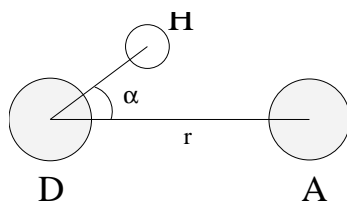


Figure 8.8: Geometrical Hydrogen bond criterion.

The covariance matrix is built and diagonalized by `g_covar`. The principal components and overlap (any many more things) can be plotted and analyzed with `g_anaeig`. The cosine content can be calculated with `g_analyze`.

8.11 Dihedral principal component analysis

`g_angle`, `g_covar`, `g_anaeig`

Principal component analysis can be performed in dihedral space [96] using GROMACS. You start by defining the dihedral angles of your interest in an index file, either using `mk_angndx` or otherwise. Then you use the `g_angle` program with the `-or` flag to produce a new `trr` file containing the cosine and sine of each dihedral angle in two coordinates respectively. That is, in the `trr` file you will have a series of numbers corresponding to: $\cos(\phi_1)$, $\sin(\phi_1)$, $\cos(\phi_2)$, $\sin(\phi_2)$, ..., $\cos(\phi_n)$, $\sin(\phi_n)$, the array is padded with zeros if necessary. Then you can use this `trr` file as input for the `g_covar` program and perform principal component analysis as usual. For this work you will need to generate a reference file (`tpr`, `gro`, `pdb` etc.) containing the same number of “atoms” as the new `trr` file, that is for n dihedrals you need $2n/3$ atoms (rounded up if not an integer number). You should use the `-nofit` option for `g_covar` since the coordinates in the dummy reference file do not correspond in any way to the information in the `trr` file. Analysis of the results is done using `g_anaeig`.

8.12 Hydrogen bonds

`g_hbond`

The program `g_hbond` analyses the *hydrogen bonds* (H-bonds) between all possible donors D and acceptors A. To determine if an H-bond exists, a geometrical criterion is used, see also Fig. 8.8:

$$\begin{aligned} r &\leq r_{HB} = 0.35\text{nm} \\ \alpha &\leq \alpha_{HB} = 30^\circ \end{aligned} \quad (8.33)$$

The value of $r_{HB} = 0.35$ nm corresponds to the first minimum of the rdf of SPC-water (see also Fig. 8.3).

The program `g_hbond` analyses all hydrogen bonds existing between two groups of atoms (which must be either identical or non-overlapping) or in specified Donor Hydrogen Acceptor triplets, in the following ways:

- Donor-Acceptor distance (r) distribution of all H-bonds

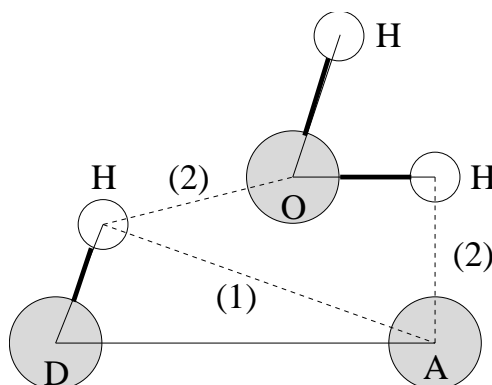


Figure 8.9: Insertion of water into an H-bond. (1) Normal H-bond between two residues. (2) H-bonding bridge via a water molecule.

- Hydrogen-Donor-Acceptor angle (α) distribution of all H-bonds
- The total number of H-bonds in each time frame
- The number of H-bonds in time between residues, divided into groups $n-n+i$ where n and $n+i$ stand for residue numbers and i goes from 0 to 6. The group for $i = 6$ also includes all H-bonds for $i > 6$. These groups include the $n-n+3$, $n-n+4$ and $n-n+5$ H-bonds which provide a measure for the formation of α -helices or β -turns or strands.
- The lifetime of the H-bonds is calculated from the average over all autocorrelation functions of the existence functions (either 0 or 1) of all H-bonds:

$$C(\tau) = \langle s_i(t) s_i(t + \tau) \rangle \quad (8.34)$$

with $s_i(t) = \{0, 1\}$ for H-bond i at time t . The integral of $C(\tau)$ gives a rough estimate of the average H-bond lifetime τ_{HB} :

$$\tau_{HB} = \int_0^{\infty} C(\tau) d\tau \quad (8.35)$$

Both the integral and the complete auto correlation function $C(\tau)$ will be output, so that more sophisticated analysis (*e.g.* using multi-exponential fits) can be used to get better estimates for τ_{HB} . A more complicate analysis is given in ref. [97], one of the more fancy option is the Luzar and Chandler analysis of hydrogen bond kinetics [98, 99].

- An H-bond existence map can be generated of dimensions $\# H\text{-bonds} \times \# \text{frames}$.
- Index groups are output containing the analyzed groups, all donor-hydrogen atom pairs and acceptor atoms in these groups, donor-hydrogen-acceptor triplets involved in hydrogen bonds between the analyzed groups and all solvent atoms involved in insertion.
- Solvent insertion into H-bonds can be analyzed, see Fig. 8.9. In this case an additional group identifying the solvent must be selected. The occurrence of insertion will be indicated in the existence map. Note that insertion into and existence of a specific H-bond can occur simultaneously and will also be indicated as such in the existence map.

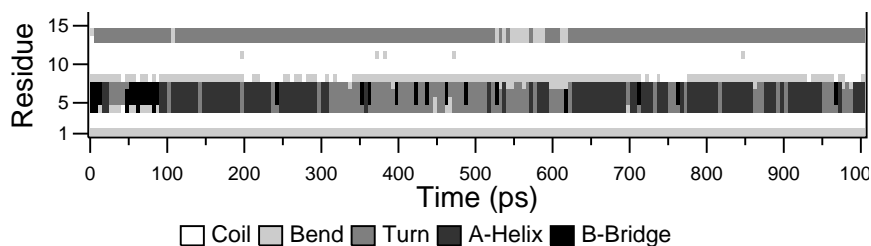


Figure 8.10: Analysis of the secondary structure elements of a peptide in time.

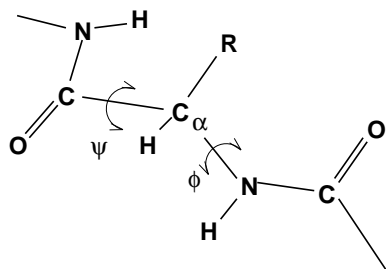


Figure 8.11: Definition of the dihedral angles ϕ and ψ of the protein backbone.

8.13 Protein related items

`do_dssp`, `g_rama`, `xrama`, `wheel`

To analyze structural changes of a protein, you can calculate the radius of gyration or the minimum residue distances during time (see sec. 8.8), or calculate the RMSD (sec. 8.9).

You can also look at the changing of *secondary structure elements* during your run. For this you can use the program `do_dssp`, which is an interface for the commercial program `dssp` [100]. For further information, see the `dssp`-manual. A typical output plot of `do_dssp` is given in Fig. 8.10.

One other important analysis of proteins is the so called *Ramachandran plot*. This is the projection of the structure on the two dihedral angles ϕ and ψ of the protein backbone, see Fig. 8.11.

To evaluate this Ramachandran plot you can use the program `g_rama`. A typical output is given in Fig. 8.12.

It is also possible to generate an animation of the Ramachandran plot in time. This can be of help for analyzing certain dihedral transitions in your protein. You can use the program `xrama` for this.

When studying α -helices it is useful to have a *helical wheel* projection of your peptide, to see whether a peptide is amphipatic. This can be done using the `wheel` program. Two examples are plotted in Fig. 8.13.

8.14 Interface related items

`g_order`, `g_density`, `g_potential`, `g_coord`

When simulating molecules with long carbon tails, it can be interesting to calculate their average

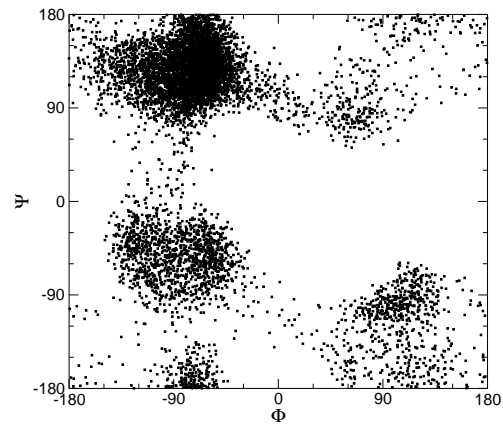


Figure 8.12: Ramachandran plot of a small protein.

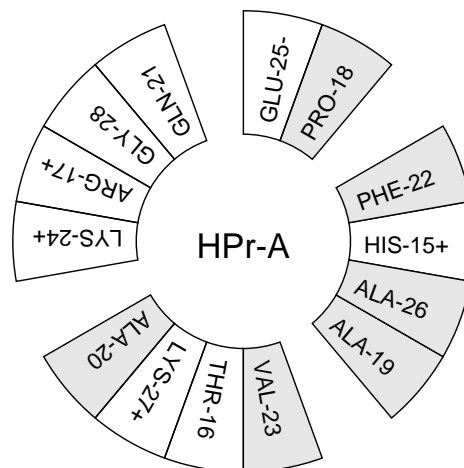


Figure 8.13: Helical wheel projection of the N-terminal helix of HPr.

orientation. There are several flavors of order parameters, most of which are related. The program `g_order` can calculate order parameters using the equation

$$S_z = \frac{3}{2} \langle \cos^2 \theta_z \rangle - \frac{1}{2} \quad (8.36)$$

where θ_z is the angle between the z -axis of the simulation box and the molecular axis under consideration. The latter is defined as the vector from C_{n-1} to C_{n+1} . The parameters S_x and S_y are defined in the same way. The brackets imply averaging over time and molecules. Order parameters can vary between 1 (full order along the interface normal) and $-1/2$ (full order perpendicular to the normal), with a value of zero in the case of isotropic orientation.

The program can do two things for you. It can calculate the order parameter for each CH_2 segment separately, for any of three axes, or it can divide the box in slices and calculate the average value of the order parameter per segment in one slice. The first method gives an idea of the ordering of a molecule from head to tail, the second method gives an idea of the ordering as function of the box length.

The electrostatic potential (ψ) across the interface can be computed from a trajectory by evaluating the double integral of the charge density ($\rho(z)$):

$$\psi(z) - \psi(-\infty) = - \int_{-\infty}^z dz' \int_{-\infty}^{z'} \rho(z'') dz'' / \epsilon_0 \quad (8.37)$$

where the position $z = -\infty$ is far enough in the bulk phase that the field is zero. With this method, it is possible to “split” the total potential into separate contributions from lipid and water molecules. The program `g_potential` divides the box in slices and sums all charges of the atoms in each slice. It then integrates this charge density, giving the electric field, and the electric field, giving the potential. Charge density, field and potential are written to `xvgr`-input files.

The program `g_coord` is a very simple analysis program. All it does is print the coordinates of selected atoms to three files, containing respectively the x -, y - and z -coordinates of those atoms. It can also calculate the center of mass of one or more molecules and print the coordinates of the center of mass to three files. By itself, this is probably not a very useful analysis, but having the coordinates of selected molecules or atoms can be very handy for further analysis, not only in interface systems.

The program `g_pvd` calculates a lot of properties, among which the density of a group in particles per unit of volume, but not a density that takes the mass of the atoms into account. The program `g_density` also calculates the density of a group, but takes the masses into account and gives a plot of the density against a box axis. This is useful for looking at the distribution of groups or atoms across the interface.

8.15 Chemical shifts

`total, do_shift`

You can compute the NMR chemical shifts of protons with the program `do_shift`. This is just an GROMACS interface to the public domain program `total` [101]. For further information, read the article. Although there is limited support for this in GROMACS users are encouraged to use the software provided by the David Case group at Scripps because it seems to be more up-to-date.

Appendix A

Technical Details

A.1 Installation

The entire GROMACS package is Free Software, licensed under the GNU General Public License. The main distribution site is our WWW server at www.gromacs.org.

The package is mainly distributed as source code, but we also provide RPM packages for Linux. On the home page you will find all the information you need to install the package, mailing lists with archives, and several additional online resources like contributed topologies, etc. The default installation action is simply to unpack the source code and the issue

```
./configure  
make  
make install
```

The configuration script should automatically determine the best options for your platform, and it will tell you if anything is missing on your system. You will also find detailed step-by-step installation instructions on the website.

A.2 Single or Double precision

GROMACS can be compiled in either single or double precision. The default choice is single precision, but it is easy to turn on double precision by selecting the `--disable-float` option to the configuration script. Double precision will be 0 to 50% slower than single precision depending on the architecture you are running on. Double precision will use somewhat more memory and run input, energy and full-precision trajectory files will be almost twice as large. Assembly loops are available in single and double precision on Pentium 4, Opteron and Itanium processors. On PowerPC processors containing the AltiVec unit only single precision is possible. On older Athlon and Pentium 3 processors only the single precision code is available, due to hardware limitations. All other processors use either C or Fortran code for the compute intensive inner loops.

The energies in single precision are accurate up to the last decimal, the last one or two decimals of the forces are non-significant. The virial is less accurate than the forces, since the virial is only one

order of magnitude larger than the size of each element in the sum over all atoms (sec. B.1). In most cases this is not really a problem, since the fluctuations in de virial can be 2 orders of magnitude larger than the average. In periodic charged systems these errors are often negligible. Especially cut-off's for the Coulomb interactions cause large errors in the energies, forces and virial. Even when using a reaction-field or lattice sum method the errors are larger than or comparable to the errors due to the single precision. Since MD is chaotic, trajectories with very similar starting conditions will diverge rapidly, the divergence is faster in single precision than in double precision. For most simulations single precision is accurate enough. In some cases double precision is required to get reasonable results:

- normal mode analysis, for the conjugate gradient or l-bfgs minimization and the calculation and diagonalization of the Hessian
- calculation of the constraint force between two large groups of atoms
- energy conservation (this can only be done without temperature coupling and without cut-off's)

A.3 Porting GROMACS

The GROMACS system is designed with portability as a major design goal. However there are a number of things we assume to be present on the system GROMACS is being ported on. We assume the following features:

1. A UNIX-like operating system (BSD 4.x or SYSTEM V rev.3 or higher) or UNIX-like libraries running under e.g. CygWin
2. an ANSI C compiler
3. optionally a Fortran-77 compiler or Fortran-90 compiler for faster (on some computers) inner loop routines
4. optionally the Nasm assembler to use the assembly innerloops on x86 processors.

There are some additional features in the package that require extra stuff to be present, but it is checked for in the configuration script and you will be warned if anything important is missing.

That's the requirements for a single processor system. If you want to compile GROMACS for a multiple processor environment you also need a MPI library (Message-Passing Interface) to perform the parallel communication. This is always shipped with supercomputers, and for workstations you can find links to free MPI implementations through the GROMACS homepage at www.gromacs.org.

A.3.1 Multi-processor Optimization

If you want to, you could write your own optimized communication (perhaps using specific libraries for your hardware) instead of MPI. This should never be necessary for normal use (we

haven't heard of a modern computer where it isn't possible to run MPI), but if you absolutely want to do it, here are some clues.

The interface between the communication routines and the rest of the GROMACS system is described in the file `$GMXHOME/src/include/network.h`. We will give a short description of the different routines below.

extern void gmx_tx(int pid, void *buf, int bufsize);

This routine, when called with the destination processor number, a pointer to a (byte oriented) transfer buffer, and the size of the buffer will send the buffer to the indicated processor (in our case always the neighboring processor). The routine does **not** wait until the transfer is finished.

extern void gmx_tx_wait(int pid);

This routine waits until the previous, or the ongoing transmission is finished.

extern void gmx_txs(int pid, void *buf, int bufsize);

This routine implements a synchronous send by calling the a-synchronous routine and then the wait. It might come in handy to code this differently.

extern void gmx_rx(int pid, void *buf, int bufsize);

extern void gmx_rx_wait(int pid);

extern void gmx_rxs(int pid, void *buf, int bufsize);

The very same routines for receiving a buffer and waiting until the reception is finished.

extern void gmx_init(int pid, int nprocs);

This routine initializes the different devices needed to do the communication. In general it sets up the communication hardware (if it is accessible) or does an initialize call to the lower level communication subsystem.

extern void gmx_stat(FILE *fp, char *msg);

With this routine we can diagnose the ongoing communication. In the current implementation it prints the various contents of the hardware communication registers of the (Intel i860) multiprocessor boards to a file.

A.4 Environment Variables

GROMACS programs may be influenced by the use of environment variables. First of all, the variables set in the `GMXRC` file are essential for running and compiling GROMACS. Other variables are:

1. `DUMPNL`, dump neighbor list. If set to a positive number the *entire* neighbor list is printed in the log file (may be many megabytes). Mainly for debugging purposes, but may also be handy for porting to other platforms.
2. `GMX_NO_QUOTES`, if this is explicitly set, no cool quotes will be printed at the end of a program

3. `WHERE`, when set print debugging info on line numbers.
4. `LOG_BUFS`, the size of the buffer for file I/O. When set to 0, all file I/O will be unbuffered and therefore very slow. This can be handy for debugging purposes, because it ensures that all files are always totally up-to-date.
5. `GMXNPRI`, for SGI systems only. When set, gives the default non-degrading priority (`npri`) for `mdrun`, `nmrn`, `g_covar` and `g_nmeig`, e.g. setting `setenv GMXNPRI 250` causes all runs to be performed at near-lowest priority by default.
6. `GMX_VIEW_XPM`, `GMX_VIEW_XVG`, `GMX_VIEW_EPS` and `GMX_VIEW_PDB`, commands used to automatically view resp. `.xvg`, `.xpm`, `.eps` and `.pdb` file types; they default to `xv`, `xmgrace`, `ghostview` and `rasmol`. Set to empty to disable automatic viewing of a particular file type. The command will be forked off and run in the background at the same priority as the GROMACS tool (which might not be what you want). Be careful not to use a command which blocks the terminal (e.g. `vi`), since multiple instances might be run.
7. `GMXTIMEUNIT` the time unit used in output files, can be anything in `fs`, `ps`, `ns`, `us`, `ms`, `s`, `m` or `h`.

Some other environment variables are specific to one program, such as `TOTAL` for the `do_shift` program, and `DSPP` for the `do_dspp` program.

A.5 Running GROMACS in parallel

If you have installed the MPI (Message Passing Interface) on your computer(s) you can compile GROMACS with this library to run simulations in parallel. All supercomputers are shipped with MPI libraries optimized for that particular platform, and if you are using a cluster of workstations there are several good free MPI implementations. You can find updated links to these on the [gromacs homepage](http://www.gromacs.org) www.gromacs.org. Once you have an MPI library installed it's trivial to compile GROMACS with MPI support: Just set the option `--enable-mpi` to the configure script and recompile. (But don't forget to make `distclean` before running `configure` if you have previously compiled with a different configuration.) If you are using a supercomputer you might also want to turn off the default nicing of the `mdrun` process with the `--disable-nice` option.

There is usually a program called `mpirun` with which you can fire up the parallel processes. A typical command line looks like:

```
% mpirun -p goofus,doofus,fred 10 mdrun -s topol -v -N 30
this runs on each of the machines goofus,doofus,fred with 10 processes on each1.
```

If you have a single machine with multiple processors you don't have to use the `mpirun` command, but you can do with an extra option to `mdrun`:

```
% mdrun -np 8 -s topol -v -N 8
```

In this example MPI reads the first option from the command line. Since `mdrun` also wants to know the number of processes you have to type it twice.

Check your local manuals (or online manual) for exact details of your MPI implementation.

¹Example taken from Silicon Graphics manual

If you are interested in programming MPI yourself, you can find manuals and reference literature on the internet.

Appendix B

Some implementation details

In this chapter we will present some implementation details. This is far from complete, but we deemed it necessary to clarify some things that would otherwise be hard to understand.

B.1 Single Sum Virial in GROMACS.

The virial Ξ can be written in full tensor form as:

$$\Xi = -\frac{1}{2} \sum_{i<j}^N \mathbf{r}_{ij} \otimes \mathbf{F}_{ij} \quad (\text{B.1})$$

where \otimes denotes the *direct product* of two vectors¹. When this is computed in the inner loop of an MD program 9 multiplications and 9 additions are needed².

Here it is shown how it is possible to extract the virial calculation from the inner loop [102].

B.1.1 Virial.

In a system with Periodic Boundary Conditions, the periodicity must be taken into account for the virial:

$$\Xi = -\frac{1}{2} \sum_{i<j}^N \mathbf{r}_{ij}^n \otimes \mathbf{F}_{ij} \quad (\text{B.2})$$

where \mathbf{r}_{ij}^n denotes the distance vector of the *nearest image* of atom i from atom j . In this definition we add a *shift vector* δ_i to the position vector \mathbf{r}_i of atom i . The difference vector \mathbf{r}_{ij}^n is thus equal to:

$$\mathbf{r}_{ij}^n = \mathbf{r}_i + \delta_i - \mathbf{r}_j \quad (\text{B.3})$$

or in shorthand:

$$\mathbf{r}_{ij}^n = \mathbf{r}_i^n - \mathbf{r}_j \quad (\text{B.4})$$

¹ $(\mathbf{u} \otimes \mathbf{v})^{\alpha\beta} = \mathbf{u}_\alpha \mathbf{v}_\beta$

²The calculation of Lennard-Jones and Coulomb forces is about 50 floating point operations.

In a triclinic system there are 27 possible images of i , when truncated octahedron is used there are 15 possible images.

B.1.2 Virial from non-bonded forces.

Here the derivation for the single sum virial in the *non-bonded force* routine is given. $i \neq j$ in all formulae below.

$$\Xi = -\frac{1}{2} \sum_{i < j}^N \mathbf{r}_{ij}^n \otimes \mathbf{F}_{ij} \quad (\text{B.5})$$

$$= -\frac{1}{4} \sum_{i=1}^N \sum_{j=1}^N (\mathbf{r}_i + \delta_i - \mathbf{r}_j) \otimes \mathbf{F}_{ij} \quad (\text{B.6})$$

$$= -\frac{1}{4} \sum_{i=1}^N \sum_{j=1}^N (\mathbf{r}_i + \delta_i) \otimes \mathbf{F}_{ij} - \mathbf{r}_j \otimes \mathbf{F}_{ij} \quad (\text{B.7})$$

$$= -\frac{1}{4} \left(\sum_{i=1}^N \sum_{j=1}^N (\mathbf{r}_i + \delta_i) \otimes \mathbf{F}_{ij} - \sum_{i=1}^N \sum_{j=1}^N \mathbf{r}_j \otimes \mathbf{F}_{ij} \right) \quad (\text{B.8})$$

$$= -\frac{1}{4} \left(\sum_{i=1}^N (\mathbf{r}_i + \delta_i) \otimes \sum_{j=1}^N \mathbf{F}_{ij} - \sum_{j=1}^N \mathbf{r}_j \otimes \sum_{i=1}^N \mathbf{F}_{ij} \right) \quad (\text{B.9})$$

$$= -\frac{1}{4} \left(\sum_{i=1}^N (\mathbf{r}_i + \delta_i) \otimes \mathbf{F}_i + \sum_{j=1}^N \mathbf{r}_j \otimes \mathbf{F}_j \right) \quad (\text{B.10})$$

$$= -\frac{1}{4} \left(2 \sum_{i=1}^N \mathbf{r}_i \otimes \mathbf{F}_i + \sum_{i=1}^N \delta_i \otimes \mathbf{F}_i \right) \quad (\text{B.11})$$

In these formulae we introduced

$$\mathbf{F}_i = \sum_{j=1}^N \mathbf{F}_{ij} \quad (\text{B.12})$$

$$\mathbf{F}_j = \sum_{i=1}^N \mathbf{F}_{ji} \quad (\text{B.13})$$

which is the total force on i resp. j . Because we use Newton's third law

$$\mathbf{F}_{ij} = -\mathbf{F}_{ji} \quad (\text{B.14})$$

we must in the implementation double the term containing the shift δ_i .

B.1.3 The intramolecular shift (mol-shift).

For the bonded-forces and shake it is possible to make a *mol-shift* list, in which the periodicity is stored. We simply have an array `mshift` in which for each atom an index in the `shiftvec` array is stored.

The algorithm to generate such a list can be derived from graph theory, considering each particle in a molecule as a bead in a graph, the bonds as edges.

- 1 represent the bonds and atoms as bidirectional graph
- 2 make all atoms white
- 3 make one of the white atoms black (atom i) and put it in the central box
- 4 make all of the neighbors of i that are currently white, grey
- 5 pick one of the grey atoms (atom j), give it the correct periodicity with respect to any of its black neighbors and make it black
- 6 make all of the neighbors of j that are currently white, grey
- 7 if any grey atom remains, goto [5]
- 8 if any white atom remains, goto [3]

Using this algorithm we can

- optimize the bonded force calculation as well as shake
- calculate the virial from the bonded forces in the single sum way again

Find a representation of the bonds as a bidirectional graph.

B.1.4 Virial from Covalent Bonds.

The covalent bond force gives a contribution to the virial, we have

$$b = \|\mathbf{r}_{ij}^n\| \quad (\text{B.15})$$

$$V_b = \frac{1}{2}k_b(b - b_0)^2 \quad (\text{B.16})$$

$$\mathbf{F}_i = -\nabla V_b \quad (\text{B.17})$$

$$= k_b(b - b_0)\frac{\mathbf{r}_{ij}^n}{b} \quad (\text{B.18})$$

$$\mathbf{F}_j = -\mathbf{F}_i \quad (\text{B.19})$$

The virial contribution from the bonds then is

$$\Xi_b = -\frac{1}{2}(\mathbf{r}_i^n \otimes \mathbf{F}_i + \mathbf{r}_j \otimes \mathbf{F}_j) \quad (\text{B.20})$$

$$= -\frac{1}{2}\mathbf{r}_{ij}^n \otimes \mathbf{F}_i \quad (\text{B.21})$$

B.1.5 Virial from Shake.

An important contribution to the virial comes from shake. Satisfying the constraints a force \mathbf{G} is exerted on the particles shaken. If this force does not come out of the algorithm (as in standard shake) it can be calculated afterwards (when using *leap-frog*) by:

$$\Delta \mathbf{r}_i = \mathbf{r}_i(t + \Delta t) - [\mathbf{r}_i(t) + \mathbf{v}_i(t - \frac{\Delta t}{2})\Delta t + \frac{\mathbf{F}_i}{m_i}\Delta t^2] \quad (\text{B.22})$$

$$\mathbf{G}_i = \frac{m_i \Delta \mathbf{r}_i}{\Delta t^2} \quad (\text{B.23})$$

but this does not help us in the general case. Only when no periodicity is needed (like in rigid water) this can be used, otherwise we must add the virial calculation in the inner loop of shake.

When it *is* applicable the virial can be calculated in the single sum way:

$$\Xi = -\frac{1}{2} \sum_i^{N_c} \mathbf{r}_i \otimes \mathbf{F}_i \quad (\text{B.24})$$

where N_c is the number of constrained atoms.

B.2 Optimizations

Here we describe some of the algorithmic optimizations used in GROMACS, apart from parallelism. One of these, the implementation of the $1.0/\sqrt{x}$ function is treated separately in sec. B.3. The most important other optimizations are described below.

B.2.1 Inner Loops for Water

GROMACS users special inner loop to calculate non-bonded interactions for water molecules with other atoms, and yet another set of loops for interactions between pairs of water molecules. This very optimized loop assumes a water model similar to SPC [55], *i.e.*:

1. There are three atoms in the molecule.
2. The first atom has Lennard-Jones (sec. 4.1.1) and coulomb (sec. 4.1.3) interactions.
3. Atoms two and three have only coulomb interactions, and equal charges.

The loop also works for the SPC/E [103] and TIP3P [48] water models. For more complicated molecules there is a general solvent loop assuming (note the order):

1. At the beginning of the molecule topology there is an arbitrary number of atoms with Lennard-Jones and coulomb interactions.
2. Then we have an arbitrary number of atoms with coulomb interactions only.
3. And finally there can be an arbitrary number of atoms with Lennard-Jones interactions only.

Note that this loop provides much less optimization than the water loop, but it is slightly better than the default routine.

The gain of these implementations is that there are more floating point operations in a single loop, which implies that some compilers can schedule the code better. However, it turns out that even some of the most advanced compilers have problems with scheduling, implying that manual tweaking is necessary to get optimum performance. This may include common-subexpression elimination, or moving code around.

B.2.2 Fortran Code

Unfortunately, Fortran compilers are still better than C-compilers, for most machines anyway. For some machines (*e.g.* SGI Power Challenge) the difference may be up to a factor of 3, in the case of vector computers this may be even larger. Therefore, some of the routines that take up a lot of computer time have been translated into Fortran and even assembly code for Intel and AMD x86 processors. In most cases, the Fortran or assembly loops should be selected automatically by the configure script when appropriate, but you can also tweak this by setting options to the configure script.

B.3 Computation of the 1.0/sqrt function.

B.3.1 Introduction.

The GROMACS project started with the development of a $1/\sqrt{x}$ processor which calculates

$$Y(x) = \frac{1}{\sqrt{x}} \quad (\text{B.25})$$

As the project continued, the Intel i860 processor was used to implement GROMACS, which now turned into almost a full software project. The $1/\sqrt{x}$ processor was implemented using a Newton-Raphson iteration scheme for one step. For this it needed lookup tables to provide the initial approximation. The $1/\sqrt{x}$ function makes it possible to use two almost independent tables for the exponent seed and the fraction seed with the IEEE floating point representation.

B.3.2 General

According to [104] the $1/\sqrt{x}$ can be calculated using the Newton-Raphson iteration scheme. The inverse function is

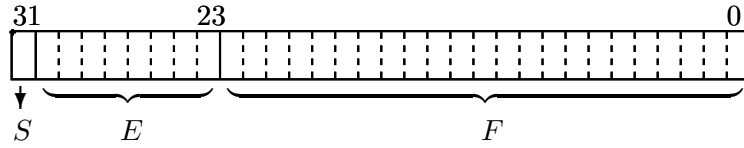
$$X(y) = \frac{1}{y^2} \quad (\text{B.26})$$

So instead of calculating

$$Y(a) = q \quad (\text{B.27})$$

the equation

$$X(q) - a = 0 \quad (\text{B.28})$$



$$Value = (-1)^S (2^{E-127}) (1.F)$$

Figure B.1: IEEE single precision floating point format

can now be solved using Newton-Raphson. An iteration is performed by calculating

$$y_{n+1} = y_n - \frac{f(y_n)}{f'(y_n)} \quad (\text{B.29})$$

The absolute error ε , in this approximation is defined by

$$\varepsilon \equiv y_n - q \quad (\text{B.30})$$

using Taylor series expansion to estimate the error results in

$$\varepsilon_{n+1} = -\frac{\varepsilon_n^2 f''(y_n)}{2 f'(y_n)} \quad (\text{B.31})$$

according to [104] equation (3.2). This is an estimation of the absolute error.

B.3.3 Applied to floating point numbers

Floating point numbers in IEEE 32 bit single precision format have a nearly constant relative error of $\Delta x/x = 2^{-24}$. As seen earlier in the Taylor series expansion equation (eqn. B.31), the error in every iteration step is absolute and in general dependent of y . If the error is expressed as a relative error ε_r the following holds

$$\varepsilon_{r_{n+1}} \equiv \frac{\varepsilon_{n+1}}{y} \quad (\text{B.32})$$

and so

$$\varepsilon_{r_{n+1}} = -\left(\frac{\varepsilon_n}{y}\right)^2 y \frac{f''}{2f'} \quad (\text{B.33})$$

for the function $f(y) = y^{-2}$ the term $y f''/2f'$ is constant (equal to $-3/2$) so the relative error ε_{r_n} is independent of y .

$$\varepsilon_{r_{n+1}} = \frac{3}{2} (\varepsilon_{r_n})^2 \quad (\text{B.34})$$

The conclusion of this is that the function $1/\sqrt{x}$ can be calculated with a specified accuracy.

B.3.4 Specification of the lookup table

To calculate the function $1/\sqrt{x}$ using the previously mentioned iteration scheme, it is clear that the first estimation of the solution must be accurate enough to get precise results. The requirements for the calculation are

- Maximum possible accuracy with the used IEEE format
- Use only one iteration step for maximum speed

The first requirement states that the result of $1/\sqrt{x}$ may have a relative error ε_r equal to the ε_r of a IEEE 32 bit single precision floating point number. From this the $1/\sqrt{x}$ of the initial approximation can be derived, rewriting the definition of the relative error for succeeding steps, equation (eqn. B.34)

$$\frac{\varepsilon_n}{y} = \sqrt{\varepsilon_{r_{n+1}} \frac{2f'}{yf''}} \quad (\text{B.35})$$

So for the lookup table the needed accuracy is

$$\frac{\Delta Y}{Y} = \sqrt{\frac{2}{3} 2^{-24}} \quad (\text{B.36})$$

which defines the width of the table that must be ≥ 13 bit.

At this point the relative error ε_{r_n} of the lookup table is known. From this the maximum relative error in the argument can be calculated as follows. The absolute error Δx is defined as

$$\Delta x \equiv \frac{\Delta Y}{Y'} \quad (\text{B.37})$$

and thus

$$\frac{\Delta x}{Y} = \frac{\Delta Y}{Y} (Y')^{-1} \quad (\text{B.38})$$

and thus

$$\Delta x = \text{constant} \frac{Y}{Y'} \quad (\text{B.39})$$

for the $1/\sqrt{x}$ function $Y/Y' \sim x$ holds, so $\Delta x/x = \text{constant}$. This is a property of the used floating point representation as earlier mentioned. The needed accuracy of the argument of the lookup table follows from

$$\frac{\Delta x}{x} = -2 \frac{\Delta Y}{Y} \quad (\text{B.40})$$

so, using the floating point accuracy, equation (eqn. B.36)

$$\frac{\Delta x}{x} = -2 \sqrt{\frac{2}{3} 2^{-24}} \quad (\text{B.41})$$

This defines the length of the lookup table which should be ≥ 12 bit.

B.3.5 Separate exponent and fraction computation

The used IEEE 32 bit single precision floating point format specifies that a number is represented by a exponent and a fraction. The previous section specifies for every possible floating point number the lookup table length and width. Only the size of the fraction of a floating point number defines the accuracy. The conclusion from this can be that the size of the lookup table is length of lookup table, earlier specified, times the size of the exponent ($2^{12}2^8, 1Mb$). The $1/\sqrt{x}$ function has the property that the exponent is independent of the fraction. This becomes clear if the floating point representation is used. Define

$$x \equiv (-1)^S (2^{E-127})(1.F) \quad (\text{B.42})$$

see Fig. B.1 where $0 \leq S \leq 1$, $0 \leq E \leq 255$, $1 \leq 1.F < 2$ and S, E, F integer (normalization conditions). The sign bit (S) can be omitted because $1/\sqrt{x}$ is only defined for $x > 0$. The $1/\sqrt{x}$ function applied to x results in

$$y(x) = \frac{1}{\sqrt{x}} \quad (\text{B.43})$$

or

$$y(x) = \frac{1}{\sqrt{(2^{E-127})(1.F)}} \quad (\text{B.44})$$

this can be rewritten as

$$y(x) = (2^{E-127})^{-1/2} (1.F)^{-1/2} \quad (\text{B.45})$$

Define

$$(2^{E'-127}) \equiv (2^{E-127})^{-1/2} \quad (\text{B.46})$$

$$1.F' \equiv (1.F)^{-1/2} \quad (\text{B.47})$$

then $\frac{1}{\sqrt{2}} < 1.F' \leq 1$ holds, so the condition $1 \leq 1.F' < 2$ which is essential for normalized real representation is not valid anymore. By introducing an extra term this can be corrected. Rewrite the $1/\sqrt{x}$ function applied to floating point numbers, equation (eqn. B.45) as

$$y(x) = (2^{\frac{127-E}{2}-1})(2(1.F)^{-1/2}) \quad (\text{B.48})$$

and

$$(2^{E'-127}) \equiv (2^{\frac{127-E}{2}-1}) \quad (\text{B.49})$$

$$1.F' \equiv 2(1.F)^{-1/2} \quad (\text{B.50})$$

then $\sqrt{2} < 1.F' \leq 2$ holds. This is not the exact valid range as defined for normalized floating point numbers in equation (eqn. B.42). The value 2 causes the problem. By mapping this value on the nearest representation < 2 this can be solved. The small error that is introduced by this approximation is within the allowable range.

The integer representation of the exponent is the next problem. Calculating $(2^{\frac{127-E}{2}-1})$ introduces a fractional result if $(127 - E) = \text{odd}$. This is again easily accounted for by splitting up the calculation into an odd and an even part. For $(127 - E) = \text{even}$ E' in equation (eqn. B.49) can be exactly calculated in integer arithmetic as a function of E .

$$E' = \frac{127 - E}{2} + 126 \quad (\text{B.51})$$

For $(127 - E) = \text{odd}$ equation (eqn. B.45) can be rewritten as

$$y(x) = (2^{\frac{127-E-1}{2}})(\frac{1.F}{2})^{-1/2} \quad (\text{B.52})$$

thus

$$E' = \frac{126 - E}{2} + 127 \quad (\text{B.53})$$

which also can be calculated exactly in integer arithmetic. Note that the fraction is automatically corrected for its range earlier mentioned, so the exponent does not need an extra correction.

The conclusions from this are:

- The fraction and exponent lookup table are independent. The fraction lookup table exists of two tables (odd and even exponent) so the odd/even information of the exponent (lsb bit) has to be used to select the right table.
- The exponent table is an 256 x 8 bit table, initialized for *odd* and *even*.

B.3.6 Implementation

The lookup tables can be generated by a small C program, which uses floating point numbers and operations with IEEE 32 bit single precision format. Note that because of the *oddeven* information that is needed, the fraction table is twice the size earlier specified (13 bit i.s.o. 12 bit).

The function according to equation (eqn. B.29) has to be implemented. Applied to the $1/\sqrt{x}$ function, equation (eqn. B.28) leads to

$$f = a - \frac{1}{y^2} \quad (\text{B.54})$$

and so

$$f' = \frac{2}{y^3} \quad (\text{B.55})$$

so

$$y_{n+1} = y_n - \frac{a - \frac{1}{y_n^2}}{\frac{2}{y_n^3}} \quad (\text{B.56})$$

or

$$y_{n+1} = \frac{y_n}{2}(3 - ay_n^2) \quad (\text{B.57})$$

Where y_0 can be found in the lookup tables, and y_1 gives the result to the maximum accuracy. It is clear that only one iteration extra (in double precision) is needed for a double precision result.

B.4 Modifying GROMACS

The following files have to be edited in case you want to add a bonded potential of any type.

1. include/bondf.h

2. `include/types/idef.h`
3. `include/types/nrn.h`
4. `include/types/enums.h`
5. `include/grompp.h`
6. `src/kernel/topdirs.c`
7. `src/gmxlib/tpxio.c`
8. `src/gmxlib/bondfree.c`
9. `src/gmxlib/ifunc.c`
10. `src/gmxlib/nrn.c`
11. `src/kernel/convparm.c`
12. `src/kernel/topdirs.c`
13. `src/kernel/topio.c`

Appendix C

Averages and fluctuations

C.1 Formulae for averaging

Note: this section was taken from ref [105].

When analyzing a MD trajectory averages $\langle x \rangle$ and fluctuations

$$\langle (\Delta x)^2 \rangle^{\frac{1}{2}} = \langle [x - \langle x \rangle]^2 \rangle^{\frac{1}{2}} \quad (\text{C.1})$$

of a quantity x are to be computed. The variance σ_x of a series of N_x values, $\{x_i\}$, can be computed from

$$\sigma_x = \sum_{i=1}^{N_x} x_i^2 - \frac{1}{N_x} \left(\sum_{i=1}^{N_x} x_i \right)^2 \quad (\text{C.2})$$

Unfortunately this formula is numerically not very accurate, especially when $\sigma_x^{\frac{1}{2}}$ is small compared to the values of x_i . The following (equivalent) expression is numerically more accurate

$$\sigma_x = \sum_{i=1}^{N_x} [x_i - \langle x \rangle]^2 \quad (\text{C.3})$$

with

$$\langle x \rangle = \frac{1}{N_x} \sum_{i=1}^{N_x} x_i \quad (\text{C.4})$$

Using eqns. C.2 and C.4 one has to go through the series of x_i values twice, once to determine $\langle x \rangle$ and again to compute σ_x , whereas eqn. C.1 requires only one sequential scan of the series $\{x_i\}$. However, one may cast eqn. C.2 in another form, containing partial sums, which allows for a sequential update algorithm. Define the partial sum

$$X_{n,m} = \sum_{i=n}^m x_i \quad (\text{C.5})$$

and the partial variance

$$\sigma_{n,m} = \sum_{i=n}^m \left[x_i - \frac{X_{n,m}}{m-n+1} \right]^2 \quad (\text{C.6})$$

It can be shown that

$$X_{n,m+k} = X_{n,m} + X_{m+1,m+k} \quad (\text{C.7})$$

and

$$\sigma_{n,m+k} = \sigma_{n,m} + \sigma_{m+1,m+k} + \frac{\left[\frac{X_{n,m}}{m-n+1} - \frac{X_{n,m+k}}{m+k-n+1} \right]^2}{k} * \quad (\text{C.8})$$

For $n = 1$ one finds

$$\sigma_{1,m+k} = \sigma_{1,m} + \sigma_{m+1,m+k} + \left[\frac{X_{1,m}}{m} - \frac{X_{1,m+k}}{m+k} \right]^2 \frac{m(m+k)}{k} \quad (\text{C.9})$$

and for $n = 1$ and $k = 1$ (eqn. C.8) becomes

$$\sigma_{1,m+1} = \sigma_{1,m} + \left[\frac{X_{1,m}}{m} - \frac{X_{1,m+1}}{m+1} \right]^2 m(m+1) \quad (\text{C.10})$$

$$= \sigma_{1,m} + \frac{[X_{1,m} - mx_{m+1}]^2}{m(m+1)} \quad (\text{C.11})$$

where we have used the relation

$$X_{1,m+1} = X_{1,m} + x_{m+1} \quad (\text{C.12})$$

Using formulae (eqn. C.11) and (eqn. C.12) the average

$$\langle x \rangle = \frac{X_{1,N_x}}{N_x} \quad (\text{C.13})$$

and the fluctuation

$$\langle (\Delta x)^2 \rangle^{\frac{1}{2}} = \left[\frac{\sigma_{1,N_x}}{N_x} \right]^{\frac{1}{2}} \quad (\text{C.14})$$

can be obtained by one sweep through the data.

C.2 Implementation

In GROMACS the instantaneous energies $E(m)$ are stored in the energy file, along with the values of $\sigma_{1,m}$ and $X_{1,m}$. Although the steps are counted from 0, for the energy and fluctuations steps are counted from 1. This means that the equations presented here are the ones that are implemented. We give somewhat lengthy derivations in this section to simplify checking of code and equations later on.

C.2.1 Part of a Simulation

It is not uncommon to perform a simulation where the first part, *e.g.* 100 ps, is taken as equilibration. However, the averages and fluctuations as printed in the log file are computed over the whole simulation. The equilibration time, which is now part of the simulation, may in such a case invalidate the averages and fluctuations, because these numbers are now dominated by the initial drift towards equilibrium.

Using eqns. C.7 and C.8 the average and standard deviation over part of the trajectory can be computed as:

$$X_{m+1,m+k} = X_{1,m+k} - X_{1,m} \quad (\text{C.15})$$

$$\sigma_{m+1,m+k} = \sigma_{1,m+k} - \sigma_{1,m} - \left[\frac{X_{1,m}}{m} - \frac{X_{1,m+k}}{m+k} \right]^2 \frac{m(m+k)}{k} \quad (\text{C.16})$$

or, more generally (with $p \geq 1$ and $q \geq p$):

$$X_{p,q} = X_{1,q} - X_{1,p-1} \quad (\text{C.17})$$

$$\sigma_{p,q} = \sigma_{1,q} - \sigma_{1,p-1} - \left[\frac{X_{1,p-1}}{p-1} - \frac{X_{1,q}}{q} \right]^2 \frac{(p-1)q}{q-p+1} \quad (\text{C.18})$$

Note that implementation of this is not entirely trivial, since energies are not stored every time step of the simulation. We therefore have to construct $X_{1,p-1}$ and $\sigma_{1,p-1}$ from the information at time p using eqns. C.11 and C.12:

$$X_{1,p-1} = X_{1,p} - x_p \quad (\text{C.19})$$

$$\sigma_{1,p-1} = \sigma_{1,p} - \frac{[X_{1,p-1} - (p-1)x_p]^2}{(p-1)p} \quad (\text{C.20})$$

C.2.2 Combining two simulations

Another frequently occurring problem is, that the fluctuations of two simulations must be combined. Consider the following example: we have two simulations (A) of n and (B) of m steps, in which the second simulation is a continuation of the first. However, the second simulation starts numbering from 1 instead of from $n+1$. For the partial sum this is no problem, we have to add $X_{1,n}^A$ from run A:

$$X_{1,n+m}^{AB} = X_{1,n}^A + X_{1,m}^B \quad (\text{C.21})$$

When we want to compute the partial variance from the two components we have to make a correction $\Delta\sigma$:

$$\sigma_{1,n+m}^{AB} = \sigma_{1,n}^A + \sigma_{1,m}^B + \Delta\sigma \quad (\text{C.22})$$

if we define x_i^{AB} as the combined and renumbered set of data points we can write:

$$\sigma_{1,n+m}^{AB} = \sum_{i=1}^{n+m} \left[x_i^{AB} - \frac{X_{1,n+m}^{AB}}{n+m} \right]^2 \quad (\text{C.23})$$

and thus

$$\sum_{i=1}^{n+m} \left[x_i^{AB} - \frac{X_{1,n+m}^{AB}}{n+m} \right]^2 = \sum_{i=1}^n \left[x_i^A - \frac{X_{1,n}^A}{n} \right]^2 + \sum_{i=1}^m \left[x_i^B - \frac{X_{1,m}^B}{m} \right]^2 + \Delta\sigma \quad (\text{C.24})$$

or

$$\begin{aligned} & \sum_{i=1}^{n+m} \left[(x_i^{AB})^2 - 2x_i^{AB} \frac{X_{1,n+m}^{AB}}{n+m} + \left(\frac{X_{1,n+m}^{AB}}{n+m} \right)^2 \right] - \\ & \sum_{i=1}^n \left[(x_i^A)^2 - 2x_i^A \frac{X_{1,n}^A}{n} + \left(\frac{X_{1,n}^A}{n} \right)^2 \right] - \\ & \sum_{i=1}^m \left[(x_i^B)^2 - 2x_i^B \frac{X_{1,m}^B}{m} + \left(\frac{X_{1,m}^B}{m} \right)^2 \right] = \Delta\sigma \end{aligned} \quad (\text{C.25})$$

all the x_i^2 terms drop out, and the terms independent of the summation counter i can be simplified:

$$\begin{aligned} & \frac{(X_{1,n+m}^{AB})^2}{n+m} - \frac{(X_{1,n}^A)^2}{n} - \frac{(X_{1,m}^B)^2}{m} - \\ & 2 \frac{X_{1,n+m}^{AB}}{n+m} \sum_{i=1}^{n+m} x_i^{AB} + 2 \frac{X_{1,n}^A}{n} \sum_{i=1}^n x_i^A + 2 \frac{X_{1,m}^B}{m} \sum_{i=1}^m x_i^B = \Delta\sigma \end{aligned} \quad (\text{C.26})$$

we recognize the three partial sums on the second line and use eqn. C.21 to obtain:

$$\Delta\sigma = \frac{(mX_{1,n}^A - nX_{1,m}^B)^2}{nm(n+m)} \quad (\text{C.27})$$

if we check this by inserting $m = 1$ we get back eqn. C.11

C.2.3 Summing energy terms

The g_energy program can also sum energy terms into one, *e.g.* potential + kinetic = total. For the partial averages this is again easy if we have S energy components s :

$$X_{m,n}^S = \sum_{i=m}^n \sum_{s=1}^S x_i^s = \sum_{s=1}^S \sum_{i=m}^n x_i^s = \sum_{s=1}^S X_{m,n}^s \quad (\text{C.28})$$

For the fluctuations it is less trivial again, considering for example that the fluctuation in potential and kinetic energy should cancel. Nevertheless we can try the same approach as before by writing:

$$\sigma_{m,n}^S = \sum_{s=1}^S \sigma_{m,n}^s + \Delta\sigma \quad (\text{C.29})$$

if we fill in eqn. C.6:

$$\sum_{i=m}^n \left[\left(\sum_{s=1}^S x_i^s \right) - \frac{X_{m,n}^S}{m-n+1} \right]^2 = \sum_{s=1}^S \sum_{i=m}^n \left[(x_i^s) - \frac{X_{m,n}^s}{m-n+1} \right]^2 + \Delta\sigma \quad (\text{C.30})$$

which we can expand to:

$$\begin{aligned} & \sum_{i=m}^n \left[\sum_{s=1}^S (x_i^s)^2 + \left(\frac{X_{m,n}^S}{m-n+1} \right)^2 - 2 \left(\frac{X_{m,n}^S}{m-n+1} \sum_{s=1}^S x_i^s + \sum_{s=1}^S \sum_{s'=s+1}^S x_i^s x_i^{s'} \right) \right] \\ - & \sum_{s=1}^S \sum_{i=m}^n \left[(x_i^s)^2 - 2 \frac{X_{m,n}^s}{m-n+1} x_i^s + \left(\frac{X_{m,n}^s}{m-n+1} \right)^2 \right] = \Delta\sigma \end{aligned} \quad (\text{C.31})$$

the terms with $(x_i^s)^2$ cancel, so that we can simplify to:

$$\begin{aligned} & \frac{\left(X_{m,n}^S \right)^2}{m-n+1} - 2 \frac{X_{m,n}^S}{m-n+1} \sum_{i=m}^n \sum_{s=1}^S x_i^s - 2 \sum_{i=m}^n \sum_{s=1}^S \sum_{s'=s+1}^S x_i^s x_i^{s'} - \\ & \sum_{s=1}^S \sum_{i=m}^n \left[-2 \frac{X_{m,n}^s}{m-n+1} x_i^s + \left(\frac{X_{m,n}^s}{m-n+1} \right)^2 \right] = \Delta\sigma \end{aligned} \quad (\text{C.32})$$

or

$$- \frac{\left(X_{m,n}^S \right)^2}{m-n+1} - 2 \sum_{i=m}^n \sum_{s=1}^S \sum_{s'=s+1}^S x_i^s x_i^{s'} + \sum_{s=1}^S \frac{\left(X_{m,n}^s \right)^2}{m-n+1} = \Delta\sigma \quad (\text{C.33})$$

If we now expand the first term using eqn. C.28 we obtain:

$$- \frac{\left(\sum_{s=1}^S X_{m,n}^s \right)^2}{m-n+1} - 2 \sum_{i=m}^n \sum_{s=1}^S \sum_{s'=s+1}^S x_i^s x_i^{s'} + \sum_{s=1}^S \frac{\left(X_{m,n}^s \right)^2}{m-n+1} = \Delta\sigma \quad (\text{C.34})$$

which we can reformulate to:

$$- 2 \left[\sum_{s=1}^S \sum_{s'=s+1}^S X_{m,n}^s X_{m,n}^{s'} + \sum_{i=m}^n \sum_{s=1}^S \sum_{s'=s+1}^S x_i^s x_i^{s'} \right] = \Delta\sigma \quad (\text{C.35})$$

or

$$- 2 \left[\sum_{s=1}^S X_{m,n}^s \sum_{s'=s+1}^S X_{m,n}^{s'} + \sum_{s=1}^S \sum_{i=m}^n x_i^s \sum_{s'=s+1}^S x_i^{s'} \right] = \Delta\sigma \quad (\text{C.36})$$

which gives

$$- 2 \sum_{s=1}^S \left[X_{m,n}^s \sum_{s'=s+1}^S \sum_{i=m}^n x_i^{s'} + \sum_{i=m}^n x_i^s \sum_{s'=s+1}^S x_i^{s'} \right] = \Delta\sigma \quad (\text{C.37})$$

Since we need all data points i to evaluate this, in general this is not possible. We can then make an estimate of $\sigma_{m,n}^S$ using only the data points that are available using the left hand side of eqn. C.30. While the average can be computed using all time steps in the simulation, the accuracy of the fluctuations is thus limited by the frequency with which energies are saved. Since this can be easily done with a program such as xmgr this is not built-in in GROMACS.

Appendix D

Manual Pages

D.1 options

All GROMACS programs have 6 standard options, of which some are hidden by default:

Other options

<code>-h</code>	<code>bool</code>	<code>no</code>	Print help info and quit
<code>-X</code>	<code>bool</code>	<code>no</code>	Use dialog box GUI to edit command line options
<code>-nice</code>	<code>int</code>	<code>0</code>	Set the nicelevel

- If the configuration script found Motif or Lesstif on your system, you can use the graphical interface (if not, you will get an error):
`-X bool no` Use dialog box GUI to edit command line options
- When compiled on an SGI-IRIX system, all GROMACS programs have an additional option:
`-npri int 0` Set non blocking priority (try 128)
- Optional files are not used unless the option is set, in contrast to non optional files, where the default file name is used when the option is not set.
- All GROMACS programs will accept file options without a file extension or filename being specified. In such cases the default filenames will be used. With multiple input file types, such as generic structure format, the directory will be searched for files of each type with the supplied or default name. When no such file is found, or with output files the first file type will be used.
- All GROMACS programs with the exception of `mdrun`, `nmr` and `eneconv` check if the command line options are valid. If this is not the case, the program will be halted.
- Enumerated options (enum) should be used with one of the arguments listed in the option description, the argument may be abbreviated. The first match to the shortest argument in the list will be selected.
- Vector options can be used with 1 or 3 parameters. When only one parameter is supplied the two others are also set to this value.
- For many GROMACS programs, the time options can be supplied in different time units, depending on the setting of the `-tu` option.
- All GROMACS programs can read compressed or g-zipped files. There might be a problem with reading compressed `.xtc`, `.trr` and `.trj` files, but these will not compress very well anyway.

- Most GROMACS programs can process a trajectory with less atoms than the run input or structure file, but only if the trajectory consists of the first *n* atoms of the run input or structure file.
- Many GROMACS programs will accept the `-tu` option to set the time units to use in output files (e.g. for `xmgr` graphs or `xpm` matrices) and in all time options.

D.2 anadock

`anadock` analyses the results of an Autodock run and clusters the structures together, based on distance or RMSD. The docked energy and free energy estimates are analysed, and for each cluster the energy statistics are printed.

An alternative approach to this is to cluster the structures first (using `g_cluster` and then sort the clusters on either lowest energy or average energy.

Files

<code>-f</code>	<code>eiwit.pdb</code>	Input	Protein data bank file
<code>-ox</code>	<code>cluster.pdb</code>	Output	Protein data bank file
<code>-od</code>	<code>edocked.xvg</code>	Output	<code>xvgr/xmgr</code> file
<code>-of</code>	<code>efree.xvg</code>	Output	<code>xvgr/xmgr</code> file
<code>-g</code>	<code>anadock.log</code>	Output	Log file

Other options

<code>-h</code>	bool	no	Print help info and quit
<code>-nice</code>	int	0	Set the nicelevel
<code>-xvgr</code>	bool	yes	Add specific codes (legends etc.) in the output xvg files for the <code>xmgrace</code> program
<code>-free</code>	bool	no	Use Free energy estimate from autodock for sorting the classes
<code>-rms</code>	bool	yes	Cluster on RMS or distance
<code>-cutoff</code>	real	0.2	Maximum RMSD/distance for belonging to the same cluster

D.3 do_dssp

`do_dssp` reads a trajectory file and computes the secondary structure for each time frame calling the `dssp` program. If you do not have the `dssp` program, get it. `do_dssp` assumes that the `dssp` executable is `/usr/local/bin/dssp`. If this is not the case, then you should set an environment variable **DSSP** pointing to the `dssp` executable, e.g.:

```
setenv DSSP /opt/dssp/bin/dssp
```

The structure assignment for each residue and time is written to an `.xpm` matrix file. This file can be visualized with for instance `xv` and can be converted to postscript with `xpm2ps`. The number of residues with each secondary structure type and the total secondary structure (`-sss`) count as a function of time are also written to file (`-sc`).

Solvent accessible surface (SAS) per residue can be calculated, both in absolute values (\AA^2) and in fractions of the maximal accessible surface of a residue. The maximal accessible surface is defined as the accessible surface of a residue in a chain of glycines. **Note** that the program `g_sas` can also compute SAS and that is more efficient.

Finally, this program can dump the secondary structure in a special file `ssdump.dat` for usage in the program `g_chi`. Together these two programs can be used to analyze dihedral properties as a function of secondary structure type.

Files

-f	traj.xtc	Input	Generic trajectory: xtc trr trj gro g96 pdb
-s	topol.tpr	Input	Structure+mass(db): tpr tpb tpa gro g96 pdb xml
-n	index.ndx	Input, Opt.	Index file
-ssdump	ssdump.dat	Output, Opt.	Generic data file
-map	ss.map	Input, Lib.	File that maps matrix data to colors
-o	ss.xpm	Output	X PixMap compatible matrix file
-sc	scount.xvg	Output	xvgr/xmgr file
-a	area.xpm	Output, Opt.	X PixMap compatible matrix file
-ta	totarea.xvg	Output, Opt.	xvgr/xmgr file
-aa	averarea.xvg	Output, Opt.	xvgr/xmgr file

Other options

-h	bool	no	Print help info and quit
-nice	int	19	Set the nicelevel
-b	time	0	First frame (ps) to read from trajectory
-e	time	0	Last frame (ps) to read from trajectory
-dt	time	0	Only use frame when $t \text{ MOD } dt = \text{first time (ps)}$
-tu	enum	ps	Time unit: ps, fs, ns, us, ms or s
-w	bool	no	View output xvg, xpm, eps and pdb files
-xvgr	bool	yes	Add specific codes (legends etc.) in the output xvg files for the xmgrace program
-sss	string	HEBT	Secondary structures for structure count

- The program is very slow

D.4 editconf

editconf converts generic structure format to .gro, .g96 or .pdb.

The box can be modified with options -box, -d and -angles. Both -box and -d will center the system in the box.

Option -bt determines the box type: triclinic is a triclinic box, cubic is a rectangular box with all sides equal dodecahedron represents a rhombic dodecahedron and octahedron is a truncated octahedron. The last two are special cases of a triclinic box. The length of the three box vectors of the truncated octahedron is the shortest distance between two opposite hexagons. The volume of a dodecahedron is 0.71 and that of a truncated octahedron is 0.77 of that of a cubic box with the same periodic image distance.

Option -box requires only one value for a cubic box, dodecahedron and a truncated octahedron.

With -d and a triclinic box the size of the system in the x, y and z directions is used. With -d and cubic, dodecahedron or octahedron boxes, the dimensions are set to the diameter of the system (largest distance between atoms) plus twice the specified distance.

Option -angles is only meaningful with option -box and a triclinic box and can not be used with option -d.

When -n or -ndef is set, a group can be selected for calculating the size and the geometric center, otherwise the whole system is used.

-rotate rotates the coordinates and velocities.

-princ aligns the principal axes of the system along the coordinate axes, this may allow you to decrease the box volume, but beware that molecules can rotate significantly in a nanosecond.

Scaling is applied before any of the other operations are performed. Boxes can be scaled to give a certain density (option `-density`). A special feature of the scaling option, when the factor -1 is given in one dimension, one obtains a mirror image, mirrored in one of the plains, when one uses -1 in three dimensions a point-mirror image is obtained.

Groups are selected after all operations have been applied.

Periodicity can be removed in a crude manner. It is important that the box sizes at the bottom of your input file are correct when the periodicity is to be removed.

When writing `.pdb` files, B-factors can be added with the `-bf` option. B-factors are read from a file with with following format: first line states number of entries in the file, next lines state an index followed by a B-factor. The B-factors will be attached per residue unless an index is larger than the number of residues or unless the `-atom` option is set. Obviously, any type of numeric data can be added instead of B-factors. `-legend` will produce a row of CA atoms with B-factors ranging from the minimum to the maximum value found, effectively making a legend for viewing.

With the option `-mead` a special `pdb` (`pqr`) file for the MEAD electrostatics program (Poisson-Boltzmann solver) can be made. A further prerequisite is that the input file is a run input file. The B-factor field is then filled with the Van der Waals radius of the atoms while the occupancy field will hold the charge.

The option `-grasp` is similar, but it puts the charges in the B-factor and the radius in the occupancy.

Finally with option `-label` `editconf` can add a chain identifier to a `pdb` file, which can be useful for analysis with e.g. `rasmol`.

To convert a truncated octahedron file produced by a package which uses a cubic box with the corners cut off (such as `Gromos`) use:

```
editconf -f <in> -rotate 0 45 35.264 -bt o -box <veclen> -o <out>
```

where `veclen` is the size of the cubic box times $\sqrt{3}/2$.

Files

<code>-f</code>	<code>conf.gro</code>	Input	Generic structure: <code>gro g96 pdb tpr tpb tpa xml</code>
<code>-n</code>	<code>index.ndx</code>	Input, Opt.	Index file
<code>-o</code>	<code>out.gro</code>	Output, Opt.	Generic structure: <code>gro g96 pdb xml</code>
<code>-mead</code>	<code>mead.pqr</code>	Output, Opt.	Coordinate file for MEAD
<code>-bf</code>	<code>bfact.dat</code>	Input, Opt.	Generic data file

Other options

<code>-h</code>	<code>bool</code>	<code>no</code>	Print help info and quit
<code>-nice</code>	<code>int</code>	<code>0</code>	Set the nicelevel
<code>-w</code>	<code>bool</code>	<code>no</code>	View output <code>xvg</code> , <code>xpm</code> , <code>eps</code> and <code>pdb</code> files
<code>-ndef</code>	<code>bool</code>	<code>no</code>	Choose output from default index groups
<code>-bt</code>	<code>enum</code>		
	<code>triclinic</code>		Box type for <code>-box</code> and <code>-d</code> : <code>triclinic</code> , <code>cubic</code> , <code>dodecahedron</code> or <code>octahedron</code>
<code>-box</code>	<code>vector</code>	<code>0 0 0</code>	Box vector lengths (a,b,c)
<code>-angles</code>	<code>vector</code>	<code>90 90</code>	Angles between the box vectors (bc,ac,ab)
<code>-d</code>	<code>real</code>	<code>0</code>	Distance between the solute and the box
<code>-c</code>	<code>bool</code>	<code>no</code>	Center molecule in box (implied by <code>-box</code> and <code>-d</code>)
<code>-center</code>	<code>vector</code>	<code>0 0 0</code>	Coordinates of geometrical center
<code>-translate</code>	<code>vector</code>	<code>0 0 0</code>	Translation
<code>-rotate</code>	<code>vector</code>	<code>0 0 0</code>	Rotation around the X, Y and Z axes in degrees
<code>-princ</code>	<code>bool</code>	<code>no</code>	Orient molecule(s) along their principal axes
<code>-scale</code>	<code>vector</code>	<code>1 1 1</code>	Scaling factor
<code>-density</code>	<code>real</code>	<code>1000</code>	Density (g/l) of the output box achieved by scaling
<code>-vol</code>	<code>bool</code>	<code>yes</code>	Compute and print volume of the box
<code>-pbc</code>	<code>bool</code>	<code>no</code>	Remove the periodicity (make molecule whole again)

<code>-grasp</code>	bool	no	Store the charge of the atom in the B-factor field and the radius of the atom in the occupancy field
<code>-rvdw</code>	real	0.12	Default Van der Waals radius (in nm) if one can not be found in the database or if no parameters are present in the topology file
<code>-sig56</code>	real	0	Use $r_{min}/2$ (minimum in the Van der Waals potential) rather than $\sigma/2$
<code>-vdwread</code>	bool	no	Read the Van der Waals radii from the file <code>vdwradii.dat</code> rather than computing the radii based on the force field
<code>-atom</code>	bool	no	Force B-factor attachment per atom
<code>-legend</code>	bool	no	Make B-factor legend
<code>-label</code>	string	A	Add chain label for all residues

- For complex molecules, the periodicity removal routine may break down, in that case you can use `trjconv`

D.5 *eneconv*

With *multiple files* specified for the `-f` option:

Concatenates several energy files in sorted order. In case of double time frames the one in the later file is used. By specifying `-settime` you will be asked for the start time of each file. The input files are taken from the command line, such that the command `eneconv -o fixed.edr *.edr` should do the trick.

With *one file* specified for `-f`:

Reads one energy file and writes another, applying the `-dt`, `-offset`, `-t0` and `-settime` options and converting to a different format if necessary (indicated by file extensions).

`-settime` is applied first, then `-dt/-offset` followed by `-b` and `-e` to select which frames to write.

Files

<code>-f</code>	<code>ener.edr</code>	Input, Mult.	Generic energy: <code>edr ene</code>
<code>-o</code>	<code>fixed.edr</code>	Output	Generic energy: <code>edr ene</code>

Other options

<code>-h</code>	bool	no	Print help info and quit
<code>-nice</code>	int	19	Set the nicelevel
<code>-b</code>	real	-1	First time to use
<code>-e</code>	real	-1	Last time to use
<code>-dt</code>	real	0	Only write out frame when $t \text{ MOD } dt = \text{offset}$
<code>-offset</code>	real	0	Time offset for <code>-dt</code> option
<code>-settime</code>	bool	no	Change starting time interactively
<code>-sort</code>	bool	yes	Sort energy files (not frames)
<code>-scalefac</code>	real	1	Multiply energy component by this factor
<code>-error</code>	bool	yes	Stop on errors in the file

- When combining trajectories the σ and E^2 (necessary for statistics) are not updated correctly. Only the actual energy is correct. One thus has to compute statistics in another way.

D.6 *ffscan*

The `ffscan` program performs a single point energy and force calculation in which the force field is modified. This way a range of parameters can be changed and tested for reproduction of e.g. quantum chemical or

experimental data. A grid scan over the parameters is done as specified using command line arguments. All parameters that reproduce the energy within a given absolute tolerance are printed to a log file.

Obviously polarizable models can be used, and shell optimisation is performed if necessary. Also, like in `mddrun` table functions can be used for user defined potential functions.

If the option `-ga` with appropriate file is passed, a genetic algorithm will be used rather than a grid scan.

Files

<code>-s</code>	<code>topol.tpr</code>	Input	Generic run input: tpr tpb tpa xml
<code>-g</code>	<code>md.log</code>	Output	Log file
<code>-table</code>	<code>table.xvg</code>	Input, Opt.	xvgr/xmgr file
<code>-parm</code>	<code>params.dat</code>	Input	Generic data file
<code>-ga</code>	<code>genalg.dat</code>	Input, Opt.	Generic data file
<code>-c</code>	<code>junk.gro</code>	Output	Coordinate file in Gromos-87 format
<code>-e</code>	<code>junk.edr</code>	Output	Generic energy: edr ene
<code>-o</code>	<code>junk.trr</code>	Output	Full precision trajectory: trr trj

Other options

<code>-h</code>	bool	no	Print help info and quit
<code>-nice</code>	int	19	Set the nicelevel
<code>-xvgr</code>	bool	yes	Add specific codes (legends etc.) in the output xvg files for the xmgrace program
<code>-tol</code>	real	0.1	Energy tolerance (kJ/mol) (zero means everything is printed)
<code>-fmax</code>	real	100	Force tolerance (zero means everything is printed)
<code>-comb</code>	bool	yes	Use combination rules
<code>-npow</code>	real	12	Power for LJ in case of table use
<code>-logeps</code>	bool	no	Use a logarithmic scale for epsilon
<code>-v</code>	bool	no	Be loud and noisy
<code>-epot</code>	real	0	Target energy (kJ/mol)
<code>-fepot</code>	real	1	Factor for scaling energy violations (0 turns energy contribution off)
<code>-pres</code>	real	1	Value for reference pressure
<code>-fpres</code>	real	0.1	Factor for scaling pressure violations (0 turns pressure contribution off)
<code>-fmsf</code>	real	0.1	Factor for scaling mean square force violations (0 turns MSF contribution off)
<code>-molsize</code>	int	1	Number of atoms per molecule
<code>-nmol</code>	int	1	Number of molecules (Epot is divided by this value!)

D.7 g_anaeig

`g_anaeig` analyzes eigenvectors. The eigenvectors can be of a covariance matrix (`g_covar`) or of a Normal Modes analysis (`g_nmeig`).

When a trajectory is projected on eigenvectors, all structures are fitted to the structure in the eigenvector file, if present, otherwise to the structure in the structure file. When no run input file is supplied, periodicity will not be taken into account. Most analyses are performed on eigenvectors `-first` to `-last`, but when `-first` is set to `-1` you will be prompted for a selection.

`-comp`: plot the vector components per atom of eigenvectors `-first` to `-last`.

`-rmsf`: plot the RMS fluctuation per atom of eigenvectors `-first` to `-last` (requires `-eig`).

`-proj`: calculate projections of a trajectory on eigenvectors `-first` to `-last`. The projections of a trajectory on the eigenvectors of its covariance matrix are called principal components (pc's). It is often useful to check the cosine content the pc's, since the pc's of random diffusion are cosines with the number

of periods equal to half the pc index. The cosine content of the pc's can be calculated with the program *g_analyze*.

`-2d`: calculate a 2d projection of a trajectory on eigenvectors `-first` and `-last`.

`-3d`: calculate a 3d projection of a trajectory on the first three selected eigenvectors.

`-filt`: filter the trajectory to show only the motion along eigenvectors `-first` to `-last`.

`-extr`: calculate the two extreme projections along a trajectory on the average structure and interpolate `-nframes` frames between them, or set your own extremes with `-max`. The eigenvector `-first` will be written unless `-first` and `-last` have been set explicitly, in which case all eigenvectors will be written to separate files. Chain identifiers will be added when writing a `.pdb` file with two or three structures (you can use `rasmol -nmrpdb` to view such a `pdb` file).

Overlap calculations between covariance analysis:

NOTE: the analysis should use the same fitting structure

`-over`: calculate the subspace overlap of the eigenvectors in file `-v2` with eigenvectors `-first` to `-last` in file `-v`.

`-inpr`: calculate a matrix of inner-products between eigenvectors in files `-v` and `-v2`. All eigenvectors of both files will be used unless `-first` and `-last` have been set explicitly.

When `-v`, `-eig`, `-v2` and `-eig2` are given, a single number for the overlap between the covariance matrices is generated. The formulas are:

$$\text{difference} = \sqrt{\text{tr}((\sqrt{M1} - \sqrt{M2})^2)}$$

$$\text{normalized overlap} = 1 - \text{difference} / \sqrt{\text{tr}(M1) + \text{tr}(M2)}$$

$$\text{shape overlap} = 1 - \sqrt{\text{tr}((\sqrt{M1/\text{tr}(M1)} - \sqrt{M2/\text{tr}(M2)})^2)}$$

where *M1* and *M2* are the two covariance matrices and *tr* is the trace of a matrix. The numbers are proportional to the overlap of the square root of the fluctuations. The normalized overlap is the most useful number, it is 1 for identical matrices and 0 when the sampled subspaces are orthogonal.

Files

<code>-v</code>	<code>eigenvec.trr</code>	Input	Full precision trajectory: <code>tr trj</code>
<code>-v2</code>	<code>eigenvec2.trr</code>	Input, Opt.	Full precision trajectory: <code>tr trj</code>
<code>-f</code>	<code>traj.xtc</code>	Input, Opt.	Generic trajectory: <code>xtc trr trj gro g96 pdb</code>
<code>-s</code>	<code>topol.tpr</code>	Input, Opt.	Structure+mass(db): <code>tpr tpb tpa gro g96 pdb xml</code>
<code>-n</code>	<code>index.ndx</code>	Input, Opt.	Index file
<code>-eig</code>	<code>eigenval.xvg</code>	Input, Opt.	<code>xvgr/xmgr</code> file
<code>-eig2</code>	<code>eigenval2.xvg</code>	Input, Opt.	<code>xvgr/xmgr</code> file
<code>-comp</code>	<code>eigcomp.xvg</code>	Output, Opt.	<code>xvgr/xmgr</code> file
<code>-rmsf</code>	<code>eigrmsf.xvg</code>	Output, Opt.	<code>xvgr/xmgr</code> file
<code>-proj</code>	<code>proj.xvg</code>	Output, Opt.	<code>xvgr/xmgr</code> file
<code>-2d</code>	<code>2dproj.xvg</code>	Output, Opt.	<code>xvgr/xmgr</code> file
<code>-3d</code>	<code>3dproj.pdb</code>	Output, Opt.	Generic structure: <code>gro g96 pdb xml</code>
<code>-filt</code>	<code>filtered.xtc</code>	Output, Opt.	Generic trajectory: <code>xtc trr trj gro g96 pdb</code>
<code>-extr</code>	<code>extreme.pdb</code>	Output, Opt.	Generic trajectory: <code>xtc trr trj gro g96 pdb</code>
<code>-over</code>	<code>overlap.xvg</code>	Output, Opt.	<code>xvgr/xmgr</code> file
<code>-inpr</code>	<code>inprod.xpm</code>	Output, Opt.	X PixMap compatible matrix file

Other options

<code>-h</code>	<code>bool</code>	<code>no</code>	Print help info and quit
<code>-nice</code>	<code>int</code>	<code>19</code>	Set the nicelevel
<code>-b</code>	<code>time</code>	<code>0</code>	First frame (ps) to read from trajectory
<code>-e</code>	<code>time</code>	<code>0</code>	Last frame (ps) to read from trajectory
<code>-dt</code>	<code>time</code>	<code>0</code>	Only use frame when <code>t MOD dt = first time (ps)</code>
<code>-tu</code>	<code>enum</code>	<code>ps</code>	Time unit: <code>ps, fs, ns, us, ms</code> or <code>s</code>

<code>-w</code>	bool	no	View output xvg, xpm, eps and pdb files
<code>-xvgr</code>	bool	yes	Add specific codes (legends etc.) in the output xvg files for the xmgrace program
<code>-first</code>	int	1	First eigenvector for analysis (-1 is select)
<code>-last</code>	int	8	Last eigenvector for analysis (-1 is till the last)
<code>-skip</code>	int	1	Only analyse every nr-th frame
<code>-max</code>	real	0	Maximum for projection of the eigenvector on the average structure, max=0 gives the extremes
<code>-nframes</code>	int	2	Number of frames for the extremes output
<code>-split</code>	bool	no	Split eigenvector projections where time is zero

D.8 g_analyze

`g_analyze` reads an ascii file and analyzes data sets. A line in the input file may start with a time (see option `-time`) and any number of y values may follow. Multiple sets can also be read when they are separated by `&` (option `-n`), in this case only one y value is read from each line. All lines starting with `#` and `@` are skipped. All analyses can also be done for the derivative of a set (option `-d`).

All options, except for `-av` and `-power` assume that the points are equidistant in time.

`g_analyze` always shows the average and standard deviation of each set. For each set it also shows the relative deviation of the third and fourth cumulant from those of a Gaussian distribution with the same standard deviation.

Option `-ac` produces the autocorrelation function(s).

Option `-cc` plots the resemblance of set *i* with a cosine of *i*/2 periods. The formula is:

$$2 \int_0^T y(t) \cos(i \pi t) dt / \int_0^T y(t) y(t) dt$$

This is useful for principal components obtained from covariance analysis, since the principal components of random diffusion are pure cosines.

Option `-msd` produces the mean square displacement(s).

Option `-dist` produces distribution plot(s).

Option `-av` produces the average over the sets. Error bars can be added with the option `-errbar`. The errorbars can represent the standard deviation, the error (assuming the points are independent) or the interval containing 90% of the points, by discarding 5% of the points at the top and the bottom.

Option `-ee` produces error estimates using block averaging. A set is divided in a number of blocks and averages are calculated for each block. The error for the total average is calculated from the variance between averages of the *m* blocks B_i as follows: $\text{error}^2 = \text{Sum} (B_i - \langle B \rangle)^2 / (m*(m-1))$. These errors are plotted as a function of the block size. Also an analytical block average curve is plotted, assuming that the autocorrelation is a sum of two exponentials. The analytical curve for the block average is:

$$f(t) = \text{sigma} \sqrt{2/T} (a (\tau_1 ((\exp(-t/\tau_1) - 1) \tau_1/t + 1)) + (1-a) (\tau_2 ((\exp(-t/\tau_2) - 1) \tau_2/t + 1)))),$$

where *T* is the total time. *a*, τ_1 and τ_2 are obtained by fitting $f^2(t)$ to error^2 . When the actual block average is very close to the analytical curve, the error is $\text{sigma} * \sqrt{2/T} (a \tau_1 + (1-a) \tau_2)$. The complete derivation is given in B. Hess, J. Chem. Phys. 116:209-217, 2002.

Option `-filter` prints the RMS high-frequency fluctuation of each set and over all sets with respect to a filtered average. The filter is proportional to $\cos(\pi t/\text{len})$ where *t* goes from $-\text{len}/2$ to $\text{len}/2$. *len* is supplied with the option `-filter`. This filter reduces oscillations with period $\text{len}/2$ and *len* by a factor of 0.79 and 0.33 respectively.

Option `-power` fits the data to $b \hat{t}^a$, which is accomplished by fitting to $a t + b$ on log-log scale. All points after the first zero or negative value are ignored.

Files

-f	graph.xvg	Input	xvgr/xmgr file
-ac	autocorr.xvg	Output, Opt.	xvgr/xmgr file
-msd	msd.xvg	Output, Opt.	xvgr/xmgr file
-cc	coscont.xvg	Output, Opt.	xvgr/xmgr file
-dist	distr.xvg	Output, Opt.	xvgr/xmgr file
-av	average.xvg	Output, Opt.	xvgr/xmgr file
-ee	errest.xvg	Output, Opt.	xvgr/xmgr file
-g	fitlog.log	Output, Opt.	Log file

Other options

-h	bool	no	Print help info and quit
-nice	int	0	Set the nicelevel
-w	bool	no	View output xvg, xpm, eps and pdb files
-xvgr	bool	yes	Add specific codes (legends etc.) in the output xvg files for the xmgrace program
-time	bool	yes	Expect a time in the input
-b	real	-1	First time to read from set
-e	real	-1	Last time to read from set
-n	int	1	Read # sets seperated by &
-d	bool	no	Use the derivative
-bw	real	0.1	Binwidth for the distribution
-errbar	enum	none	Error bars for -av: none, stddev, error or 90
-integrate	bool	no	Integrate data function(s) numerically using trapezium rule
-aver_start	real	0	Start averaging the integral from here
-xydy	bool	no	Interpret second data set as error in the y values for integrating
-filter	real	0	Print the high-frequency fluctuation after filtering with a cosine filter of length #
-power	bool	no	Fit data to: $b \hat{t}^a$
-subav	bool	yes	Subtract the average before autocorrelating
-oneacf	bool	no	Calculate one ACF over all sets
-acflen	int	-1	Length of the ACF, default is half the number of frames
-normalize	bool	yes	Normalize ACF
-P	enum	0	Order of Legendre polynomial for ACF (0 indicates none): 0, 1, 2 or 3
-fitfn	enum	none	Fit function: none, exp, aexp, exp_exp, vac, exp5, exp7 or exp9
-ncskip	int	0	Skip N points in the output file of correlation functions
-beginfit	real	0	Time where to begin the exponential fit of the correlation function
-endfit	real	-1	Time where to end the exponential fit of the correlation function, -1 is till the end

D.9 g_angle

g_angle computes the angle distribution for a number of angles or dihedrals. This way you can check whether your simulation is correct. With option -ov you can plot the average angle of a group of angles as a function of time. With the -all option the first graph is the average, the rest are the individual angles.

With the -of option *g_angle* also calculates the fraction of trans dihedrals (only for dihedrals) as function of time, but this is probably only fun for a selected few.

With option -oc a dihedral correlation function is calculated.

It should be noted that the indexfile should contain atom-triples for angles or atom-quadruplets for dihedrals. If this is not the case, the program will crash.

With option `-or` a trajectory file is dumped containing cos andsin of selected dihedral angles which subsequently can be used as input for a PCA analysis using `g_covar`.

Files

<code>-f</code>	<code>traj.xtc</code>	Input	Generic trajectory: xtc trr trj gro g96 pdb
<code>-s</code>	<code>topol.tpr</code>	Input	Generic run input: tpr tpb tpa xml
<code>-n</code>	<code>angle.ndx</code>	Input	Index file
<code>-od</code>	<code>angdist.xvg</code>	Output	xvgr/xmgr file
<code>-ov</code>	<code>angaver.xvg</code>	Output, Opt.	xvgr/xmgr file
<code>-of</code>	<code>dihfrac.xvg</code>	Output, Opt.	xvgr/xmgr file
<code>-ot</code>	<code>dihtrans.xvg</code>	Output, Opt.	xvgr/xmgr file
<code>-oh</code>	<code>trhisto.xvg</code>	Output, Opt.	xvgr/xmgr file
<code>-oc</code>	<code>dihcorr.xvg</code>	Output, Opt.	xvgr/xmgr file
<code>-or</code>	<code>traj.trr</code>	Output, Opt.	Trajectory in portable xdr format

Other options

<code>-h</code>	<code>bool</code>	<code>no</code>	Print help info and quit
<code>-nice</code>	<code>int</code>	<code>19</code>	Set the nicelevel
<code>-b</code>	<code>time</code>	<code>0</code>	First frame (ps) to read from trajectory
<code>-e</code>	<code>time</code>	<code>0</code>	Last frame (ps) to read from trajectory
<code>-dt</code>	<code>time</code>	<code>0</code>	Only use frame when $t \text{ MOD } dt = \text{first time (ps)}$
<code>-w</code>	<code>bool</code>	<code>no</code>	View output xvg, xpm, eps and pdb files
<code>-xvgr</code>	<code>bool</code>	<code>yes</code>	Add specific codes (legends etc.) in the output xvg files for the xmgrace program
<code>-type</code>	<code>enum</code>	<code>angle</code>	Type of angle to analyse: <code>angle</code> , <code>dihedral</code> , <code>improper</code> or <code>ryckaert-bellemans</code>
<code>-all</code>	<code>bool</code>	<code>no</code>	Plot all angles separately in the averages file, in the order of appearance in the index file.
<code>-binwidth</code>	<code>real</code>	<code>1</code>	binwidth (degrees) for calculating the distribution
<code>-periodic</code>	<code>bool</code>	<code>yes</code>	Print dihedral angles modulo 360 degrees
<code>-chandler</code>	<code>bool</code>	<code>no</code>	Use Chandler correlation function ($N[\text{trans}] = 1$, $N[\text{gauche}] = 0$) rather than cosine correlation function. Trans is defined as $\phi < -60$ or $\phi > 60$.
<code>-avercorr</code>	<code>bool</code>	<code>no</code>	Average the correlation functions for the individual angles/dihedrals
<code>-acflen</code>	<code>int</code>	<code>-1</code>	Length of the ACF, default is half the number of frames
<code>-normalize</code>	<code>bool</code>	<code>yes</code>	Normalize ACF
<code>-P</code>	<code>enum</code>	<code>0</code>	Order of Legendre polynomial for ACF (0 indicates none): 0, 1, 2 or 3
<code>-fitfn</code>	<code>enum</code>	<code>none</code>	Fit function: <code>none</code> , <code>exp</code> , <code>aexp</code> , <code>exp-exp</code> , <code>vac</code> , <code>exp5</code> , <code>exp7</code> or <code>exp9</code>
<code>-ncskip</code>	<code>int</code>	<code>0</code>	Skip N points in the output file of correlation functions
<code>-beginfit</code>	<code>real</code>	<code>0</code>	Time where to begin the exponential fit of the correlation function
<code>-endfit</code>	<code>real</code>	<code>-1</code>	Time where to end the exponential fit of the correlation function, -1 is till the end

- Counting transitions only works for dihedrals with multiplicity 3

D.10 g_bond

`g_bond` makes a distribution of bond lengths. If all is well a gaussian distribution should be made when using a harmonic potential. bonds are read from a single group in the index file in order `i1-j1 i2-j2` thru `in-jn`.

`-tol` gives the half-width of the distribution as a fraction of the bondlength (`-blen`). That means, for a bond of 0.2 a tol of 0.1 gives a distribution from 0.18 to 0.22.

Option `-d` plots all the distances as a function of time. This requires a structure file for the atom and residue names in the output. If however the option `-averdist` is given (as well or separately) the average bond length is plotted instead.

Files

<code>-f</code>	<code>traj.xtc</code>	Input	Generic trajectory: xtc trr trj gro g96 pdb
<code>-n</code>	<code>index.ndx</code>	Input	Index file
<code>-s</code>	<code>topol.tpr</code>	Input, Opt.	Structure+mass(db): tpr tpb tpa gro g96 pdb xml
<code>-o</code>	<code>bonds.xvg</code>	Output	xvgr/xmgr file
<code>-l</code>	<code>bonds.log</code>	Output, Opt.	Log file
<code>-d</code>	<code>distance.xvg</code>	Output, Opt.	xvgr/xmgr file

Other options

<code>-h</code>	bool	no	Print help info and quit
<code>-nice</code>	int	19	Set the nicelevel
<code>-b</code>	time	0	First frame (ps) to read from trajectory
<code>-e</code>	time	0	Last frame (ps) to read from trajectory
<code>-dt</code>	time	0	Only use frame when $t \text{ MOD } dt = \text{first time (ps)}$
<code>-w</code>	bool	no	View output xvg, xpm, eps and pdb files
<code>-xvgr</code>	bool	yes	Add specific codes (legends etc.) in the output xvg files for the xmgrace program
<code>-blen</code>	real	-1	Bond length. By default length of first bond
<code>-tol</code>	real	0.1	Half width of distribution as fraction of blen
<code>-aver</code>	bool	yes	Average bond length distributions
<code>-averdist</code>	bool	yes	Average distances (turns on -d)

- It should be possible to get bond information from the topology.

D.11 *g_bundle*

g_bundle analyzes bundles of axes. The axes can be for instance helix axes. The program reads two index groups and divides both of them in `-na` parts. The centers of mass of these parts define the tops and bottoms of the axes. Several quantities are written to file: the axis length, the distance and the z-shift of the axis mid-points with respect to the average center of all axes, the total tilt, the radial tilt and the lateral tilt with respect to the average axis.

With options `-ok`, `-okr` and `-okl` the total, radial and lateral kinks of the axes are plotted. An extra index group of kink atoms is required, which is also divided into `-na` parts. The kink angle is defined as the angle between the kink-top and the bottom-kink vectors.

With option `-oa` the top, mid (or kink when `-ok` is set) and bottom points of each axis are written to a pdb file each frame. The residue numbers correspond to the axis numbers. When viewing this file with *rasmol*, use the command line option `-nmrpdb`, and type `set axis true` to display the reference axis.

Files

<code>-f</code>	<code>traj.xtc</code>	Input	Generic trajectory: xtc trr trj gro g96 pdb
<code>-s</code>	<code>topol.tpr</code>	Input	Structure+mass(db): tpr tpb tpa gro g96 pdb xml
<code>-n</code>	<code>index.ndx</code>	Input, Opt.	Index file
<code>-ol</code>	<code>bun_len.xvg</code>	Output	xvgr/xmgr file
<code>-od</code>	<code>bun_dist.xvg</code>	Output	xvgr/xmgr file
<code>-oz</code>	<code>bun_z.xvg</code>	Output	xvgr/xmgr file
<code>-ot</code>	<code>bun_tilt.xvg</code>	Output	xvgr/xmgr file

-otr	bun_tiltr.xvg	Output	xvgr/xmgr file
-otl	bun_tiltl.xvg	Output	xvgr/xmgr file
-ok	bun_kink.xvg	Output, Opt.	xvgr/xmgr file
-okr	bun_kinkr.xvg	Output, Opt.	xvgr/xmgr file
-okl	bun_kinkl.xvg	Output, Opt.	xvgr/xmgr file
-oa	axes.pdb	Output, Opt.	Protein data bank file

Other options

-h	bool	no	Print help info and quit
-nice	int	19	Set the nicelevel
-b	time	0	First frame (ps) to read from trajectory
-e	time	0	Last frame (ps) to read from trajectory
-dt	time	0	Only use frame when t MOD dt = first time (ps)
-tu	enum	ps	Time unit: ps, fs, ns, us, ms or s
-xvgr	bool	yes	Add specific codes (legends etc.) in the output xvg files for the xmgrace program
-na	int	0	Number of axes
-z	bool	no	Use the Z-axis as reference iso the average axis

D.12 g_chi

g_chi computes phi, psi, omega and chi dihedrals for all your amino acid backbone and sidechains. It can compute dihedral angle as a function of time, and as histogram distributions. The distributions (dihedral)(RESIDUE).xvg) are cumulative over all residues of each type.

If option `-corr` is given, the program will calculate dihedral autocorrelation functions. The function used is $C(t) = \langle \cos(\chi(\tau)) \cos(\chi(\tau+t)) \rangle$. The use of cosines rather than angles themselves, resolves the problem of periodicity. (Van der Spoel & Berendsen (1997), **Biophys. J.** **72**, 2032-2041). Separate files for each dihedral of each residue (corr(dihedral)(RESIDUE)(nresnr).xvg) are output, as well as a file containing the information for all residues (argument of `-corr`).

With option `-all`, the angles themselves as a function of time for each residue are printed to separate files (dihedral)(RESIDUE)(nresnr).xvg. These can be in radians or degrees.

A log file (argument `-g`) is also written. This contains

- information about the number of residues of each type.
- The NMR 3J coupling constants from the Karplus equation.
- a table for each residue of the number of transitions between rotamers per nanosecond, and the order parameter S2 of each dihedral.
- a table for each residue of the rotamer occupancy.

All rotamers are taken as 3-fold, except for omegas and chi-dihedrals to planar groups (i.e. chi2 of aromatics asp and asn, chi3 of glu and gln, and chi4 of arg), which are 2-fold. "rotamer 0" means that the dihedral was not in the core region of each rotamer. The width of the core region can be set with `-core_rotamer`

The S2 order parameters are also output to an xvg file (argument `-o`) and optionally as a pdb file with the S2 values as B-factor (argument `-p`). The total number of rotamer transitions per timestep (argument `-ot`), the number of transitions per rotamer (argument `-rt`), and the 3J couplings (argument `-jc`), can also be written to .xvg files.

If `-chi_prod` is set (and `maxchi > 0`), cumulative rotamers, e.g. $1+9(\chi_1-1)+3(\chi_2-1)+(\chi_3-1)$ (if the residue has three 3-fold dihedrals and `maxchi >= 3`) are calculated. As before, if any dihedral is not in the core region, the rotamer is taken to be 0. The occupancies of these cumulative rotamers (starting with rotamer 0) are written to the file that is the argument of `-cp`, and if the `-all` flag is given, the

rotamers as functions of time are written to `chiproduct(RESIDUE)(nresnr).xvg` and their occupancies to `histo-chiproduct(RESIDUE)(nresnr).xvg`.

The option `-r` generates a contour plot of the average omega angle as a function of the phi and psi angles, that is, in a Ramachandran plot the average omega angle is plotted using color coding.

Files

<code>-s</code>	<code>conf.gro</code>	Input	Generic structure: gro g96 pdb tpr tpb tpa xml
<code>-f</code>	<code>traj.xtc</code>	Input	Generic trajectory: xtc trr trj gro g96 pdb
<code>-o</code>	<code>order.xvg</code>	Output	xvgr/xmgr file
<code>-p</code>	<code>order.pdb</code>	Output, Opt.	Protein data bank file
<code>-ss</code>	<code>ssdump.dat</code>	Input, Opt.	Generic data file
<code>-jc</code>	<code>Jcoupling.xvg</code>	Output	xvgr/xmgr file
<code>-corr</code>	<code>dihcorr.xvg</code>	Output, Opt.	xvgr/xmgr file
<code>-g</code>	<code>chi.log</code>	Output	Log file
<code>-ot</code>	<code>dihtrans.xvg</code>	Output, Opt.	xvgr/xmgr file
<code>-oh</code>	<code>trhisto.xvg</code>	Output, Opt.	xvgr/xmgr file
<code>-rt</code>	<code>restrans.xvg</code>	Output, Opt.	xvgr/xmgr file
<code>-ch</code>	<code>chiprodhisto.xvg</code>	Output, Opt.	xvgr/xmgr file

Other options

<code>-h</code>	bool	no	Print help info and quit
<code>-nice</code>	int	19	Set the nicelevel
<code>-b</code>	time	0	First frame (ps) to read from trajectory
<code>-e</code>	time	0	Last frame (ps) to read from trajectory
<code>-dt</code>	time	0	Only use frame when $t \text{ MOD } dt = \text{first time (ps)}$
<code>-w</code>	bool	no	View output xvg, xpm, eps and pdb files
<code>-xvgr</code>	bool	yes	Add specific codes (legends etc.) in the output xvg files for the xmgrace program
<code>-r0</code>	int	1	starting residue
<code>-phi</code>	bool	no	Output for Phi dihedral angles
<code>-psi</code>	bool	no	Output for Psi dihedral angles
<code>-omega</code>	bool	no	Output for Omega dihedrals (peptide bonds)
<code>-rama</code>	bool	no	Generate Phi/Psi and Chi1/Chi2 ramachandran plots
<code>-viol</code>	bool	no	Write a file that gives 0 or 1 for violated Ramachandran angles
<code>-periodic</code>	bool	yes	Print dihedral angles modulo 360 degrees
<code>-all</code>	bool	no	Output separate files for every dihedral.
<code>-rad</code>	bool	no	in angle vs time files, use radians rather than degrees.
<code>-shift</code>	bool	no	Compute chemical shifts from Phi/Psi angles
<code>-binwidth</code>	int	1	bin width for histograms (degrees)
<code>-core_rotamer</code>	real	0.5	only the central <code>-core_rotamer*(360/multiplicity)</code> belongs to each rotamer (the rest is assigned to rotamer 0)
<code>-maxchi</code>	enum	0	calculate first ndih Chi dihedrals: 0, 1, 2, 3, 4, 5 or 6
<code>-normhisto</code>	bool	yes	Normalize histograms
<code>-ramomega</code>	bool	no	compute average omega as a function of phi/psi and plot it in an xpm plot
<code>-bfact</code>	real	-1	B-factor value for pdb file for atoms with no calculated dihedral order parameter
<code>-chi_prod</code>	bool	no	compute a single cumulative rotamer for each residue
<code>-HChi</code>	bool	no	Include dihedrals to sidechain hydrogens
<code>-bmax</code>	real	0	Maximum B-factor on any of the atoms that make up a dihedral, for the dihedral angle to be considered in the statistics. Applies to database work where a number of X-Ray structures is analyzed. <code>-bmax <= 0</code> means no limit.
<code>-acflen</code>	int	-1	Length of the ACF, default is half the number of frames

<code>-normalize</code>	bool	yes	Normalize ACF
	<code>-P</code>	enum	0 Order of Legendre polynomial for ACF (0 indicates none): 0, 1, 2 or 3
	<code>-fitfn</code>	enum	none Fit function: none, exp, aexp, exp_exp, vac, exp5, exp7 or exp9
	<code>-ncskip</code>	int	0 Skip N points in the output file of correlation functions
	<code>-beginfit</code>	real	0 Time where to begin the exponential fit of the correlation function
	<code>-endfit</code>	real	-1 Time where to end the exponential fit of the correlation function, -1 is till the end

- Produces MANY output files (up to about 4 times the number of residues in the protein, twice that if autocorrelation functions are calculated). Typically several hundred files are output.
- Phi and psi dihedrals are calculated in a non-standard way, using H-N-CA-C for phi instead of C(-)-N-CA-C, and N-CA-C-O for psi instead of N-CA-C-N(+). This causes (usually small) discrepancies with the output of other tools like `g_rama`.
- `-r0` option does not work properly
- Rotamers with multiplicity 2 are printed in `chi.log` as if they had multiplicity 3, with the 3rd (`g(+)`) always having probability 0

D.13 `g_cluster`

`g_cluster` can cluster structures with several different methods. Distances between structures can be determined from a trajectory or read from an XPM matrix file with the `-dm` option. RMS deviation after fitting or RMS deviation of atom-pair distances can be used to define the distance between structures.

single linkage: add a structure to a cluster when its distance to any element of the cluster is less than `cutoff`.

Jarvis Patrick: add a structure to a cluster when this structure and a structure in the cluster have each other as neighbors and they have a least `P` neighbors in common. The neighbors of a structure are the `M` closest structures or all structures within `cutoff`.

Monte Carlo: reorder the RMSD matrix using Monte Carlo.

diagonalization: diagonalize the RMSD matrix.

gromos: use algorithm as described in Daura *et al.* (*Angew. Chem. Int. Ed.* **1999**, 38, pp 236-240). Count number of neighbors using cut-off, take structure with largest number of neighbors with all its neighbors as cluster and eliminate it from the pool of clusters. Repeat for remaining structures in pool.

When the clustering algorithm assigns each structure to exactly one cluster (single linkage, Jarvis Patrick and gromos) and a trajectory file is supplied, the structure with the smallest average distance to the others or the average structure or all structures for each cluster will be written to a trajectory file. When writing all structures, separate numbered files are made for each cluster.

Two output files are always written:

`-o` writes the RMSD values in the upper left half of the matrix and a graphical depiction of the clusters in the lower right half. When `-minstruct = 1` the graphical depiction is black when two structures are in the same cluster. When `-minstruct > 1` different colors will be used for each cluster.

`-g` writes information on the options used and a detailed list of all clusters and their members.

Additionally, a number of optional output files can be written:

`-dist` writes the RMSD distribution.

`-ev` writes the eigenvectors of the RMSD matrix diagonalization.

`-sz` writes the cluster sizes.

`-tr` writes a matrix of the number transitions between cluster pairs.

-ntr writes the total number of transitions to or from each cluster.
 -clid writes the cluster number as a function of time.
 -cl writes average (with option -av) or central structure of each cluster or writes numbered files with cluster members for a selected set of clusters (with option -wcl, depends on -nst and -rmsmin).

Files

-f	traj.xtc	Input, Opt.	Generic trajectory: xtc trr trj gro g96 pdb
-s	topol.tpr	Input, Opt.	Structure+mass(db): tpr tpb tpa gro g96 pdb xml
-n	index.ndx	Input, Opt.	Index file
-dm	rmsd.xpm	Input, Opt.	X PixMap compatible matrix file
-o	rmsd-clust.xpm	Output	X PixMap compatible matrix file
-g	cluster.log	Output	Log file
-dist	rmsd-dist.xvg	Output, Opt.	xvgr/xmgr file
-ev	rmsd-eig.xvg	Output, Opt.	xvgr/xmgr file
-sz	clust-size.xvg	Output, Opt.	xvgr/xmgr file
-tr	clust-trans.xpm	Output, Opt.	X PixMap compatible matrix file
-nt	clust-trans.xvg	Output, Opt.	xvgr/xmgr file
-clid	clust-id.xvg	Output, Opt.	xvgr/xmgr file
-cl	clusters.pdb	Output, Opt.	Generic trajectory: xtc trr trj gro g96 pdb

Other options

-h	bool	no	Print help info and quit
-nice	int	19	Set the nicelevel
-b	time	0	First frame (ps) to read from trajectory
-e	time	0	Last frame (ps) to read from trajectory
-dt	time	0	Only use frame when $t \text{ MOD } dt = \text{first time (ps)}$
-tu	enum	ps	Time unit: ps, fs, ns, us, ms or s
-w	bool	no	View output xvg, xpm, eps and pdb files
-xvgr	bool	yes	Add specific codes (legends etc.) in the output xvg files for the xmgrace program
-dista	bool	no	Use RMSD of distances instead of RMS deviation
-nlevels	int	40	Discretize RMSD matrix in # levels
-cutoff	real	0.1	RMSD cut-off (nm) for two structures to be neighbor
-fit	bool	yes	Use least squares fitting before RMSD calculation
-max	real	-1	Maximum level in RMSD matrix
-skip	int	1	Only analyze every nr-th frame
-av	bool	no	Write average iso middle structure for each cluster
-wcl	int	0	Write all structures for first # clusters to numbered files
-nst	int	1	Only write all structures if more than # per cluster
-rmsmin	real	0	minimum rms difference with rest of cluster for writing structures
-method	enum	linkage	Method for cluster determination: linkage, jarvis-patrick, monte-carlo, diagonalization or gromos
-minstruct	int	1	Minimum number of structures in cluster for coloring in the xpm file
-binary	bool	no	Treat the RMSD matrix as consisting of 0 and 1, where the cut-off is given by -cutoff
-M	int	10	Number of nearest neighbors considered for Jarvis-Patrick algorithm, 0 is use cutoff
-P	int	3	Number of identical nearest neighbors required to form a cluster
-seed	int	1993	Random number seed for Monte Carlo clustering algorithm
-niter	int	10000	Number of iterations for MC
-kT	real	0.001	Boltzmann weighting factor for Monte Carlo optimization (zero turns off uphill steps)

D.14 g_clustsize

This program computes the size distributions of molecular/atomic clusters in the gas phase. The output is given in the form of a XPM file. The total number of clusters is written to a XVG file.

When the `-mol` option is given clusters will be made out of molecules rather than atoms, which allows clustering of large molecules. In this case an index file would still contain atom numbers or your calculation will die with a SEGV.

When velocities are present in your trajectory, the temperature of the largest cluster will be printed in a separate xvg file assuming that the particles are free to move. If you are using constraints, please correct the temperature. For instance water simulated with SHAKE or SETTLE will yield a temperature that is 1.5 times too low. You can compensate for this with the `-ndf` option. Remember to take the removal of center of mass motion into account.

The `-mc` option will produce an index file containing the atom numbers of the largest cluster.

Files

<code>-f</code>	<code>traj.xtc</code>	Input	Generic trajectory: xtc trr trj gro g96 pdb
<code>-s</code>	<code>topol.tpr</code>	Input, Opt.	Portable xdr run input file
<code>-n</code>	<code>index.ndx</code>	Input, Opt.	Index file
<code>-o</code>	<code>csize.xpm</code>	Output	X PixMap compatible matrix file
<code>-ow</code>	<code>csizew.xpm</code>	Output	X PixMap compatible matrix file
<code>-nc</code>	<code>nclust.xvg</code>	Output	xvgr/xmgr file
<code>-mc</code>	<code>maxclust.xvg</code>	Output	xvgr/xmgr file
<code>-ac</code>	<code>avclust.xvg</code>	Output	xvgr/xmgr file
<code>-h</code>	<code>histo-clust.xvg</code>	Output	xvgr/xmgr file
<code>-temp</code>	<code>temp.xvg</code>	Output, Opt.	xvgr/xmgr file
<code>-mcn</code>	<code>maxclust.ndx</code>	Output, Opt.	Index file

Other options

<code>-h</code>	bool	no	Print help info and quit
<code>-nice</code>	int	19	Set the nicelevel
<code>-b</code>	time	0	First frame (ps) to read from trajectory
<code>-e</code>	time	0	Last frame (ps) to read from trajectory
<code>-dt</code>	time	0	Only use frame when $t \text{ MOD } dt = \text{first time (ps)}$
<code>-tu</code>	enum	ps	Time unit: ps, fs, ns, us, ms or s
<code>-w</code>	bool	no	View output xvg, xpm, eps and pdb files
<code>-xvgr</code>	bool	yes	Add specific codes (legends etc.) in the output xvg files for the xmgrace program
<code>-cut</code>	real	0.35	Largest distance (nm) to be considered in a cluster
<code>-mol</code>	bool	no	Cluster molecules rather than atoms (needs tpr file)
<code>-pbc</code>	bool	yes	Use periodic boundary conditions
<code>-nskip</code>	int	0	Number of frames to skip between writing
<code>-nlevels</code>	int	20	Number of levels of grey in xpm output
<code>-ndf</code>	int	-1	Number of degrees of freedom of the entire system for temperature calculation. If not set the number of atoms times three is used.
<code>-rgblo</code>	vector	1 1 0	RGB values for the color of the lowest occupied cluster size
<code>-rgbhi</code>	vector	0 0 1	RGB values for the color of the highest occupied cluster size

D.15 g_confirms

`g_confirms` computes the root mean square deviation (RMSD) of two structures after LSQ fitting the second structure on the first one. The two structures do NOT need to have the same number of atoms, only the

two index groups used for the fit need to be identical. With `-name` only matching atom names from the selected groups will be used for the fit and RMSD calculation. This can be useful when comparing mutants of a protein.

The superimposed structures are written to file. In a `.pdb` file the two structures will be written as separate models (use `rasmol -nmrpdb`). Also in a `.pdb` file, B-factors calculated from the atomic MSD values can be written with `-bfac`.

Files

<code>-f1</code>	<code>conf1.gro</code>	Input	Structure+mass(db): tpr tpb tpa gro g96 pdb xml
<code>-f2</code>	<code>conf2.gro</code>	Input	Generic structure: gro g96 pdb tpr tpb tpa xml
<code>-o</code>	<code>fit.pdb</code>	Output	Generic structure: gro g96 pdb xml
<code>-n1</code>	<code>fit1.ndx</code>	Input, Opt.	Index file
<code>-n2</code>	<code>fit2.ndx</code>	Input, Opt.	Index file
<code>-no</code>	<code>match.ndx</code>	Output, Opt.	Index file

Other options

<code>-h</code>	bool	no	Print help info and quit
<code>-nice</code>	int	19	Set the nicelevel
<code>-w</code>	bool	no	View output xvg, xpm, eps and pdb files
<code>-one</code>	bool	no	Only write the fitted structure to file
<code>-mw</code>	bool	yes	Mass-weighted fitting and RMSD
<code>-pbc</code>	bool	no	Try to make molecules whole again
<code>-fit</code>	bool	yes	Do least squares superposition of the target structure to the reference
<code>-name</code>	bool	no	Only compare matching atom names
<code>-bfac</code>	bool	no	Output B-factors from atomic MSD values

D.16 *g_covar*

`g_covar` calculates and diagonalizes the (mass-weighted) covariance matrix. All structures are fitted to the structure in the structure file. When this is not a run input file periodicity will not be taken into account. When the fit and analysis groups are identical and the analysis is non mass-weighted, the fit will also be non mass-weighted.

The eigenvectors are written to a trajectory file (`-v`). When the same atoms are used for the fit and the covariance analysis, the reference structure for the fit is written first with `t=-1`. The average (or reference when `-ref` is used) structure is written with `t=0`, the eigenvectors are written as frames with the eigenvector number as timestamp.

The eigenvectors can be analyzed with `g_anaeig`.

Option `-ascii` writes the whole covariance matrix to an ASCII file. The order of the elements is: `x1x1, x1y1, x1z1, x1x2, ...`

Option `-xpm` writes the whole covariance matrix to an xpm file.

Option `-xpm` writes the atomic covariance matrix to an xpm file, i.e. for each atom pair the sum of the `xx`, `yy` and `zz` covariances is written.

Files

<code>-f</code>	<code>traj.xtc</code>	Input	Generic trajectory: xtc trr trj gro g96 pdb
<code>-s</code>	<code>topol.tpr</code>	Input	Structure+mass(db): tpr tpb tpa gro g96 pdb xml
<code>-n</code>	<code>index.ndx</code>	Input, Opt.	Index file
<code>-o</code>	<code>eigenval.xvg</code>	Output	xvgr/xmgr file
<code>-v</code>	<code>eigenvec.trr</code>	Output	Full precision trajectory: trr trj

-av	average.pdb	Output	Generic structure: gro g96 pdb xml
-l	covar.log	Output	Log file
-ascii	covar.dat	Output, Opt.	Generic data file
-xpm	covar.xpm	Output, Opt.	X PixMap compatible matrix file
-xpma	covara.xpm	Output, Opt.	X PixMap compatible matrix file

Other options

-h	bool	no	Print help info and quit
-nice	int	19	Set the nicelevel
-b	time	0	First frame (ps) to read from trajectory
-e	time	0	Last frame (ps) to read from trajectory
-dt	time	0	Only use frame when t MOD dt = first time (ps)
-tu	enum	ps	Time unit: ps, fs, ns, us, ms or s
-xvgr	bool	yes	Add specific codes (legends etc.) in the output xvgr files for the xmgrace program
-fit	bool	yes	Fit to a reference structure
-ref	bool	no	Use the deviation from the conformation in the structure file instead of from the average
-mwa	bool	no	Mass-weighted covariance analysis
-last	int	-1	Last eigenvector to write away (-1 is till the last)
-pbc	bool	yes	Apply corrections for periodic boundary conditions

D.17 g_density

Compute partial densities across the box, using an index file. Densities in kg/m^3 , number densities or electron densities can be calculated. For electron densities, a file describing the number of electrons for each type of atom should be provided using `-ei`. It should look like:

2

atomname = nrelectrons

atomname = nrelectrons

The first line contains the number of lines to read from the file. There should be one line for each unique atom name in your system. The number of electrons for each atom is modified by its atomic partial charge.

Files

-f	traj.xtc	Input	Generic trajectory: xtc trr trj gro g96 pdb
-n	index.ndx	Input, Opt.	Index file
-s	topol.tpr	Input	Generic run input: tpr tpb tpa xml
-ei	electrons.dat	Input, Opt.	Generic data file
-o	density.xvg	Output	xvgr/xmgr file

Other options

-h	bool	no	Print help info and quit
-nice	int	19	Set the nicelevel
-b	time	0	First frame (ps) to read from trajectory
-e	time	0	Last frame (ps) to read from trajectory
-dt	time	0	Only use frame when t MOD dt = first time (ps)
-w	bool	no	View output xvgr, xpm, eps and pdb files
-xvgr	bool	yes	Add specific codes (legends etc.) in the output xvgr files for the xmgrace program
-d	string	Z	Take the normal on the membrane in direction X, Y or Z.
-sl	int	50	Divide the box in #nr slices.

<code>-number</code>	bool	no	Calculate number density instead of mass density. Hydrogens are not counted!
<code>-ed</code>	bool	no	Calculate electron density instead of mass density
<code>-count</code>	bool	no	Only count atoms in slices, no densities. Hydrogens are not counted
<code>-ng</code>	int	1	Number of groups to compute densities of
<code>-symm</code>	bool	no	Symmetrize the density along the axis, with respect to the center. Useful for bilayers.
<code>-center</code>	bool	no	Shift the center of mass along the axis to zero. This means if your axis is Z and your box is bX, bY, bZ, the center of mass will be at bX/2, bY/2, 0.

- When calculating electron densities, atomnames are used instead of types. This is bad.
- When calculating number densities, atoms with names that start with H are not counted. This may be surprising if you use hydrogens with names like OP3.

D.18 *g_densmap*

g_densmap computes 2D number-density maps. It can make planar and axial-radial density maps. The output `.xpm` file can be visualized with for instance `xv` and can be converted to postscript with `xpm2ps`.

The default analysis is a 2-D number-density map for a selected group of atoms in the *x-z* plane. The averaging direction can be changed with the option `-aver`. The grid spacing is set with the option `-bin`. When `-n1` or `-n2` is non-zero, the grid size is set by this option. Box size fluctuations are properly taken into account.

When options `-amax` and `-rmax` are set, an axial-radial number-density map is made. Three groups should be supplied, the centers of mass of the first two groups define the axis, the third defines the analysis group. The axial direction goes from `-amax` to `+amax`, where the center is defined as the midpoint between the centers of mass and the positive direction goes from the first to the second center of mass. The radial direction goes from 0 to `rmax` or from `-rmax` to `+rmax` when the `-mirror` option has been set.

The normalization of the output is set with the `-unit` option. The default produces a true number density. Unit `nm-2` leaves out the normalization for the averaging or the angular direction. Option `count` produces the count for each grid cell. When you do not want the scale in the output to go from zero to the maximum density, you can set the maximum with the option `-dmax`.

Files

<code>-f</code>	<code>traj.xtc</code>	Input	Generic trajectory: <code>xtc trr trj gro g96 pdb</code>
<code>-s</code>	<code>topol.tpr</code>	Input, Opt.	Structure+mass(db): <code>tpr tpb tpa gro g96 pdb xml</code>
<code>-n</code>	<code>index.ndx</code>	Input, Opt.	Index file
<code>-o</code>	<code>densmap.xpm</code>	Output	X PixMap compatible matrix file

Other options

<code>-h</code>	bool	no	Print help info and quit
<code>-nice</code>	int	19	Set the nicelevel
<code>-b</code>	time	0	First frame (ps) to read from trajectory
<code>-e</code>	time	0	Last frame (ps) to read from trajectory
<code>-dt</code>	time	0	Only use frame when $t \text{ MOD } dt = \text{first time (ps)}$
<code>-w</code>	bool	no	View output <code>xvg</code> , <code>xpm</code> , <code>eps</code> and <code>pdb</code> files
<code>-bin</code>	real	0.02	Grid size (nm)
<code>-aver</code>	enum	z	The direction to average over: <code>z</code> , <code>y</code> or <code>x</code>
<code>-n1</code>	int	0	Number of grid cells in the first direction
<code>-n2</code>	int	0	Number of grid cells in the second direction

-amax	real	0	Maximum axial distance from the center
-rmax	real	0	Maximum radial distance
-mirror	bool	no	Add the mirror image below the axial axis
-unit	enum	nm-3	Unit for the output: nm-3, nm-2 or count
-dmin	real	0	Minimum density in output
-dmax	real	0	Maximum density in output (0 means calculate it)

D.19 g_dielectric

dielectric calculates frequency dependent dielectric constants from the autocorrelation function of the total dipole moment in your simulation. This ACF can be generated by g_dipoles. For an estimate of the error you can run g_statistics on the ACF, and use the output thus generated for this program. The functional forms of the available functions are:

One parameter : $y = \text{Exp}[-a1 x]$ Two parameters : $y = a2 \text{Exp}[-a1 x]$ Three parameter: $y = a2 \text{Exp}[-a1 x] + (1 - a2) \text{Exp}[-a3 x]$ Startvalues for the fit procedure can be given on the commandline. It is also possible to fix parameters at their start value, use -fix with the number of the parameter you want to fix.

Three output files are generated, the first contains the ACF, an exponential fit to it with 1, 2 or 3 parameters, and the numerical derivative of the combination data/fit. The second file contains the real and imaginary parts of the frequency-dependent dielectric constant, the last gives a plot known as the Cole-Cole plot, in which the imaginary component is plotted as a function of the real component. For a pure exponential relaxation (Debye relaxation) the latter plot should be one half of a circle

Files

-f	Mtot.svg	Input	xvgr/xmgr file
-d	deriv.svg	Output	xvgr/xmgr file
-o	epsw.svg	Output	xvgr/xmgr file
-c	cole.svg	Output	xvgr/xmgr file

Other options

-h	bool	no	Print help info and quit
-nice	int	19	Set the nicelevel
-b	time	0	First frame (ps) to read from trajectory
-e	time	0	Last frame (ps) to read from trajectory
-dt	time	0	Only use frame when $t \text{ MOD } dt = \text{first time (ps)}$
-w	bool	no	View output xvg, xpm, eps and pdb files
-xvgr	bool	yes	Add specific codes (legends etc.) in the output xvg files for the xmgrace program
-fft	bool	no	use fast fourier transform for correlation function
-x1	bool	yes	use first column as X axis rather than first data set
-eint	real	5	Time were to end the integration of the data and start to use the fit
-bfit	real	5	Begin time of fit
-efit	real	500	End time of fit
-tail	real	500	Length of function including data and tail from fit
-A	real	0.5	Start value for fit parameter A
-tau1	real	10	Start value for fit parameter tau1
-tau2	real	1	Start value for fit parameter tau2
-eps0	real	80	Epsilon 0 of your liquid
-epsRF	real	78.5	Epsilon of the reaction field used in your simulation. A value of 0 means infinity.
-fix	int	0	Fix parameters at their start values, A (2), tau1 (1), or tau2 (4)
-ffn	enum	none	Fit function: none, exp, aexp, exp_exp, vac, exp5, exp7 or exp9

`-nsmooth` int 3 Number of points for smoothing

D.20 *g_dih*

g_dih can do two things. The default is to analyze dihedral transitions by merely computing all the dihedral angles defined in your topology for the whole trajectory. When a dihedral flips over to another minimum an angle/time plot is made.

The other option is to discretize the dihedral space into a number of bins, and group each conformation in dihedral space in the appropriate bin. The output is then given as a number of dihedral conformations sorted according to occupancy.

Files

<code>-f</code>	<code>traj.xtc</code>	Input	Generic trajectory: xtc trr trj gro g96 pdb
<code>-s</code>	<code>topol.tpr</code>	Input	Generic run input: tpr tpb tpa xml
<code>-o</code>	<code>hello.out</code>	Output	Generic output file

Other options

<code>-h</code>	bool	no	Print help info and quit
<code>-nice</code>	int	19	Set the nicelevel
<code>-b</code>	time	0	First frame (ps) to read from trajectory
<code>-e</code>	time	0	Last frame (ps) to read from trajectory
<code>-dt</code>	time	0	Only use frame when $t \text{ MOD } dt = \text{first time (ps)}$
<code>-w</code>	bool	no	View output xvg, xpm, eps and pdb files
<code>-sa</code>	bool	no	Perform cluster analysis in dihedral space instead of analysing dihedral transitions.
<code>-mult</code>	int	-1	multiplicity for dihedral angles (by default read from topology)

D.21 *g_dipoles*

g_dipoles computes the total dipole plus fluctuations of a simulation system. From this you can compute e.g. the dielectric constant for low dielectric media

The file `Mtot.xvg` contains the total dipole moment of a frame, the components as well as the norm of the vector. The file `aver.xvg` contains $\langle \text{orMuor}^2 \rangle$ and $\langle \text{orMuor} \rangle^2$ during the simulation. The file `dipdist.xvg` contains the distribution of dipole moments during the simulation. The `mu_max` is used as the highest value in the distribution graph.

Furthermore the dipole autocorrelation function will be computed when option `-corr` is used. The output file name is given with the `-c` option. The correlation functions can be averaged over all molecules (`mol`), plotted per molecule separately (`molsep`) or it can be computed over the total dipole moment of the simulation box (`total`).

Option `-g` produces a plot of the distance dependent Kirkwood G-factor, as well as the average cosine of the angle between the dipoles as a function of the distance. The plot also includes `gOO` and `hOO` according to Nymand & Linse, JCP 112 (2000) pp 6386-6395. In the same plot we also include the energy per scale computed by taking the inner product of the dipoles divided by the distance to the third power.

EXAMPLES

```
g_dipoles -corr mol -P1 -o dip_sqr -mu 2.273 -mumax 5.0 -nofft
```

This will calculate the autocorrelation function of the molecular dipoles using a first order Legendre polynomial of the angle of the dipole vector and itself a time t later. For this calculation 1001 frames will be

used. Further the dielectric constant will be calculated using an epsilonRF of infinity (default), temperature of 300 K (default) and an average dipole moment of the molecule of 2.273 (SPC). For the distribution function a maximum of 5.0 will be used.

Files

-enx	ener.edr	Input, Opt.	Generic energy: edr ene
-f	traj.xtc	Input	Generic trajectory: xtc trr trj gro g96 pdb
-s	topol.tpr	Input	Generic run input: tpr tpb tpa xml
-n	index.ndx	Input, Opt.	Index file
-o	Mtot.xvg	Output	xvgr/xmgr file
-eps	epsilon.xvg	Output	xvgr/xmgr file
-a	aver.xvg	Output	xvgr/xmgr file
-d	dipdist.xvg	Output	xvgr/xmgr file
-c	dipcorr.xvg	Output, Opt.	xvgr/xmgr file
-g	gkr.xvg	Output, Opt.	xvgr/xmgr file
-adip	adip.xvg	Output, Opt.	xvgr/xmgr file
-dip3d	dip3d.xvg	Output, Opt.	xvgr/xmgr file
-cos	cosaver.xvg	Output, Opt.	xvgr/xmgr file
-q	quadrupole.xvg	Output, Opt.	xvgr/xmgr file
-slab	slab.xvg	Output, Opt.	xvgr/xmgr file

Other options

-h	bool	no	Print help info and quit
-nice	int	19	Set the nicelevel
-b	time	0	First frame (ps) to read from trajectory
-e	time	0	Last frame (ps) to read from trajectory
-dt	time	0	Only use frame when t MOD dt = first time (ps)
-w	bool	no	View output xvg, xpm, eps and pdb files
-xvgr	bool	yes	Add specific codes (legends etc.) in the output xvg files for the xmgrace program
-mu	real	-1	dipole of a single molecule (in Debye)
-mumax	real	5	max dipole in Debye (for histogram)
-epsilonRF	real	0	epsilon of the reaction field used during the simulation, needed for dielectric constant calculation. WARNING: 0.0 means infinity (default)
-skip	int	0	Skip steps in the output (but not in the computations)
-temp	real	300	Average temperature of the simulation (needed for dielectric constant calculation)
-corr	enum	none	Correlation function to calculate: none, mol, molsep or total
-pairs	bool	yes	Calculate orcos thetaor between all pairs of molecules. May be slow
-axis	string	Z	Take the normal on the computational box in direction X, Y or Z.
-sl	int	10	Divide the box in #nr slices.
-gkratom	int	0	Use the n-th atom of a molecule (starting from 1) to calculate the distance between molecules rather than the center of charge (when 0) in the calculation of distance dependent Kirkwood factors
-acflen	int	-1	Length of the ACF, default is half the number of frames
-normalize	bool	yes	Normalize ACF
-P	enum	0	Order of Legendre polynomial for ACF (0 indicates none): 0, 1, 2 or 3
-fitfn	enum	none	Fit function: none, exp, aexp, exp_exp, vac, exp5, exp7 or exp9
-ncskip	int	0	Skip N points in the output file of correlation functions
-beginfit	real	0	Time where to begin the exponential fit of the correlation function
-endfit	real	-1	Time where to end the exponential fit of the correlation function, -1 is till the end

D.22 *g_disre*

g_disre computes violations of distance restraints. If necessary all protons can be added to a protein molecule using the *protonate* program.

The program allways computes the instantaneous violations rather than time-averaged, because this analysis is done from a trajectory file afterwards it does not make sense to use time averaging. However, the time averaged values per restraint are given in the log file.

An index file may be used to select specific restraints for printing.

When the optional `-q` flag is given a pdb file coloured by the amount of average violations.

When the `-c` option is given, an index file will be read containing the frames in your trajectory corresponding to the clusters (defined in another manner) that you want to analyze. For these clusters the program will compute average violations using the *thisd* power averaging algorithm and print them in the log file.

Files

<code>-s</code>	<code>topol.tpr</code>	Input	Generic run input: tpr tpb tpa xml
<code>-f</code>	<code>traj.xtc</code>	Input	Generic trajectory: xtc trr trj gro g96 pdb
<code>-ds</code>	<code>drsum.xvg</code>	Output	xvgr/xmgr file
<code>-da</code>	<code>draver.xvg</code>	Output	xvgr/xmgr file
<code>-dn</code>	<code>drnum.xvg</code>	Output	xvgr/xmgr file
<code>-dm</code>	<code>drmax.xvg</code>	Output	xvgr/xmgr file
<code>-dr</code>	<code>restr.xvg</code>	Output	xvgr/xmgr file
<code>-l</code>	<code>disres.log</code>	Output	Log file
<code>-n</code>	<code>viol.ndx</code>	Input, Opt.	Index file
<code>-q</code>	<code>viol.pdb</code>	Output, Opt.	Protein data bank file
<code>-c</code>	<code>clust.ndx</code>	Input, Opt.	Index file

Other options

<code>-h</code>	bool	no	Print help info and quit
<code>-nice</code>	int	19	Set the nicelevel
<code>-b</code>	time	0	First frame (ps) to read from trajectory
<code>-e</code>	time	0	Last frame (ps) to read from trajectory
<code>-dt</code>	time	0	Only use frame when $t \text{ MOD } dt = \text{first time (ps)}$
<code>-w</code>	bool	no	View output xvg, xpm, eps and pdb files
<code>-xvgr</code>	bool	yes	Add specific codes (legends etc.) in the output xvg files for the xmgrace program
<code>-ntop</code>	int	0	Number of large violations that are stored in the log file every step

D.23 *g_dist*

g_dist can calculate the distance between the centers of mass of two groups of atoms as a function of time. The total distance and its x, y and z components are plotted.

Or when `-dist` is set, print all the atoms in group 2 that are closer than a certain distance to the center of mass of group 1.

Other programs that calculate distances are *g_mindist* and *g_bond*.

Files

<code>-f</code>	<code>traj.xtc</code>	Input	Generic trajectory: xtc trr trj gro g96 pdb
<code>-s</code>	<code>topol.tpr</code>	Input	Generic run input: tpr tpb tpa xml
<code>-n</code>	<code>index.ndx</code>	Input, Opt.	Index file
<code>-o</code>	<code>dist.xvg</code>	Output, Opt.	xvgr/xmgr file

Other options

-h	bool	no	Print help info and quit
-nice	int	19	Set the nicelevel
-b	time	0	First frame (ps) to read from trajectory
-e	time	0	Last frame (ps) to read from trajectory
-dt	time	0	Only use frame when $t \text{ MOD } dt = \text{first time (ps)}$
-xvgr	bool	yes	Add specific codes (legends etc.) in the output xvg files for the xmgrace program
-dist	real	0	Print all atoms in group 2 closer than dist to the center of mass of group 1

D.24 g_dyndom

`g_dyndom` reads a pdb file output from DynDom <http://md.chem.rug.nl/~steve/DynDom/dyndom.home.html>. It reads the coordinates, and the coordinates of the rotation axis furthermore it reads an index file containing the domains. Furthermore it takes the first and last atom of the arrow file as command line arguments (head and tail) and finally it takes the translation vector (given in DynDom info file) and the angle of rotation (also as command line arguments). If the angle determined by DynDom is given, one should be able to recover the second structure used for generating the DynDom output. Because of limited numerical accuracy this should be verified by computing an all-atom RMSD (using `g_confirms`) rather than by file comparison (using `diff`).

The purpose of this program is to interpolate and extrapolate the rotation as found by DynDom. As a result unphysical structures with long or short bonds, or overlapping atoms may be produced. Visual inspection, and energy minimization may be necessary to validate the structure.

Files

-f	dyndom.pdb	Input	Protein data bank file
-o	rotated.xtc	Output	Generic trajectory: xtc trr trj gro g96 pdb
-n	domains.ndx	Input	Index file

Other options

-h	bool	no	Print help info and quit
-nice	int	0	Set the nicelevel
-firstangle	real	0	Angle of rotation about rotation vector
-lastangle	real	0	Angle of rotation about rotation vector
-nframe	int	11	Number of steps on the pathway
-maxangle	real	0	DymDom dtermined angle of rotation about rotation vector
-trans	real	0	Translation (Aangstroem) along rotation vector (see DynDom info file)
-head	vector	0 0 0	First atom of the arrow vector
-tail	vector	0 0 0	Last atom of the arrow vector

D.25 g_enemat

`g_enemat` extracts an energy matrix from the energy file (`-f`). With `-groups` a file must be supplied with on each line a group of atoms to be used. For these groups matrix of interaction energies will be extracted from the energy file by looking for energy groups with names corresponding to pairs of groups of atoms. E.g. if your `-groups` file contains:

```
2
Protein
```

SOL

then energy groups with names like 'Coul-SR:Protein-SOL' and 'LJ:Protein-SOL' are expected in the energy file (although *g_enemat* is most useful if many groups are analyzed simultaneously). Matrices for different energy types are written out separately, as controlled by the `-[no]coul`, `-[no]coulr`, `-[no]coul14`, `-[no]lj`, `-[no]lj14`, `-[no]bham` and `-[no]free` options. Finally, the total interaction energy per group can be calculated (`-etot`).

An approximation of the free energy can be calculated using: $E(\text{free}) = E_0 + kT \log(\langle \exp((E-E_0)/kT) \rangle)$, where ' $\langle \rangle$ ' stands for time-average. A file with reference free energies can be supplied to calculate the free energy difference with some reference state. Group names (e.g. residue names) in the reference file should correspond to the group names as used in the `-groups` file, but a appended number (e.g. residue number) in the `-groups` will be ignored in the comparison.

Files

<code>-f</code>	<code>ener.edr</code>	Input, Opt.	Generic energy: <code>edr ene</code>
<code>-groups</code>	<code>groups.dat</code>	Input	Generic data file
<code>-eref</code>	<code>eref.dat</code>	Input, Opt.	Generic data file
<code>-emat</code>	<code>emat.xpm</code>	Output	X PixMap compatible matrix file
<code>-etot</code>	<code>energy.xvg</code>	Output	<code>xvgr/xmgr</code> file

Other options

<code>-h</code>	bool	no	Print help info and quit
<code>-nice</code>	int	19	Set the nicelevel
<code>-b</code>	time	0	First frame (ps) to read from trajectory
<code>-e</code>	time	0	Last frame (ps) to read from trajectory
<code>-dt</code>	time	0	Only use frame when $t \text{ MOD } dt = \text{first time (ps)}$
<code>-w</code>	bool	no	View output <code>xvg</code> , <code>xpm</code> , <code>eps</code> and <code>pdb</code> files
<code>-xvgr</code>	bool	yes	Add specific codes (legends etc.) in the output <code>xvg</code> files for the <code>xmgrace</code> program
<code>-sum</code>	bool	no	Sum the energy terms selected rather than display them all
<code>-skip</code>	int	0	Skip number of frames between data points
<code>-mean</code>	bool	yes	with <code>-groups</code> extracts matrix of mean energies in stead of matrix for each timestep
<code>-nlevels</code>	int	20	number of levels for matrix colors
<code>-max</code>	real	1e+20	max value for energies
<code>-min</code>	real	-1e+20	min value for energies
<code>-coul</code>	bool	yes	extract Coulomb SR energies
<code>-coulr</code>	bool	no	extract Coulomb LR energies
<code>-coul14</code>	bool	no	extract Coulomb 1-4 energies
<code>-lj</code>	bool	yes	extract Lennard-Jones SR energies
<code>-lj</code>	bool	no	extract Lennard-Jones LR energies
<code>-lj14</code>	bool	no	extract Lennard-Jones 1-4 energies
<code>-bhamsr</code>	bool	no	extract Buckingham SR energies
<code>-bhamlr</code>	bool	no	extract Buckingham LR energies
<code>-free</code>	bool	yes	calculate free energy
<code>-temp</code>	real	300	reference temperature for free energy calculation

D.26 *g_energy*

g_energy extracts energy components or distance restraint data from an energy file. The user is prompted to interactively select the energy terms she wants.

Average and RMSD are calculated with full precision from the simulation (see printed manual). Drift is calculated by performing a LSQ fit of the data to a straight line. Total drift is drift multiplied by total time.

When the `-viol` option is set, the time averaged violations are plotted and the running time-averaged and instantaneous sum of violations are recalculated. Additionally running time-averaged and instantaneous distances between selected pairs can be plotted with the `-pairs` option.

Options `-ora`, `-ort`, `-oda`, `-odr` and `-odt` are used for analyzing orientation restraint data. The first two options plot the orientation, the last three the deviations of the orientations from the experimental values. The options that end on an 'a' plot the average over time as a function of restraint. The options that end on a 't' prompt the user for restraint label numbers and plot the data as a function of time. Option `-odr` plots the RMS deviation as a function of restraint. When the run used time or ensemble averaged orientation restraints, option `-orinst` can be used to analyse the instantaneous, not ensemble-averaged orientations and deviations instead of the time and ensemble averages.

Option `-oten` plots the eigenvalues of the molecular order tensor for each orientation restraint experiment. With option `-ovec` also the eigenvectors are plotted.

With `-fee` an estimate is calculated for the free-energy difference with an ideal gas state:

$$\Delta A = A(N,V,T) - A_{\text{idgas}}(N,V,T) = kT \ln \langle e^{\hat{U}_{\text{pot}}/kT} \rangle$$

$$\Delta G = G(N,p,T) - G_{\text{idgas}}(N,p,T) = kT \ln \langle e^{\hat{U}_{\text{pot}}/kT} \rangle$$

where k is Boltzmann's constant, T is set by `-fetemp` and the average is over the ensemble (or time in a trajectory). Note that this is in principle only correct when averaging over the whole (Boltzmann) ensemble and using the potential energy. This also allows for an entropy estimate using:

$$\Delta S(N,V,T) = S(N,V,T) - S_{\text{idgas}}(N,V,T) = (\langle U_{\text{pot}} \rangle - \Delta A)/T$$

$$\Delta S(N,p,T) = S(N,p,T) - S_{\text{idgas}}(N,p,T) = (\langle U_{\text{pot}} \rangle + pV - \Delta G)/T$$

When a second energy file is specified (`-f2`), a free energy difference is calculated $dF = -kT \ln \langle e^{\Delta(E_B - EA)/kT} \rangle$, where EA and EB are the energies from the first and second energy files, and the average is over the ensemble A . **NOTE** that the energies must both be calculated from the same trajectory.

Files

<code>-f</code>	<code>ener.edr</code>	Input	Generic energy: <code>edr ene</code>
<code>-f2</code>	<code>ener.edr</code>	Input, Opt.	Generic energy: <code>edr ene</code>
<code>-s</code>	<code>topol.tpr</code>	Input, Opt.	Generic run input: <code>tpr tpb tpa xml</code>
<code>-o</code>	<code>energy.xvg</code>	Output	<code>xvgr/xmgr</code> file
<code>-viol</code>	<code>violaver.xvg</code>	Output, Opt.	<code>xvgr/xmgr</code> file
<code>-pairs</code>	<code>pairs.xvg</code>	Output, Opt.	<code>xvgr/xmgr</code> file
<code>-ora</code>	<code>orienta.xvg</code>	Output, Opt.	<code>xvgr/xmgr</code> file
<code>-ort</code>	<code>orientt.xvg</code>	Output, Opt.	<code>xvgr/xmgr</code> file
<code>-oda</code>	<code>orideva.xvg</code>	Output, Opt.	<code>xvgr/xmgr</code> file
<code>-odr</code>	<code>oridevr.xvg</code>	Output, Opt.	<code>xvgr/xmgr</code> file
<code>-odt</code>	<code>oridevt.xvg</code>	Output, Opt.	<code>xvgr/xmgr</code> file
<code>-oten</code>	<code>oriten.xvg</code>	Output, Opt.	<code>xvgr/xmgr</code> file
<code>-corr</code>	<code>enecorr.xvg</code>	Output, Opt.	<code>xvgr/xmgr</code> file
<code>-vis</code>	<code>visco.xvg</code>	Output, Opt.	<code>xvgr/xmgr</code> file
<code>-ravg</code>	<code>runavgdf.xvg</code>	Output, Opt.	<code>xvgr/xmgr</code> file

Other options

<code>-h</code>	bool	no	Print help info and quit
<code>-nice</code>	int	19	Set the nicelevel
<code>-b</code>	time	0	First frame (ps) to read from trajectory
<code>-e</code>	time	0	Last frame (ps) to read from trajectory
<code>-w</code>	bool	no	View output <code>xvg</code> , <code>xpm</code> , <code>eps</code> and <code>pdb</code> files
<code>-xvgr</code>	bool	yes	Add specific codes (legends etc.) in the output <code>xvg</code> files for the <code>xmgrace</code> program

-fee	bool	no	Do a free energy estimate
-fetemp	real	300	Reference temperature for free energy calculation
-zero	real	0	Subtract a zero-point energy
-sum	bool	no	Sum the energy terms selected rather than display them all
-dp	bool	no	Print energies in high precision
-mutot	bool	no	Compute the total dipole moment from the components
-skip	int	0	Skip number of frames between data points
-aver	bool	no	Print also the $X_{1,t}$ and $\sigma_{1,t}$, only if only 1 energy is requested
-nmol	int	1	Number of molecules in your sample: the energies are divided by this number
-ndf	int	3	Number of degrees of freedom per molecule. Necessary for calculating the heat capacity
-fluc	bool	no	Calculate autocorrelation of energy fluctuations rather than energy itself
-orinst	bool	no	Analyse instantaneous orientation data
-ovec	bool	no	Also plot the eigenvectors with -oten
-acflen	int	-1	Length of the ACF, default is half the number of frames
-normalize	bool	yes	Normalize ACF
-P	enum	0	Order of Legendre polynomial for ACF (0 indicates none): 0, 1, 2 or 3
-fitfn	enum	none	Fit function: none, exp, aexp, exp_exp, vac, exp5, exp7 or exp9
-ncskip	int	0	Skip N points in the output file of correlation functions
-beginfit	real	0	Time where to begin the exponential fit of the correlation function
-endfit	real	-1	Time where to end the exponential fit of the correlation function, -1 is till the end

D.27 *g_filter*

g_filter performs frequency filtering on a trajectory. The filter shape is $\cos(\pi t/A) + 1$ from $-A$ to $+A$, where A is given by the option `-nf` times the time step in the input trajectory. This filter reduces fluctuations with period A by 85%, with period $2*A$ by 50% and with period $3*A$ by 17% for low-pass filtering. Both a low-pass and high-pass filtered trajectory can be written.

Option `-ol` writes a low-pass filtered trajectory. A frame is written every `nf` input frames. This ratio of filter length and output interval ensures a good suppression of aliasing of high-frequency motion, which is useful for making smooth movies. Also averages of properties which are linear in the coordinates are preserved, since all input frames are weighted equally in the output. When all frames are needed, use the `-all` option.

Option `-oh` writes a high-pass filtered trajectory. The high-pass filtered coordinates are added to the coordinates from the structure file. When using high-pass filtering use `-fit` or make sure you use a trajectory which has been fitted on the coordinates in the structure file.

Files

-f	traj.xtc	Input	Generic trajectory: xtc trr trj gro g96 pdb
-s	topol.tpr	Input, Opt.	Structure+mass(db): tpr tpb tpa gro g96 pdb xml
-n	index.ndx	Input, Opt.	Index file
-ol	lowpass.xtc	Output, Opt.	Generic trajectory: xtc trr trj gro g96 pdb
-oh	highpass.xtc	Output, Opt.	Generic trajectory: xtc trr trj gro g96 pdb

Other options

-h	bool	no	Print help info and quit
-nice	int	19	Set the nicelevel
-b	time	0	First frame (ps) to read from trajectory
-e	time	0	Last frame (ps) to read from trajectory

-dt	time	0	Only use frame when $t \text{ MOD } dt = \text{first time (ps)}$
-w	bool	no	View output xvg, xpm, eps and pdb files
-nf	int	10	Sets the filter length as well as the output interval for low-pass filtering
-all	bool	no	Write all low-pass filtered frames
-nojump	bool	yes	Remove jumps of atoms across the box
-fit	bool	no	Fit all frames to a reference structure

D.28 g_gyrate

`g_gyrate` computes the radius of gyration of a group of atoms and the radii of gyration about the x, y and z axes, as a function of time. The atoms are explicitly mass weighted. With the `-nmol` option the radius of gyration will be calculated for multiple molecules by splitting the analysis group in equally sized parts.

Files

-f	traj.xtc	Input	Generic trajectory: xtc trr trj gro g96 pdb
-s	topol.tpr	Input	Structure+mass(db): tpr tpb tpa gro g96 pdb xml
-n	index.ndx	Input, Opt.	Index file
-o	gyrate.xvg	Output	xvgr/xmgr file
-acf	moi-acf.xvg	Output, Opt.	xvgr/xmgr file

Other options

-h	bool	no	Print help info and quit
-nice	int	19	Set the nicelevel
-b	time	0	First frame (ps) to read from trajectory
-e	time	0	Last frame (ps) to read from trajectory
-dt	time	0	Only use frame when $t \text{ MOD } dt = \text{first time (ps)}$
-w	bool	no	View output xvg, xpm, eps and pdb files
-xvgr	bool	yes	Add specific codes (legends etc.) in the output xvg files for the xmgrace program
-nmol	int	1	The number of molecules to analyze
-q	bool	no	Use absolute value of the charge of an atom as weighting factor instead of mass
-p	bool	no	Calculate the radii of gyration about the principal axes.
-moi	bool	no	Calculate the moments of inertia (defined by the principal axes).
-acflen	int	-1	Length of the ACF, default is half the number of frames
-normalize	bool	yes	Normalize ACF
-P	enum	0	Order of Legendre polynomial for ACF (0 indicates none): 0, 1, 2 or 3
-fitfn	enum	none	Fit function: none, exp, aexp, exp_exp, vac, exp5, exp7 or exp9
-ncskip	int	0	Skip N points in the output file of correlation functions
-beginfit	real	0	Time where to begin the exponential fit of the correlation function
-endfit	real	-1	Time where to end the exponential fit of the correlation function, -1 is till the end

D.29 g_h2order

Compute the orientation of water molecules with respect to the normal of the box. The program determines the average cosine of the angle between the dipole moment of water and an axis of the box. The box is divided in slices and the average orientation per slice is printed. Each water molecule is assigned to a slice, per time frame, based on the position of the oxygen. When `-nm` is used the angle between the water dipole and the axis from the center of mass to the oxygen is calculated instead of the angle between the dipole and a box axis.

Files

-f	traj.xtc	Input	Generic trajectory: xtc trr trj gro g96 pdb
-n	index.ndx	Input	Index file
-nm	index.ndx	Input, Opt.	Index file
-s	topol.tpr	Input	Generic run input: tpr tpb tpa xml
-o	order.xvg	Output	xvgr/xmgr file

Other options

-h	bool	no	Print help info and quit
-nice	int	19	Set the nicelevel
-b	time	0	First frame (ps) to read from trajectory
-e	time	0	Last frame (ps) to read from trajectory
-dt	time	0	Only use frame when $t \text{ MOD } dt = \text{first time (ps)}$
-w	bool	no	View output xvg, xpm, eps and pdb files
-xvgr	bool	yes	Add specific codes (legends etc.) in the output xvg files for the xmgrace program
-d	string	Z	Take the normal on the membrane in direction X, Y or Z.
-sl	int	0	Calculate order parameter as function of boxlength, dividing the box in #nr slices.

- The program assigns whole water molecules to a slice, based on the firstatom of three in the index file group. It assumes an order O,H,H.Name is not important, but the order is. If this demand is not met, assigning molecules to slices is different.

D.30 g_hbond

g_hbond computes and analyzes hydrogen bonds. Hydrogen bonds are determined based on cutoffs for the angle Acceptor - Donor - Hydrogen (zero is extended) and the distance Hydrogen - Acceptor. OH and NH groups are regarded as donors, O is an acceptor always, N is an acceptor by default, but this can be switched using *-nitacc*. Dummy hydrogen atoms are assumed to be connected to the first preceding non-hydrogen atom.

You need to specify two groups for analysis, which must be either identical or non-overlapping. All hydrogen bonds between the two groups are analyzed.

If you set *-shell*, you will be asked for an additional index group which should contain exactly one atom. In this case, only hydrogen bonds between atoms within the shell distance from the one atom are considered.

```
[ selected ]
20 21 24
25 26 29
1 3 6
```

Note that the triplets need not be on separate lines. Each atom triplet specifies a hydrogen bond to be analyzed, note also that no check is made for the types of atoms.

-ins turns on computing solvent insertion into hydrogen bonds. In this case an additional group must be selected, specifying the solvent molecules.

Output:

- num: number of hydrogen bonds as a function of time.
- ac: average over all autocorrelations of the existence functions (either 0 or 1) of all hydrogen bonds.
- dist: distance distribution of all hydrogen bonds.
- ang: angle distribution of all hydrogen bonds.

-hx: the number of n-n+i hydrogen bonds as a function of time where n and n+i stand for residue numbers and i ranges from 0 to 6. This includes the n-n+3, n-n+4 and n-n+5 hydrogen bonds associated with helices in proteins.

-hbn: all selected groups, donors, hydrogens and acceptors for selected groups, all hydrogen bonded atoms from all groups and all solvent atoms involved in insertion.

-hbm: existence matrix for all hydrogen bonds over all frames, this also contains information on solvent insertion into hydrogen bonds. Ordering is identical to that in -hbn index file.

-dan: write out the number of donors and acceptors analyzed for each timeframe. This is especially useful when using -shell.

-nhbdist: compute the number of HBonds per hydrogen in order to compare results to Raman Spectroscopy.

Note: options -ac, -life, -hbn and -hbm require an amount of memory proportional to the total numbers of donors times the total number of acceptors in the selected group(s).

Files

-f	traj.xtc	Input	Generic trajectory: xtc trr trj gro g96 pdb
-s	topol.tpr	Input	Generic run input: tpr tpb tpa xml
-n	index.ndx	Input, Opt.	Index file
-num	hnum.xvg	Output	xvgr/xmgr file
-ac	hbac.xvg	Output, Opt.	xvgr/xmgr file
-dist	hbdist.xvg	Output, Opt.	xvgr/xmgr file
-ang	hbang.xvg	Output, Opt.	xvgr/xmgr file
-hx	hbhelix.xvg	Output, Opt.	xvgr/xmgr file
-hbn	hbond.ndx	Output, Opt.	Index file
-hbm	hmap.xpm	Output, Opt.	X PixMap compatible matrix file
-don	donor.xvg	Output, Opt.	xvgr/xmgr file
-dan	danum.xvg	Output, Opt.	xvgr/xmgr file
-life	hblife.xvg	Output, Opt.	xvgr/xmgr file
-nhbdist	nhbdist.xvg	Output, Opt.	xvgr/xmgr file

Other options

-h	bool	no	Print help info and quit
-nice	int	19	Set the nicelevel
-b	time	0	First frame (ps) to read from trajectory
-e	time	0	Last frame (ps) to read from trajectory
-dt	time	0	Only use frame when t MOD dt = first time (ps)
-xvgr	bool	yes	Add specific codes (legends etc.) in the output xvg files for the xmgrace program
-ins	bool	no	Analyze solvent insertion
-a	real	30	Cutoff angle (degrees, Acceptor - Donor - Hydrogen)
-r	real	0.35	Cutoff radius (nm, X - Acceptor, see next option)
-da	bool	yes	Use distance Donor-Acceptor (if TRUE) or Hydrogen-Acceptor (FALSE)
-abin	real	1	Binwidth angle distribution (degrees)
-rbin	real	0.005	Binwidth distance distribution (nm)
-nitacc	bool	yes	Regard nitrogen atoms as acceptors
-contact	bool	no	Do not look for hydrogen bonds, but merely for contacts within the cut-off distance
-shell	real	-1	when > 0, only calculate hydrogen bonds within # nm shell around one particle
-fitstart	real	1	Time from which to start fitting the correlation functions in order to obtain the forward and backward rate constants for HB breaking and formation
-temp	real	298.15	Temperature (K) for computing the Gibbs energy corresponding to HB breaking and reforming

-dump	int	0	Dump the first N hydrogen bond ACFs in a single xvg file for debugging
-max_hb	real	0	Theoretical maximum number of hydrogen bonds used for normalizing HB autocorrelation function. Can be useful in case the program estimates it wrongly
-merge	bool	yes	H-bonds between the same donor and acceptor, but with different hydrogen are treated as a single H-bond. Mainly important for the ACF.
-acflen	int	-1	Length of the ACF, default is half the number of frames
-normalize	bool	yes	Normalize ACF
-P	enum	0	Order of Legendre polynomial for ACF (0 indicates none): 0, 1, 2 or 3
-fitfn	enum	none	Fit function: none, exp, aexp, exp_exp, vac, exp5, exp7 or exp9
-ncskip	int	0	Skip N points in the output file of correlation functions
-beginfit	real	0	Time where to begin the exponential fit of the correlation function
-endfit	real	-1	Time where to end the exponential fit of the correlation function, -1 is till the end

- The option `-sel` that used to work on selected hbonds is out of order, and therefore not available for the time being.

D.31 *g_helix*

g_helix computes all kind of helix properties. First, the peptide is checked to find the longest helical part. This is determined by Hydrogen bonds and Phi/Psi angles. That bit is fitted to an ideal helix around the Z-axis and centered around the origin. Then the following properties are computed:

1. Helix radius (file *radius.xvg*). This is merely the RMS deviation in two dimensions for all Calpha atoms. It is calced as $\sqrt{(\text{SUM } i(x^2(i)+y^2(i)))/N}$, where N is the number of backbone atoms. For an ideal helix the radius is 0.23 nm
2. Twist (file *twist.xvg*). The average helical angle per residue is calculated. For alpha helix it is 100 degrees, for 3-10 helices it will be smaller, for 5-helices it will be larger.
3. Rise per residue (file *rise.xvg*). The helical rise per residue is plotted as the difference in Z-coordinate between Ca atoms. For an ideal helix this is 0.15 nm
4. Total helix length (file *len-ahx.xvg*). The total length of the helix in nm. This is simply the average rise (see above) times the number of helical residues (see below).
5. Number of helical residues (file *n-ahx.xvg*). The title says it all.
6. Helix Dipole, backbone only (file *dip-ahx.xvg*).
7. RMS deviation from ideal helix, calculated for the Calpha atoms only (file *rms-ahx.xvg*).
8. Average Calpha-Calpha dihedral angle (file *phi-ahx.xvg*).
9. Average Phi and Psi angles (file *phipsi.xvg*).
10. Ellipticity at 222 nm according to *Hirst and Brooks*

Files

-s	<i>topol.tpr</i>	Input	Generic run input: tpr tpb tpa xml
-n	<i>index.ndx</i>	Input	Index file
-f	<i>traj.xtc</i>	Input	Generic trajectory: xtc trr trj gro g96 pdb
-to	<i>gtraj.g87</i>	Output, Opt.	Gromos-87 ASCII trajectory format
-cz	<i>zconf.gro</i>	Output	Generic structure: gro g96 pdb xml
-co	<i>waver.gro</i>	Output	Generic structure: gro g96 pdb xml

Other options

-h	bool	no	Print help info and quit
-nice	int	19	Set the nicelevel
-b	time	0	First frame (ps) to read from trajectory

-e	time	0	Last frame (ps) to read from trajectory
-dt	time	0	Only use frame when t MOD dt = first time (ps)
-w	bool	no	View output xvg, xpm, eps and pdb files
-r0	int	1	The first residue number in the sequence
-q	bool	no	Check at every step which part of the sequence is helical
-F	bool	yes	Toggle fit to a perfect helix
-db	bool	no	Print debug info
-prop	enum	RAD	Select property to weight eigenvectors with. WARNING experimental stuff: RAD, TWIST, RISE, LEN, NHX, DIP, RMS, CPHI, RMSA, PHI, PSI, HB3, HB4, HB5 or CD222
-ev	bool	no	Write a new 'trajectory' file for ED
-ahxstart	int	0	First residue in helix
-ahxend	int	0	Last residue in helix

D.32 g_lie

g_lie computes a free energy estimate based on an energy analysis from. One needs an energy file with the following components: Coul (A-B) LJ-SR (A-B) etc.

Files

-f	ener.edr	Input	Generic energy: edr ene
-o	lie.xvg	Output	xvgr/xmgr file

Other options

-h	bool	no	Print help info and quit
-nice	int	19	Set the nicelevel
-b	time	0	First frame (ps) to read from trajectory
-e	time	0	Last frame (ps) to read from trajectory
-dt	time	0	Only use frame when t MOD dt = first time (ps)
-w	bool	no	View output xvg, xpm, eps and pdb files
-xvgr	bool	yes	Add specific codes (legends etc.) in the output xvg files for the xmgrace program
-Elj	real	0	Lennard-Jones interaction between ligand and solvent
-Eqq	real	0	Coulomb interaction between ligand and solvent
-Clj	real	0.181	Factor in the LIE equation for Lennard-Jones component of energy
-Cqq	real	0.5	Factor in the LIE equation for Coulomb component of energy
-ligand	string	none	Name of the ligand in the energy file

D.33 g_mdmat

g_mdmat makes distance matrices consisting of the smallest distance between residue pairs. With -frames these distance matrices can be stored as a function of time, to be able to see differences in tertiary structure as a function of time. If you choose your options unwise, this may generate a large output file. Default only an averaged matrix over the whole trajectory is output. Also a count of the number of different atomic contacts between residues over the whole trajectory can be made. The output can be processed with xpm2ps to make a PostScript (tm) plot.

Files

-f	traj.xtc	Input	Generic trajectory: xtc trr trj gro g96 pdb
-s	topol.tpr	Input	Structure+mass(db): tpr tpb tpa gro g96 pdb xml

-n	index.ndx	Input, Opt.	Index file
-mean	dm.xpm	Output	X PixMap compatible matrix file
-frames	dmf.xpm	Output, Opt.	X PixMap compatible matrix file
-no	num.xvg	Output, Opt.	xvgr/xmgr file

Other options

-h	bool	no	Print help info and quit
-nice	int	19	Set the nicelevel
-b	time	0	First frame (ps) to read from trajectory
-e	time	0	Last frame (ps) to read from trajectory
-dt	time	0	Only use frame when $t \text{ MOD } dt = \text{first time (ps)}$
-xvgr	bool	yes	Add specific codes (legends etc.) in the output xvg files for the xmgrace program
-t	real	1.5	trunc distance
-nlevels	int	40	Discretize distance in # levels

D.34 g_mindist

g_mindist computes the distance between one group and a number of other groups. Both the minimum distance (between any pair of atoms from the respective groups) and the number of contacts within a given distance are written to two separate output files. With `-or`, minimum distances to each residue in the first group are determined and plotted as a function of residue number.

With option `-pi` the minimum distance of a group to its periodic image is plotted. This is useful for checking if a protein has seen its periodic image during a simulation. Only one shift in each direction is considered, giving a total of 26 shifts. It also plots the maximum distance within the group and the lengths of the three box vectors.

Other programs that calculate distances are *g_dist* and *g_bond*.

Files

-f	traj.xtc	Input	Generic trajectory: xtc trr trj gro g96 pdb
-s	topol.tpr	Input, Opt.	Structure+mass(db): tpr tpb tpa gro g96 pdb xml
-n	index.ndx	Input, Opt.	Index file
-od	mindist.xvg	Output	xvgr/xmgr file
-on	numcont.xvg	Output, Opt.	xvgr/xmgr file
-o	atm-pair.out	Output, Opt.	Generic output file
-ox	mindist.xtc	Output, Opt.	Generic trajectory: xtc trr trj gro g96 pdb
-or	mindistres.xvg	Output, Opt.	xvgr/xmgr file

Other options

-h	bool	no	Print help info and quit
-nice	int	19	Set the nicelevel
-b	time	0	First frame (ps) to read from trajectory
-e	time	0	Last frame (ps) to read from trajectory
-dt	time	0	Only use frame when $t \text{ MOD } dt = \text{first time (ps)}$
-tu	enum	ps	Time unit: ps, fs, ns, us, ms or s
-w	bool	no	View output xvg, xpm, eps and pdb files
-xvgr	bool	yes	Add specific codes (legends etc.) in the output xvg files for the xmgrace program
-matrix	bool	no	Calculate half a matrix of group-group distances
-max	bool	no	Calculate *maximum* distance instead of minimum
-d	real	0.6	Distance for contacts

-pi	bool	no	Calculate minimum distance with periodic images
-split	bool	no	Split graph where time is zero
-ng	int	1	Number of secondary groups to compute distance to a central group

D.35 g_morph

`g_morph` does a linear interpolation of conformations in order to create intermediates. Of course these are completely unphysical, but that you may try to justify yourself. Output is in the form of a generic trajectory. The number of intermediates can be controlled with the `-ninterm` flag. The first and last flag correspond to the way of interpolating: 0 corresponds to input structure 1 while 1 corresponds to input structure 2. If you specify `first < 0` or `last > 1` extrapolation will be on the path from input structure `x1` to `x2`. In general the coordinates of the intermediate `x(i)` out of `N` total intermediates correspond to:

$$x(i) = x1 + (first + (i/(N-1)) * (last - first)) * (x2 - x1)$$

Finally the RMSD with respect to both input structures can be computed if explicitly selected (`-or` option). In that case an index file may be read to select what group RMS is computed from.

Files

-f1	conf1.gro	Input	Generic structure: gro g96 pdb tpr tpb tpa xml
-f2	conf2.gro	Input	Generic structure: gro g96 pdb tpr tpb tpa xml
-o	interm.xtc	Output	Generic trajectory: xtc trr trj gro g96 pdb
-or rms-	interm.xvg	Output, Opt.	xvgr/xmgr file
-n	index.ndx	Input, Opt.	Index file

Other options

-h	bool	no	Print help info and quit
-nice	int	0	Set the nicelevel
-w	bool	no	View output xvg, xpm, eps and pdb files
-xvgr	bool	yes	Add specific codes (legends etc.) in the output xvg files for the xmgrace program
-ninterm	int	11	Number of intermediates
-first	real	0	Corresponds to first generated structure (0 is input <code>x0</code> , see above)
-last	real	1	Corresponds to last generated structure (1 is input <code>x1</code> , see above)
-fit	bool	yes	Do a least squares fit of the second to the first structure before interpolating

D.36 g_msd

`g_msd` computes the mean square displacement (MSD) of atoms from their initial positions. This provides an easy way to compute the diffusion constant using the Einstein relation. The time between additional starting points for the MSD calculation is set with `-trestart`. The diffusion constant is calculated by least squares fitting a straight line through the MSD from `-beginfit` to `-endfit`. An error estimate given, which is the difference of the diffusion coefficients obtained from fits over the two halves of the fit interval.

Option `-mol` plots the MSD for molecules, this implies `-mw`, i.e. for each individual molecule a diffusion constant is computed. When using an index file, it should contain molecule numbers instead of atom numbers. Using this option one also gets an accurate error estimate based on the statistics between individual molecules. Since one usually is interested in self-diffusion at infinite dilution this is probably the most useful number.

Files

-f	traj.xtc	Input	Generic trajectory: xtc trr trj gro g96 pdb
-s	topol.tpr	Input	Structure+mass(db): tpr tpb tpa gro g96 pdb xml
-n	index.ndx	Input, Opt.	Index file
-o	msd.xvg	Output	xvgr/xmgr file
-mol	diff_mol.xvg	Output, Opt.	xvgr/xmgr file

Other options

-h	bool	no	Print help info and quit
-nice	int	19	Set the nicelevel
-b	time	0	First frame (ps) to read from trajectory
-e	time	0	Last frame (ps) to read from trajectory
-dt	time	0	Only use frame when t MOD dt = first time (ps)
-tu	enum	ps	Time unit: ps, fs, ns, us, ms or s
-w	bool	no	View output xvg, xpm, eps and pdb files
-xvgr	bool	yes	Add specific codes (legends etc.) in the output xvg files for the xmgrace program
-type	enum	no	Compute diffusion coefficient in one direction: no, x, y or z
-lateral	enum	no	Calculate the lateral diffusion in a plane perpendicular to: no, x, y or z
-ngroup	int	1	Number of groups to calculate MSD for
-mw	bool	yes	Mass weighted MSD
-trestart	time	10	Time between restarting points in trajectory (ps)
-beginfit	time	-1	Start time for fitting the MSD (ps), -1 is 10%
-endfit	time	-1	End time for fitting the MSD (ps), -1 is 90%

D.37 g_nmeig

g_nmeig calculates the eigenvectors/values of a (Hessian) matrix, which can be calculated with *mdrun*. The eigenvectors are written to a trajectory file (-v). The structure is written first with t=0. The eigenvectors are written as frames with the eigenvector number as timestamp. The eigenvectors can be analyzed with *g_anaeig*. An ensemble of structures can be generated from the eigenvectors with *g_nmens*. When mass weighting is used, the generated eigenvectors will be scaled back to plain cartesian coordinates before generating the output - in this case they will no longer be exactly orthogonal in the standard cartesian norm (But in the mass weighted norm they would be).

Files

-f	hessian.mtx	Input	Hessian matrix
-s	topol.tpr	Input	Structure+mass(db): tpr tpb tpa gro g96 pdb xml
-of	eigenfreq.xvg	Output	xvgr/xmgr file
-ol	eigenval.xvg	Output	xvgr/xmgr file
-v	eigenvec.trr	Output	Full precision trajectory: trr trj

Other options

-h	bool	no	Print help info and quit
-nice	int	19	Set the nicelevel
-xvgr	bool	yes	Add specific codes (legends etc.) in the output xvg files for the xmgrace program
-m	bool	yes	Divide elements of Hessian by product of sqrt(mass) of involved atoms prior to diagonalization. This should be used for 'Normal Modes' analysis
-first	int	1	First eigenvector to write away
-last	int	50	Last eigenvector to write away

D.38 g_nmens

`g_nmens` generates an ensemble around an average structure in a subspace which is defined by a set of normal modes (eigenvectors). The eigenvectors are assumed to be mass-weighted. The position along each eigenvector is randomly taken from a Gaussian distribution with variance $kT/\text{eigenvalue}$.

By default the starting eigenvector is set to 7, since the first six normal modes are the translational and rotational degrees of freedom.

Files

-v	eigenvec.trr	Input	Full precision trajectory: trr trj
-e	eigenval.xvg	Input	xvgr/xmgr file
-s	topol.tpr	Input	Structure+mass(db): tpr tpb tpa gro g96 pdb xml
-n	index.ndx	Input, Opt.	Index file
-o	ensemble.xtc	Output	Generic trajectory: xtc trr trj gro g96 pdb

Other options

-h	bool	no	Print help info and quit
-nice	int	19	Set the nicelevel
-xvgr	bool	yes	Add specific codes (legends etc.) in the output xvg files for the xmgrace program
-temp	real	300	Temperature in Kelvin
-seed	int	-1	Random seed, -1 generates a seed from time and pid
-num	int	100	Number of structures to generate
-first	int	7	First eigenvector to use (-1 is select)
-last	int	-1	Last eigenvector to use (-1 is till the last)

D.39 g_nmtraj

`g_nmtraj` generates a virtual trajectory from an eigenvector, corresponding to a harmonic cartesian oscillation around the average structure. The eigenvectors should normally be mass-weighted, but you can use non-weighted eigenvectors to generate orthogonal motions. The output frames are written as a trajectory file covering an entire period, and the first frame is the average structure. If you write the trajectory in (or convert to) PDB format you can view it directly in PyMol and also render a photorealistic movie. Motion amplitudes are calculated from the eigenvalues and a preset temperature, assuming equipartition of the energy over all modes. To make the motion clearly visible in PyMol you might want to amplify it by setting an unrealistic high temperature. However, be aware that both the linear cartesian displacements and mass weighting will lead to serious structure deformation for high amplitudes - this is simply a limitation of the cartesian normal mode model. By default the selected eigenvector is set to 7, since the first six normal modes are the translational and rotational degrees of freedom.

Files

-s	topol.tpr	Input	Structure+mass(db): tpr tpb tpa gro g96 pdb xml
-v	eigenvec.trr	Input	Full precision trajectory: trr trj
-o	nmtraj.xtc	Output	Generic trajectory: xtc trr trj gro g96 pdb

Other options

-h	bool	no	Print help info and quit
-nice	int	19	Set the nicelevel
-eignr	int	7	Eigenvector to use (first is 1)
-temp	real	300	Temperature in Kelvin
-amplitude	real	0.25	Amplitude for modes with eigenvalue ≤ 0
-nframes	int	30	Number of frames to generate

D.40 *g_order*

Compute the order parameter per atom for carbon tails. For atom *i* the vector *i-1*, *i+1* is used together with an axis. The index file has to contain a group with all equivalent atoms in all tails for each atom the order parameter has to be calculated for. The program can also give all diagonal elements of the order tensor and even calculate the deuterium order parameter *Scd* (default). If the option *-s*only is given, only one order tensor component (specified by the *-d* option) is given and the order parameter per slice is calculated as well. If *-s*only is not selected, all diagonal elements and the deuterium order parameter is given.

The tetrahedrality order parameters can be determined around an atom. Both angle and distance order parameters are calculated. See P.-L. Chau and A.J. Hardwick, *Mol. Phys.*, 93, (1998), 511-518. for more details.

Files

<i>-f</i>	<i>traj.xtc</i>	Input	Generic trajectory: <i>xtc trr trj gro g96 pdb</i>
<i>-n</i>	<i>index.ndx</i>	Input	Index file
<i>-s</i>	<i>topol.tpr</i>	Input	Generic run input: <i>tpr tpb tpa xml</i>
<i>-o</i>	<i>order.xvg</i>	Output	<i>xvgr/xmgr</i> file
<i>-od</i>	<i>deuter.xvg</i>	Output	<i>xvgr/xmgr</i> file
<i>-os</i>	<i>sliced.xvg</i>	Output	<i>xvgr/xmgr</i> file
<i>-Sg</i>	<i>sg-ang.xvg</i>	Output	<i>xvgr/xmgr</i> file
<i>-Sk</i>	<i>sk-dist.xvg</i>	Output	<i>xvgr/xmgr</i> file

Other options

<i>-h</i>	bool	no	Print help info and quit
<i>-nice</i>	int	19	Set the nicelevel
<i>-b</i>	time	0	First frame (ps) to read from trajectory
<i>-e</i>	time	0	Last frame (ps) to read from trajectory
<i>-dt</i>	time	0	Only use frame when <i>t MOD dt = first time (ps)</i>
<i>-w</i>	bool	no	View output <i>xvg</i> , <i>xpm</i> , <i>eps</i> and <i>pdb</i> files
<i>-xvgr</i>	bool	yes	Add specific codes (legends etc.) in the output <i>xvg</i> files for the <i>xmgrace</i> program
<i>-d</i>	enum	<i>z</i>	Direction of the normal on the membrane: <i>z</i> , <i>x</i> or <i>y</i>
<i>-sl</i>	int	1	Calculate order parameter as function of boxlength, dividing the box in <i>#nr</i> slices.
<i>-s</i> only	bool	no	Only give <i>Sz</i> element of order tensor. (axis can be specified with <i>-d</i>)
<i>-unsat</i>	bool	no	Calculate order parameters for unsaturated carbons. Note that this cannot be mixed with normal order parameters.

D.41 *g_potential*

Compute the electrostatical potential across the box. The potential is calculated by first summing the charges per slice and then integrating twice of this charge distribution. Periodic boundaries are not taken into account. Reference of potential is taken to be the left side of the box. It's also possible to calculate the potential in spherical coordinates as function of *r* by calculating a charge distribution in spherical slices and twice integrating them. *epsilon_r* is taken as 1,2 is more appropriate in many cases

Files

<i>-f</i>	<i>traj.xtc</i>	Input	Generic trajectory: <i>xtc trr trj gro g96 pdb</i>
<i>-n</i>	<i>index.ndx</i>	Input	Index file
<i>-s</i>	<i>topol.tpr</i>	Input	Generic run input: <i>tpr tpb tpa xml</i>

-o	potential.xvg	Output	xvgr/xmgr file
-oc	charge.xvg	Output	xvgr/xmgr file
-of	field.xvg	Output	xvgr/xmgr file

Other options

-h	bool	no	Print help info and quit
-nice	int	19	Set the nicelevel
-b	time	0	First frame (ps) to read from trajectory
-e	time	0	Last frame (ps) to read from trajectory
-dt	time	0	Only use frame when t MOD dt = first time (ps)
-w	bool	no	View output xvg, xpm, eps and pdb files
-xvgr	bool	yes	Add specific codes (legends etc.) in the output xvg files for the xmgrace program
-d	string	Z	Take the normal on the membrane in direction X, Y or Z.
-sl	int	10	Calculate potential as function of boxlength, dividing the box in #nr slices.
-cb	int	0	Discard first #nr slices of box for integration
-ce	int	0	Discard last #nr slices of box for integration
-tz	real	0	Translate all coordinates <distance> in the direction of the box
-spherical	bool	no	Calculate spherical thingie
-ng	int	1	Number of groups to consider

- Discarding slices for integration should not be necessary.

D.42 g_rama

`g_rama` selects the Phi/Psi dihedral combinations from your topology file and computes these as a function of time. Using simple Unix tools such as `grep` you can select out specific residues.

Files

-f	traj.xtc	Input	Generic trajectory: xtc trr trj gro g96 pdb
-s	topol.tpr	Input	Generic run input: tpr tpb tpa xml
-o	rama.xvg	Output	xvgr/xmgr file

Other options

-h	bool	no	Print help info and quit
-nice	int	19	Set the nicelevel
-b	time	0	First frame (ps) to read from trajectory
-e	time	0	Last frame (ps) to read from trajectory
-dt	time	0	Only use frame when t MOD dt = first time (ps)
-w	bool	no	View output xvg, xpm, eps and pdb files
-xvgr	bool	yes	Add specific codes (legends etc.) in the output xvg files for the xmgrace program

D.43 g_rdf

The structure of liquids can be studied by either neutron or X-ray scattering. The most common way to describe liquid structure is by a radial distribution function. However, this is not easy to obtain from a scattering experiment.

g_rdf calculates radial distribution functions in different ways. The normal method is around a (set of) particle(s), the other method is around the center of mass of a set of particles. With both methods *rdf*'s can also be calculated around axes parallel to the z-axis with option *-xy*.

If a run input file is supplied (*-s*), exclusions defined in that file are taken into account when calculating the *rdf*. The option *-cut* is meant as an alternative way to avoid intramolecular peaks in the *rdf* plot. It is however better to supply a run input file with a higher number of exclusions. For eg. benzene a topology with *nrexcl* set to 5 would eliminate all intramolecular contributions to the *rdf*. Note that all atoms in the selected groups are used, also the ones that don't have Lennard-Jones interactions.

Option *-cn* produces the cumulative number *rdf*.

To bridge the gap between theory and experiment structure factors can be computed (option *-sq*). The algorithm uses FFT, the gridspacing of which is determined by option *-grid*.

Files

<i>-f</i>	<i>traj.xtc</i>	Input	Generic trajectory: <i>xtc trr trj gro g96 pdb</i>
<i>-s</i>	<i>topol.tpr</i>	Input, Opt.	Structure+mass(db): <i>tpr tpb tpa gro g96 pdb xml</i>
<i>-n</i>	<i>index.ndx</i>	Input, Opt.	Index file
<i>-o</i>	<i>rdf.xvg</i>	Output, Opt.	<i>xvgr/xmgr</i> file
<i>-sq</i>	<i>sq.xvg</i>	Output, Opt.	<i>xvgr/xmgr</i> file
<i>-cn</i>	<i>rdf.cn.xvg</i>	Output, Opt.	<i>xvgr/xmgr</i> file
<i>-hq</i>	<i>hq.xvg</i>	Output, Opt.	<i>xvgr/xmgr</i> file

Other options

<i>-h</i>	bool	no	Print help info and quit
<i>-nice</i>	int	19	Set the nicelevel
<i>-b</i>	time	0	First frame (ps) to read from trajectory
<i>-e</i>	time	0	Last frame (ps) to read from trajectory
<i>-dt</i>	time	0	Only use frame when $t \text{ MOD } dt = \text{first time (ps)}$
<i>-w</i>	bool	no	View output <i>xvg</i> , <i>xpm</i> , <i>eps</i> and <i>pdb</i> files
<i>-xvgr</i>	bool	yes	Add specific codes (legends etc.) in the output <i>xvg</i> files for the <i>xmgrace</i> program
<i>-bin</i>	real	0.002	Binwidth (nm)
<i>-com</i>	bool	no	RDF with respect to the center of mass of first group
<i>-pbc</i>	bool	yes	Use periodic boundary conditions for computing distances
<i>-xy</i>	bool	no	Use only the x and y components of the distance
<i>-cut</i>	real	0	Shortest distance (nm) to be considered
<i>-ng</i>	int	1	Number of secondary groups to compute RDFs around a central group
<i>-fade</i>	real	0	From this distance onwards the RDF is transformed by $g'(r) = 1 + [g(r)-1] \exp(-(r/\text{fade}-1)^2)$ to make it go to 1 smoothly. If <i>fade</i> is 0.0 nothing is done.
<i>-nlevel</i>	int	20	Number of different colors in the diffraction image
<i>-startq</i>	real	0	Starting <i>q</i> (1/nm)
<i>-endq</i>	real	60	Ending <i>q</i> (1/nm)
<i>-energy</i>	real	12	Energy of the incoming X-ray (keV)

D.44 *g_rms*

g_rms compares two structures by computing the root mean square deviation (RMSD), the size-independent 'rho' similarity parameter (*rho*) or the scaled *rho* (*rhosc*), see Maiorov & Crippen, *PROTEINS* **22**, 273 (1995). This is selected by *-what*.

Each structure from a trajectory (*-f*) is compared to a reference structure. The reference structure is taken from the structure file (*-s*).

With option `-mir` also a comparison with the mirror image of the reference structure is calculated. This is useful as a reference for 'significant' values, see Maiorov & Crippen, *PROTEINS* **22**, 273 (1995).

Option `-prev` produces the comparison with a previous frame the specified number of frames ago.

Option `-m` produces a matrix in `.xpm` format of comparison values of each structure in the trajectory with respect to each other structure. This file can be visualized with for instance `xv` and can be converted to postscript with `xpm2ps`.

Option `-fit` controls the least-squares fitting of the structures on top of each other: complete fit (rotation and translation), translation only, or no fitting at all.

Option `-mw` controls whether mass weighting is done or not. If you select the option (default) and supply a valid `tpr` file masses will be taken from there, otherwise the masses will be deduced from the `atommass.dat` file in the GROMACS library directory. This is fine for proteins but not necessarily for other molecules. A default mass of 12.011 amu (Carbon) is assigned to unknown atoms. You can check whether this happens by turning on the `-debug` flag and inspecting the log file.

With `-f2`, the 'other structures' are taken from a second trajectory, this generates a comparison matrix of one trajectory versus the other.

Option `-bin` does a binary dump of the comparison matrix.

Option `-bm` produces a matrix of average bond angle deviations analogously to the `-m` option. Only bonds between atoms in the comparison group are considered.

Files

<code>-s</code>	<code>topol.tpr</code>	Input	Structure+mass(db): <code>tpr tpb tpa gro g96 pdb xml</code>
<code>-f</code>	<code>traj.xtc</code>	Input	Generic trajectory: <code>xtc trr trj gro g96 pdb</code>
<code>-f2</code>	<code>traj.xtc</code>	Input, Opt.	Generic trajectory: <code>xtc trr trj gro g96 pdb</code>
<code>-n</code>	<code>index.ndx</code>	Input, Opt.	Index file
<code>-o</code>	<code>rmsd.xvg</code>	Output	<code>xvgr/xmgr</code> file
<code>-mir</code>	<code>rmsdmir.xvg</code>	Output, Opt.	<code>xvgr/xmgr</code> file
<code>-a</code>	<code>avgrp.xvg</code>	Output, Opt.	<code>xvgr/xmgr</code> file
<code>-dist</code>	<code>rmsd-dist.xvg</code>	Output, Opt.	<code>xvgr/xmgr</code> file
<code>-m</code>	<code>rmsd.xpm</code>	Output, Opt.	X PixMap compatible matrix file
<code>-bin</code>	<code>rmsd.dat</code>	Output, Opt.	Generic data file
<code>-bm</code>	<code>bond.xpm</code>	Output, Opt.	X PixMap compatible matrix file

Other options

<code>-h</code>	bool	no	Print help info and quit
<code>-nice</code>	int	19	Set the nicelevel
<code>-b</code>	time	0	First frame (ps) to read from trajectory
<code>-e</code>	time	0	Last frame (ps) to read from trajectory
<code>-dt</code>	time	0	Only use frame when $t \text{ MOD } dt = \text{first time (ps)}$
<code>-tu</code>	enum	ps	Time unit: <code>ps, fs, ns, us, ms</code> or <code>s</code>
<code>-w</code>	bool	no	View output <code>xvg, xpm, eps</code> and <code>pdb</code> files
<code>-xvgr</code>	bool	yes	Add specific codes (legends etc.) in the output <code>xvg</code> files for the <code>xmgrace</code> program
<code>-what</code>	enum	rmsd	Structural difference measure: <code>rmsd, rho</code> or <code>rhosc</code>
<code>-pbc</code>	bool	yes	PBC check
<code>-fit</code>	enum	rot+trans	Fit to reference structure: <code>rot+trans, translation</code> or <code>none</code>
<code>-prev</code>	int	0	Compare with previous frame
<code>-split</code>	bool	no	Split graph where time is zero
<code>-skip</code>	int	1	Only write every <code>nr</code> -th frame to matrix
<code>-skip2</code>	int	1	Only write every <code>nr</code> -th frame to matrix

-max	real	-1	Maximum level in comparison matrix
-min	real	-1	Minimum level in comparison matrix
-bmax	real	-1	Maximum level in bond angle matrix
-bmin	real	-1	Minimum level in bond angle matrix
-mw	bool	yes	Use mass weighting for superposition
-nlevels	int	80	Number of levels in the matrices
-ng	int	1	Number of groups to compute RMS between

D.45 *g_rmsdist*

g_rmsdist computes the root mean square deviation of atom distances, which has the advantage that no fit is needed like in standard RMS deviation as computed by *g_rms*. The reference structure is taken from the structure file. The rmsd at time *t* is calculated as the rms of the differences in distance between atom-pairs in the reference structure and the structure at time *t*.

g_rmsdist can also produce matrices of the rms distances, rms distances scaled with the mean distance and the mean distances and matrices with NMR averaged distances ($1/r^3$ and $1/r^6$ averaging). Finally, lists of atom pairs with $1/r^3$ and $1/r^6$ averaged distance below the maximum distance (*-max*, which will default to 0.6 in this case) can be generated, by default averaging over equivalent hydrogens (all triplets of hydrogens named *[123]). Additionally a list of equivalent atoms can be supplied (*-equiv*), each line containing a set of equivalent atoms specified as residue number and name and atom name; e.g.:

```
3 SER HB1 3 SER HB2
```

Residue and atom names must exactly match those in the structure file, including case. Specifying non-sequential atoms is undefined.

Files

-f	traj.xtc	Input	Generic trajectory: xtc trr trj gro g96 pdb
-s	topol.tpr	Input	Structure+mass(db): tpr tpb tpa gro g96 pdb xml
-n	index.ndx	Input, Opt.	Index file
-equiv	equiv.dat	Input, Opt.	Generic data file
-o	distrmsd.xvg	Output	xvgr/xmgr file
-rms	rmsdist.xpm	Output, Opt.	X PixMap compatible matrix file
-scl	rmsscale.xpm	Output, Opt.	X PixMap compatible matrix file
-mean	rmsmean.xpm	Output, Opt.	X PixMap compatible matrix file
-nmr3	nmr3.xpm	Output, Opt.	X PixMap compatible matrix file
-nmr6	nmr6.xpm	Output, Opt.	X PixMap compatible matrix file
-noe	noe.dat	Output, Opt.	Generic data file

Other options

-h	bool	no	Print help info and quit
-nice	int	19	Set the nicelevel
-b	time	0	First frame (ps) to read from trajectory
-e	time	0	Last frame (ps) to read from trajectory
-dt	time	0	Only use frame when $t \text{ MOD } dt = \text{first time (ps)}$
-w	bool	no	View output xvg, xpm, eps and pdb files
-xvgr	bool	yes	Add specific codes (legends etc.) in the output xvg files for the xmgrace program
-nlevels	int	40	Discretize rms in # levels
-max	real	-1	Maximum level in matrices
-sumh	bool	yes	average distance over equivalent hydrogens

D.46 g_rmsf

`g_rmsf` computes the root mean square fluctuation (RMSF, i.e. standard deviation) of atomic positions after (optionally) fitting to a reference frame.

With option `-oq` the RMSF values are converted to B-factor values, which are written to a `pdb` file with the coordinates, of the structure file, or of a `pdb` file when `-q` is specified. Option `-ox` writes the B-factors to a file with the average coordinates.

With the option `-od` the root mean square deviation with respect to the reference structure is calculated.

With the option `aniso` `g_rmsf` will compute anisotropic temperature factors and then it will also output average coordinates and a `pdb` file with ANISOU records (corresponding to the `-oq` or `-ox` option). Please note that the U values are orientation dependent, so before comparison with experimental data you should verify that you fit to the experimental coordinates.

When a `pdb` input file is passed to the program and the `-aniso` flag is set a correlation plot of the U_{ij} will be created, if any anisotropic temperature factors are present in the `pdb` file.

With option `-dir` the average MSF (3x3) matrix is diagonalized. This shows the directions in which the atoms fluctuate the most and the least.

Files

<code>-f</code>	<code>traj.xtc</code>	Input	Generic trajectory: <code>xtc trr trj gro g96 pdb</code>
<code>-s</code>	<code>topol.tpr</code>	Input	Structure+mass(db): <code>tpr tpb tpa gro g96 pdb xml</code>
<code>-n</code>	<code>index.ndx</code>	Input, Opt.	Index file
<code>-q</code>	<code>eiwit.pdb</code>	Input, Opt.	Protein data bank file
<code>-oq</code>	<code>bfac.pdb</code>	Output, Opt.	Protein data bank file
<code>-ox</code>	<code>xaver.pdb</code>	Output, Opt.	Protein data bank file
<code>-o</code>	<code>rmsf.xvg</code>	Output	<code>xvgr/xmgr</code> file
<code>-od</code>	<code>rmsdev.xvg</code>	Output, Opt.	<code>xvgr/xmgr</code> file
<code>-oc</code>	<code>correl.xvg</code>	Output, Opt.	<code>xvgr/xmgr</code> file
<code>-dir</code>	<code>rmsf.log</code>	Output, Opt.	Log file

Other options

<code>-h</code>	<code>bool</code>	<code>no</code>	Print help info and quit
<code>-nice</code>	<code>int</code>	<code>19</code>	Set the nicelevel
<code>-b</code>	<code>time</code>	<code>0</code>	First frame (ps) to read from trajectory
<code>-e</code>	<code>time</code>	<code>0</code>	Last frame (ps) to read from trajectory
<code>-dt</code>	<code>time</code>	<code>0</code>	Only use frame when $t \text{ MOD } dt = \text{first time (ps)}$
<code>-w</code>	<code>bool</code>	<code>no</code>	View output <code>xvg</code> , <code>xpm</code> , <code>eps</code> and <code>pdb</code> files
<code>-xvgr</code>	<code>bool</code>	<code>yes</code>	Add specific codes (legends etc.) in the output <code>xvg</code> files for the <code>xmgrace</code> program
<code>-res</code>	<code>bool</code>	<code>no</code>	Calculate averages for each residue
<code>-aniso</code>	<code>bool</code>	<code>no</code>	Compute anisotropic temperature factors
<code>-fit</code>	<code>bool</code>	<code>yes</code>	Do a least squares superposition before computing RMSF. Without this you must make sure that the reference structure and the trajectory match.

D.47 g_rotacf

`g_rotacf` calculates the rotational correlation function for molecules. Three atoms (i,j,k) must be given in the index file, defining two vectors ij and jk . The rotational acf is calculated as the autocorrelation function of the vector $n = ij \times jk$, i.e. the cross product of the two vectors. Since three atoms span a plane, the order of the three atoms does not matter. Optionally, controlled by the `-d` switch, you can calculate the rotational correlation function for linear molecules by specifying two atoms (i,j) in the index file.

EXAMPLES

```
g_rotacf -P 1 -nparm 2 -fft -n index -o rotacf-x-P1 -fa expfit-x-P1 -beginfit 2.5 -endfit 20.0
```

This will calculate the rotational correlation function using a first order Legendre polynomial of the angle of a vector defined by the index file. The correlation function will be fitted from 2.5 ps till 20.0 ps to a two parameter exponential

Files

-f	traj.xtc	Input	Generic trajectory: xtc trr trj gro g96 pdb
-s	topol.tpr	Input	Generic run input: tpr tpb tpa xml
-n	index.ndx	Input	Index file
-o	rotacf.xvg	Output	xvgr/xmgr file

Other options

-h	bool	no	Print help info and quit
-nice	int	19	Set the nicelevel
-b	time	0	First frame (ps) to read from trajectory
-e	time	0	Last frame (ps) to read from trajectory
-dt	time	0	Only use frame when t MOD dt = first time (ps)
-w	bool	no	View output xvg, xpm, eps and pdb files
-xvgr	bool	yes	Add specific codes (legends etc.) in the output xvg files for the xmgrace program
-d	bool	no	Use index doublets (vectors) for correlation function instead of triplets (planes)
-aver	bool	yes	Average over molecules
-acflen	int	-1	Length of the ACF, default is half the number of frames
-normalize	bool	yes	Normalize ACF
-P	enum	0	Order of Legendre polynomial for ACF (0 indicates none): 0, 1, 2 or 3
-fitfn	enum	none	Fit function: none, exp, aexp, exp_exp, vac, exp5, exp7 or exp9
-ncskip	int	0	Skip N points in the output file of correlation functions
-beginfit	real	0	Time where to begin the exponential fit of the correlation function
-endfit	real	-1	Time where to end the exponential fit of the correlation function, -1 is till the end

D.48 g_saltbr

g_saltbr plots the difference between all combination of charged groups as a function of time. The groups are combined in different ways. A minimum distance can be given, (eg. the cut-off), then groups that are never closer than that distance will not be plotted.

Output will be in a number of fixed filenames, min-min.xvg, min-plus.xvg and plus-plus.xvg, or files for every individual ion-pair if selected

Files

-f	traj.xtc	Input	Generic trajectory: xtc trr trj gro g96 pdb
-s	topol.tpr	Input	Generic run input: tpr tpb tpa xml

Other options

-h	bool	no	Print help info and quit
-nice	int	19	Set the nicelevel
-b	time	0	First frame (ps) to read from trajectory
-e	time	0	Last frame (ps) to read from trajectory
-dt	time	0	Only use frame when t MOD dt = first time (ps)
-t	real	1000	trunc distance
-sep	bool	no	Use separate files for each interaction (may be MANY)

D.49 g_sas

`g_sas` computes hydrophobic, hydrophilic and total solvent accessible surface area. As a side effect the Connolly surface can be generated as well in a `pdb` file where the nodes are represented as atoms and the vertices connecting the nearest nodes as `CONNECT` records. The program will ask for a group for the surface calculation and a group for the output. The calculation group should always consist of all the non-solvent atoms in the system. The output group can be the whole or part of the calculation group. The area can be plotted per residue and atom as well (options `-or` and `-oa`). In combination with the latter option an `itp` file can be generated (option `-i`) which can be used to restrain surface atoms.

By default, periodic boundary conditions are taken into account, this can be turned off using the `-pbc` option.

Files

<code>-f</code>	<code>traj.xtc</code>	Input	Generic trajectory: <code>xtc trr trj gro g96 pdb</code>
<code>-s</code>	<code>topol.tpr</code>	Input	Generic run input: <code>tpr tpb tpa xml</code>
<code>-o</code>	<code>area.xvg</code>	Output	<code>xvgr/xmgr</code> file
<code>-or</code>	<code>resarea.xvg</code>	Output, Opt.	<code>xvgr/xmgr</code> file
<code>-oa</code>	<code>atomarea.xvg</code>	Output, Opt.	<code>xvgr/xmgr</code> file
<code>-q</code>	<code>connelly.pdb</code>	Output, Opt.	Protein data bank file
<code>-n</code>	<code>index.ndx</code>	Input, Opt.	Index file
<code>-i</code>	<code>surfat.itp</code>	Output, Opt.	Include file for topology

Other options

<code>-h</code>	bool	no	Print help info and quit
<code>-nice</code>	int	19	Set the nicelevel
<code>-b</code>	time	0	First frame (ps) to read from trajectory
<code>-e</code>	time	0	Last frame (ps) to read from trajectory
<code>-dt</code>	time	0	Only use frame when $t \text{ MOD } dt = \text{first time (ps)}$
<code>-w</code>	bool	no	View output <code>xvg</code> , <code>xpm</code> , <code>eps</code> and <code>pdb</code> files
<code>-xvgr</code>	bool	yes	Add specific codes (legends etc.) in the output <code>xvg</code> files for the <code>xmgrace</code> program
<code>-solsize</code>	real	0.14	Radius of the solvent probe (nm)
<code>-ndots</code>	int	24	Number of dots per sphere, more dots means more accuracy
<code>-qmax</code>	real	0.2	The maximum charge (e, absolute value) of a hydrophobic atom
<code>-f_index</code>	bool	no	Determine from a group in the index file what are the hydrophobic atoms rather than from the charge
<code>-minarea</code>	real	0.5	The minimum area (nm^2) to count an atom as a surface atom when writing a position restraint file (see help)
<code>-pbc</code>	bool	yes	Take periodicity into account
<code>-prot</code>	bool	yes	Output the protein to the <code>connelly.pdb</code> file too
<code>-dgs</code>	real	0	default value for solvation free energy per area (kJ/mol/nm^2)

D.50 g_sgangle

Compute the angle and distance between two groups. The groups are defined by a number of atoms given in an index file and may be two or three atoms in size. If `-one` is set, only one group should be specified in the index file and the angle between this group at time 0 and `t` will be computed. The angles calculated depend on the order in which the atoms are given. Giving for instance `5 6` will rotate the vector `5-6` with 180 degrees compared to giving `6 5`.

If three atoms are given, the normal on the plane spanned by those three atoms will be calculated, using the formula $\text{PIP2} \times \text{PIP3}$. The cos of the angle is calculated, using the inproduct of the two normalized vectors.

Here is what some of the file options do:

-oa: Angle between the two groups specified in the index file. If a group contains three atoms the normal to the plane defined by those three atoms will be used. If a group contains two atoms, the vector defined by those two atoms will be used.

-od: Distance between two groups. Distance is taken from the center of one group to the center of the other group.

-od1: If one plane and one vector is given, the distances for each of the atoms from the center of the plane is given seperately.

-od2: For two planes this option has no meaning.

Files

-f	traj.xtc	Input	Generic trajectory: xtc trr trj gro g96 pdb
-n	index.ndx	Input	Index file
-s	topol.tpr	Input	Generic run input: tpr tpb tpa xml
-oa	sg_angle.xvg	Output	xvgr/xmgr file
-od	sg_dist.xvg	Output, Opt.	xvgr/xmgr file
-od1	sg_dist1.xvg	Output, Opt.	xvgr/xmgr file
-od2	sg_dist2.xvg	Output, Opt.	xvgr/xmgr file

Other options

-h	bool	no	Print help info and quit
-nice	int	19	Set the nicelevel
-b	time	0	First frame (ps) to read from trajectory
-e	time	0	Last frame (ps) to read from trajectory
-dt	time	0	Only use frame when $t \text{ MOD } dt = \text{first time (ps)}$
-w	bool	no	View output xvg, xpm, eps and pdb files
-xvgr	bool	yes	Add specific codes (legends etc.) in the output xvg files for the xmgrace program
-pbc	bool	no	Use periodic boundary conditions
-one	bool	no	Only one group compute angle between vector at time zero and time t
-z	bool	no	Use the Z-axis as reference

D.51 g_sham

g_sham makes multi-dimensional free-energy, enthalpy and entropy plots. *g_sham* reads one or more xvg files and analyzes data sets. *g_sham* basic purpose is plotting Gibbs free energy landscapes by Boltzmann inverting multi-dimensional histograms, but it can also make enthalpy (option `-lsh`) and entropy (option `-lss`) plots. The histograms can be made for any quantities the user supplies. A line in the input file may start with a time (see option `-time`) and any number of y values may follow. Multiple sets can also be read when they are separated by & (option `-n`), in this case only one y value is read from each line. All lines starting with # and @ are skipped.

Option `-ge` can be used to supply a file with free energies when the ensemble is not a Boltzmann ensemble, but needs to be biased by this free energy. One free energy value is required for each (multi-dimensional) data point in the `-f` input.

Option `-ene` can be used to supply a file with energies. These energies are used as a weighting function in the single histogram analysis method due to Kumar et. al. When also temperatures are supplied (as a second column in the file) an experimental weighting scheme is applied. In addition the vales are used for making enthalpy and entropy plots.

With option `-dim` dimensions can be gives for distances. When a distance is 2- or 3-dimensional, the circumference or surface sampled by two particles increases with increasing distance. Depending on what

one would like to show, one can choose to correct the free-energy for this volume effect. The probability is normalized by r and r^2 for a dimension of 2 and 3 respectively. A value of -1 is used to indicate an angle in degrees between two vectors: a $\sin(\text{angle})$ normalization will be applied. Note that for angles between vectors the inner-product or cosine is the natural quantity to use, as it will produce bins of the same volume.

Files

-f	graph.xvg	Input	xvgr/xmgr file
-ge	gibbs.xvg	Input, Opt.	xvgr/xmgr file
-ene	esham.xvg	Input, Opt.	xvgr/xmgr file
-dist	ener.xvg	Output, Opt.	xvgr/xmgr file
-histo	edist.xvg	Output, Opt.	xvgr/xmgr file
-bin	bindex.ndx	Output, Opt.	Index file
-ls	gibbs.xpm	Output, Opt.	X PixMap compatible matrix file
-lsh	enthalpy.xpm	Output, Opt.	X PixMap compatible matrix file
-lss	entropy.xpm	Output, Opt.	X PixMap compatible matrix file
-map	map.xpm	Output, Opt.	X PixMap compatible matrix file
-ls3	gibbs3.pdb	Output, Opt.	Protein data bank file
-mdata	mapdata.xvg	Output, Opt.	xvgr/xmgr file
-g	shamlog.log	Output, Opt.	Log file

Other options

-h	bool	no	Print help info and quit
-nice	int	19	Set the nicelevel
-w	bool	no	View output xvg, xpm, eps and pdb files
-xvgr	bool	yes	Add specific codes (legends etc.) in the output xvg files for the xmgrace program
-time	bool	yes	Expect a time in the input
-b	real	-1	First time to read from set
-e	real	-1	Last time to read from set
-ttol	real	0	Tolerance on time in appropriate units (usually ps)
-n	int	1	Read # sets seperated by &
-d	bool	no	Use the derivative
-bw	real	0.1	Binwidth for the distribution
-sham	bool	yes	Turn off energy weighting even if energies are given
-tsham	real	298.15	Temperature for single histogram analysis
-pmin	real	0	Minimum probability. Anything lower than this will be set to zero
-dim	vector	1 1 1	Dimensions for distances, used for volume correction (max 3 values, dimensions > 3 will get the same value as the last)
-ngrid	vector	32 32	Number of bins for energy landscapes (max 3 values, dimensions > 3 will get the same value as the last)
-xmin	vector	0 0 0	Minimum for the axes in energy landscape (see above for > 3 dimensions)
-xmax	vector	1 1 1	Maximum for the axes in energy landscape (see above for > 3 dimensions)
-gmax	real	0	Maximum free energy in output, default is calculate
-emin	real	0	Minimum enthalpy in output, default is calculate
-emax	real	0	Maximum enthalpy in output, default is calculate
-nlevels	int	25	Number of levels for energy landscape
-mname	string		Legend label for the custom landscape

D.52 *g_sorient*

g_sorient analyzes solvent orientation around solutes. It calculates two angles between the vector from one or more reference positions to the first atom of each solvent molecule:

theta1: the angle with the vector from the first atom of the solvent molecule to the midpoint between atoms 2 and 3.

theta2: the angle with the normal of the solvent plane, defined by the same three atoms.

The reference can be a set of atoms or the center of mass of a set of atoms. The group of solvent atoms should consist of 3 atoms per solvent molecule. Only solvent molecules between *-rmin* and *-rmax* are considered for *-o* and *-no* each frame.

-o: distribution of $\cos(\theta_1)$ for $r_{\min} \leq r \leq r_{\max}$.

-no: distribution of $3\cos^2(\theta_2) - 1$ for $r_{\min} \leq r \leq r_{\max}$.

-ro: $\langle \cos(\theta_1) \rangle$ and $\langle 3\cos^2(\theta_2) - 1 \rangle$ as a function of the distance.

-co: the sum over all solvent molecules within distance *r* of $\cos(\theta_1)$ and $3\cos^2(\theta_2) - 1$ as a function of *r*.

Files

<i>-f</i>	<i>traj.xtc</i>	Input	Generic trajectory: <i>xtc trr trj gro g96 pdb</i>
<i>-s</i>	<i>topol.tpr</i>	Input	Structure+mass(db): <i>tpr tpb tpa gro g96 pdb xml</i>
<i>-n</i>	<i>index.ndx</i>	Input, Opt.	Index file
<i>-o</i>	<i>sori.xvg</i>	Output	xvgr/xmgr file
<i>-no</i>	<i>snor.xvg</i>	Output	xvgr/xmgr file
<i>-ro</i>	<i>sord.xvg</i>	Output	xvgr/xmgr file
<i>-co</i>	<i>scum.xvg</i>	Output	xvgr/xmgr file

Other options

<i>-h</i>	bool	no	Print help info and quit
<i>-nice</i>	int	19	Set the nicelevel
<i>-b</i>	time	0	First frame (ps) to read from trajectory
<i>-e</i>	time	0	Last frame (ps) to read from trajectory
<i>-dt</i>	time	0	Only use frame when $t \text{ MOD } dt = \text{first time (ps)}$
<i>-w</i>	bool	no	View output xvg, xpm, eps and pdb files
<i>-xvgr</i>	bool	yes	Add specific codes (legends etc.) in the output xvg files for the xmgrace program
<i>-com</i>	bool	no	Use the center of mass as the reference position
<i>-rmin</i>	real	0	Minimum distance
<i>-rmax</i>	real	0.5	Maximum distance
<i>-bin</i>	real	0.02	Binwidth
<i>-pbc</i>	bool	no	Check PBC for the center of mass calculation. Only necessary when your reference group consists of several molecules.

D.53 *g_tcaf*

g_tcaf computes tranverse current autocorrelations. These are used to estimate the shear viscosity η . For details see: Palmer, JCP 49 (1994) pp 359-366.

Transverse currents are calculated using the *k*-vectors (1,0,0) and (2,0,0) each also in the *y*- and *z*-direction, (1,1,0) and (1,-1,0) each also in the 2 other plains (these vectors are not independent) and (1,1,1) and the 3 other box diagonals (also not independent). For each *k*-vector the sine and cosine are used, in combination with the velocity in 2 perpendicular directions. This gives a total of $16 \cdot 2 \cdot 2 = 64$ transverse currents. One

autocorrelation is calculated fitted for each k-vector, which gives 16 tcáf's. Each of these tcáf's is fitted to $f(t) = \exp(-v)(\cosh(Wv) + 1/W \sinh(Wv))$, $v = -t/(2 \text{tau})$, $W = \sqrt{1 - 4 \text{tau} \text{eta}/\rho k^2}$, which gives 16 tau's and eta's. The fit weights decay with time as $\exp(-t/wt)$, the tcáf and fit are calculated up to time $5*wt$. The eta's should be fitted to $1 - a \text{eta}(k) k^2$, from which one can estimate the shear viscosity at $k=0$.

When the box is cubic, one can use the option `-oc`, which averages the tcáf's over all k-vectors with the same length. This results in more accurate tcáf's. Both the cubic tcáf's and fits are written to `-oc`. The cubic eta estimates are also written to `-ov`.

With option `-mol` the transverse current is determined of molecules instead of atoms. In this case the index group should consist of molecule numbers instead of atom numbers.

The k-dependent viscosities in the `-ov` file should be fitted to $\text{eta}(k) = \text{eta0} (1 - a k^2)$ to obtain the viscosity at infinite wavelength.

NOTE: make sure you write coordinates and velocities often enough. The initial, non-exponential, part of the autocorrelation function is very important for obtaining a good fit.

Files

<code>-f</code>	<code>traj.trr</code>	Input	Full precision trajectory: tr trj
<code>-s</code>	<code>topol.tpr</code>	Input, Opt.	Structure+mass(db): tpr tpb tpa gro g96 pdb xml
<code>-n</code>	<code>index.ndx</code>	Input, Opt.	Index file
<code>-ot</code>	<code>transcur.xvg</code>	Output, Opt.	xvgr/xmgr file
<code>-oa</code>	<code>tcáf_all.xvg</code>	Output	xvgr/xmgr file
<code>-o</code>	<code>tcáf.xvg</code>	Output	xvgr/xmgr file
<code>-of</code>	<code>tcáf_fit.xvg</code>	Output	xvgr/xmgr file
<code>-oc</code>	<code>tcáf_cub.xvg</code>	Output, Opt.	xvgr/xmgr file
<code>-ov</code>	<code>visc.k.xvg</code>	Output	xvgr/xmgr file

Other options

<code>-h</code>	<code>bool</code>	<code>no</code>	Print help info and quit
<code>-nice</code>	<code>int</code>	<code>19</code>	Set the nicelevel
<code>-b</code>	<code>time</code>	<code>0</code>	First frame (ps) to read from trajectory
<code>-e</code>	<code>time</code>	<code>0</code>	Last frame (ps) to read from trajectory
<code>-dt</code>	<code>time</code>	<code>0</code>	Only use frame when $t \text{ MOD } dt = \text{first time (ps)}$
<code>-w</code>	<code>bool</code>	<code>no</code>	View output xvg, xpm, eps and pdb files
<code>-xvgr</code>	<code>bool</code>	<code>yes</code>	Add specific codes (legends etc.) in the output xvg files for the xmgrace program
<code>-mol</code>	<code>bool</code>	<code>no</code>	Calculate tcáf of molecules
<code>-k34</code>	<code>bool</code>	<code>no</code>	Also use $k=(3,0,0)$ and $k=(4,0,0)$
<code>-wt</code>	<code>real</code>	<code>5</code>	Exponential decay time for the TCAF fit weights
<code>-acflen</code>	<code>int</code>	<code>-1</code>	Length of the ACF, default is half the number of frames
<code>-normalize</code>	<code>bool</code>	<code>yes</code>	Normalize ACF
<code>-P</code>	<code>enum</code>	<code>0</code>	Order of Legendre polynomial for ACF (0 indicates none): 0, 1, 2 or 3
<code>-fitfn</code>	<code>enum</code>	<code>none</code>	Fit function: none, exp, aexp, exp_exp, vac, exp5, exp7 or exp9
<code>-ncskip</code>	<code>int</code>	<code>0</code>	Skip N points in the output file of correlation functions
<code>-beginfit</code>	<code>real</code>	<code>0</code>	Time where to begin the exponential fit of the correlation function
<code>-endfit</code>	<code>real</code>	<code>-1</code>	Time where to end the exponential fit of the correlation function, -1 is till the end

D.54 g_traj

`g_traj` plots coordinates, velocities, forces and/or the box. With `-com` the coordinates, velocities and forces are calculated for the center of mass of each group. When `-mol` is set, the numbers in the index file are interpreted as molecule numbers and the same procedure as with `-com` is used for each molecule.

Option `-ot` plots the temperature of each group, provided velocities are present in the trajectory file. No corrections are made for constrained degrees of freedom! This implies `-com`.

Options `-ekt` and `-ekr` plot the translational and rotational kinetic energy of each group, provided velocities are present in the trajectory file. This implies `-com`.

Options `-cv` and `-cf` write the average velocities and average forces as temperature factors to a pdb file with the average coordinates. The temperature factors are scaled such that the maximum is 10. The scaling can be changed with the option `-scale`. To get the velocities or forces of one frame set both `-b` and `-e` to the time of desired frame. When averaging over frames you might need to use the `-nojump` option to obtain the correct average coordinates.

Option `-vd` computes a velocity distribution, i.e. the norm of the vector is plotted. In addition in the same graph the kinetic energy distribution is given.

Files

<code>-f</code>	<code>traj.xtc</code>	Input	Generic trajectory: xtc trr trj gro g96 pdb
<code>-s</code>	<code>topol.tpr</code>	Input	Structure+mass(db): tpr tpb tpa gro g96 pdb xml
<code>-n</code>	<code>index.ndx</code>	Input, Opt.	Index file
<code>-ox</code>	<code>coord.xvg</code>	Output, Opt.	xvgr/xmgr file
<code>-ov</code>	<code>veloc.xvg</code>	Output, Opt.	xvgr/xmgr file
<code>-of</code>	<code>force.xvg</code>	Output, Opt.	xvgr/xmgr file
<code>-ob</code>	<code>box.xvg</code>	Output, Opt.	xvgr/xmgr file
<code>-ot</code>	<code>temp.xvg</code>	Output, Opt.	xvgr/xmgr file
<code>-ekt</code>	<code>ektrans.xvg</code>	Output, Opt.	xvgr/xmgr file
<code>-ekr</code>	<code>ekrot.xvg</code>	Output, Opt.	xvgr/xmgr file
<code>-vd</code>	<code>veldist.xvg</code>	Output, Opt.	xvgr/xmgr file
<code>-cv</code>	<code>veloc.pdb</code>	Output, Opt.	Protein data bank file
<code>-cf</code>	<code>force.pdb</code>	Output, Opt.	Protein data bank file

Other options

<code>-h</code>	bool	no	Print help info and quit
<code>-nice</code>	int	19	Set the nicelevel
<code>-b</code>	time	0	First frame (ps) to read from trajectory
<code>-e</code>	time	0	Last frame (ps) to read from trajectory
<code>-dt</code>	time	0	Only use frame when $t \text{ MOD } dt = \text{first time (ps)}$
<code>-tu</code>	enum	ps	Time unit: ps, fs, ns, us, ms or s
<code>-w</code>	bool	no	View output xvg, xpm, eps and pdb files
<code>-xvgr</code>	bool	yes	Add specific codes (legends etc.) in the output xvg files for the xmgrace program
<code>-com</code>	bool	no	Plot data for the com of each group
<code>-mol</code>	bool	no	Index contains molecule numbers iso atom numbers
<code>-nojump</code>	bool	no	Remove jumps of atoms across the box
<code>-x</code>	bool	yes	Plot X-component
<code>-y</code>	bool	yes	Plot Y-component
<code>-z</code>	bool	yes	Plot Z-component
<code>-ng</code>	int	1	Number of groups to consider
<code>-len</code>	bool	no	Plot vector length
<code>-bin</code>	real	1	Binwidth for velocity histogram (nm/ps)
<code>-scale</code>	real	0	Scale factor for pdb output, 0 is autoscale

D.55 g_velacc

`g_velacc` computes the velocity autocorrelation function. When the `-s` option is used, the momentum autocorrelation function is calculated.

With option `-mol` the momentum autocorrelation function of molecules is calculated. In this case the index group should consist of molecule numbers instead of atom numbers.

Files

<code>-f</code>	<code>traj.trr</code>	Input	Full precision trajectory: trr trj
<code>-s</code>	<code>topol.tpr</code>	Input, Opt.	Structure+mass(db): tpr tpb tpa gro g96 pdb xml
<code>-n</code>	<code>index.ndx</code>	Input, Opt.	Index file
<code>-o</code>	<code>vac.xvg</code>	Output	xvgr/xmgr file

Other options

<code>-h</code>	bool	no	Print help info and quit
<code>-nice</code>	int	19	Set the nicelevel
<code>-b</code>	time	0	First frame (ps) to read from trajectory
<code>-e</code>	time	0	Last frame (ps) to read from trajectory
<code>-dt</code>	time	0	Only use frame when $t \text{ MOD } dt = \text{first time (ps)}$
<code>-w</code>	bool	no	View output xvg, xpm, eps and pdb files
<code>-xvgr</code>	bool	yes	Add specific codes (legends etc.) in the output xvg files for the xmgrace program
<code>-mol</code>	bool	no	Calculate vac of molecules
<code>-acflen</code>	int	-1	Length of the ACF, default is half the number of frames
<code>-normalize</code>	bool	yes	Normalize ACF
<code>-P</code>	enum	0	Order of Legendre polynomial for ACF (0 indicates none): 0, 1, 2 or 3
<code>-fitfn</code>	enum	none	Fit function: none, exp, aexp, exp_exp, vac, exp5, exp7 or exp9
<code>-ncskip</code>	int	0	Skip N points in the output file of correlation functions
<code>-beginfit</code>	real	0	Time where to begin the exponential fit of the correlation function
<code>-endfit</code>	real	-1	Time where to end the exponential fit of the correlation function, -1 is till the end

D.56 g_wham

This is an analysis program that implements the Weighted Histogram Analysis Method (WHAM). It is intended to analyze `.pdo` files generated by `mdrun` using umbrella sampling to create a potential of mean force (PMF). The options are

- `-o` name of the PMF output file
- `-hist` name of the histograms output file
- `-min` minimum coordinate to use
- `-max` maximum coordinate to use

Note: the program will throw out any data that is outside of `min - max`. The program will output the true `min` and `max` after completion, so you can use these values the next time. or you can use:

- `-noprof` only calculate `min` and `max`
- `-bins` number of bins to use in calculation

Files

<code>-o</code>	<code>profile.xvg</code>	Output	xvgr/xmgr file
<code>-hist</code>	<code>histo.xvg</code>	Output	xvgr/xmgr file

Other options

-h	bool	no	Print help info and quit
-nice	int	0	Set the nicelevel
-w	bool	no	View output xvg, xpm, eps and pdb files
-xvgr	bool	yes	Add specific codes (legends etc.) in the output xvg files for the xmgrace program
-min	real	0	Minimum coordinate in profile
-max	real	0	Maximum coordinate in profile
-bins	int	100	Number of bins in profile
-prof	bool	yes	Only calculate min and max
-temp	real	298	Temperature
-flip	bool	no	Combine halves of profile
-tol	real	0.01	Tolerance

D.57 genbox

Genbox can do one of 3 things:

1) Generate a box of solvent. Specify `-cs` and `-box`. Or specify `-cs` and `-cp` with a structure file with a box, but without atoms.

2) Solvate a solute configuration, eg. a protein, in a bath of solvent molecules. Specify `-cp` (solute) and `-cs` (solvent). The box specified in the solute coordinate file (`-cp`) is used, unless `-box` is set, which also centers the solute. The program `editconf` has more sophisticated options to change the box and center the solute. Solvent molecules are removed from the box where the distance between any atom of the solute molecule(s) and any atom of the solvent molecule is less than the sum of the VanderWaals radii of both atoms. A database (`vdwradii.dat`) of VanderWaals radii is read by the program, atoms not in the database are assigned a default distance `-vdw`.

3) Insert a number (`-nmol`) of extra molecules (`-ci`) at random positions. The program iterates until `nmol` molecules have been inserted in the box. To test whether an insertion is successful the same VanderWaals criterium is used as for removal of solvent molecules. When no appropriately sized holes (holes that can hold an extra molecule) are available the program tries for `-nmol * -try` times before giving up. Increase `-try` if you have several small holes to fill.

The default solvent is Simple Point Charge water (SPC), with coordinates from `$GMXLIB/spc216.gro`. Other solvents are also supported, as well as mixed solvents. The only restriction to solvent types is that a solvent molecule consists of exactly one residue. The residue information in the coordinate files is used, and should therefore be more or less consistent. In practice this means that two subsequent solvent molecules in the solvent coordinate file should have different residue number. The box of solute is built by stacking the coordinates read from the coordinate file. This means that these coordinates should be equilibrated in periodic boundary conditions to ensure a good alignment of molecules on the stacking interfaces.

The program can optionally rotate the solute molecule to align the longest molecule axis along a box edge. This way the amount of solvent molecules necessary is reduced. It should be kept in mind that this only works for short simulations, as eg. an alpha-helical peptide in solution can rotate over 90 degrees, within 500 ps. In general it is therefore better to make a more or less cubic box.

Setting `-shell` larger than zero will place a layer of water of the specified thickness (nm) around the solute. Hint: it is a good idea to put the protein in the center of a box first (using `editconf`).

Finally, `genbox` will optionally remove lines from your topology file in which a number of solvent molecules is already added, and adds a line with the total number of solvent molecules in your coordinate file.

Files

-cp	protein.gro	Input, Opt.	Generic structure: gro g96 pdb tpr tpb tpa xml
-cs	spc216.gro	Input, Opt., Lib	Generic structure: gro g96 pdb tpr tpb tpa xml
-ci	insert.gro	Input, Opt.	Generic structure: gro g96 pdb tpr tpb tpa xml
-o	out.gro	Output	Generic structure: gro g96 pdb xml
-p	topol.top	In/Out, Opt.	Topology file

Other options

-h	bool	no	Print help info and quit
-nice	int	19	Set the nicelevel
-box	vector	0 0 0	box size
-nmol	int	0	no of extra molecules to insert
-try	int	10	try inserting -nmol*-try times
-seed	int	1997	random generator seed
-vdwd	real	0.105	default vdwaals distance
-shell	real	0	thickness of optional water layer around solute
-maxsol	int	0	maximum number of solvent molecules to add if they fit in the box. If zero (default) this is ignored

- Molecules must be whole in the initial configurations.
- At the moment -ci only works when inserting one molecule.

D.58 genconf

genconf multiplies a given coordinate file by simply stacking them on top of each other, like a small child playing with wooden blocks. The program makes a grid of *user defined* proportions (-nbox), and interspaces the grid point with an extra space -dist.

When option -rot is used the program does not check for overlap between molecules on grid points. It is recommended to make the box in the input file at least as big as the coordinates + Van der Waals radius.

If the optional trajectory file is given, conformations are not generated, but read from this file and translated appropriately to build the grid.

Files

-f	conf.gro	Input	Generic structure: gro g96 pdb tpr tpb tpa xml
-o	out.gro	Output	Generic structure: gro g96 pdb xml
-trj	traj.xtc	Input, Opt.	Generic trajectory: xtc trr trj gro g96 pdb

Other options

-h	bool	no	Print help info and quit
-nice	int	0	Set the nicelevel
-nbox	vector	1 1 1	Number of boxes
-dist	vector	0 0 0	Distance between boxes
-seed	int	0	Random generator seed, if 0 generated from the time
-rot	bool	no	Randomly rotate conformations
-shuffle	bool	no	Random shuffling of molecules
-sort	bool	no	Sort molecules on X coord
-block	int	1	Divide the box in blocks on this number of cpus
-nmolat	int	3	Number of atoms per molecule, assumed to start from 0. If you set this wrong, it will screw up your system!
-maxrot	vector	0 90 90	Maximum random rotation
-renumber	bool	yes	Renumber residues

- The program should allow for random displacement off lattice points.

D.59 *genion*

genion replaces solvent molecules by monoatomic ions at the position of the first atoms with the most favorable electrostatic potential or at random. The potential is calculated on all atoms, using normal GROMACS particle based methods (in contrast to other methods based on solving the Poisson-Boltzmann equation). The potential is recalculated after every ion insertion. If specified in the run input file, a reaction field, shift function or user function can be used. For the user function a table file can be specified with the option `-table`. The group of solvent molecules should be continuous and all molecules should have the same number of atoms. The user should add the ion molecules to the topology file and include the file `ions.itp`. Ion names for Gromos96 should include the charge.

With the option `-pot` the potential can be written as B-factors in a `pdb` file (for visualisation using e.g. `rasmol`). The unit of the potential is 1000 kJ/(mol e), the scaling be changed with the `-scale` option.

For larger ions, e.g. sulfate we recommended to use `genbox`.

Files

<code>-s</code>	<code>topol.tpr</code>	Input	Generic run input: <code>tpr tpb tpa xml</code>
<code>-table</code>	<code>table.xvg</code>	Input, Opt.	<code>xvgr/xmgr</code> file
<code>-n</code>	<code>index.ndx</code>	Input, Opt.	Index file
<code>-o</code>	<code>out.gro</code>	Output	Generic structure: <code>gro g96 pdb xml</code>
<code>-g</code>	<code>genion.log</code>	Output	Log file
<code>-pot</code>	<code>pot.pdb</code>	Output, Opt.	Protein data bank file
<code>-p</code>	<code>topol.top</code>	In/Out, Opt.	Topology file

Other options

<code>-h</code>	bool	no	Print help info and quit
<code>-nice</code>	int	19	Set the nicelevel
<code>-xvgr</code>	bool	yes	Add specific codes (legends etc.) in the output <code>xvg</code> files for the <code>xmgrace</code> program
<code>-np</code>	int	0	Number of positive ions
<code>-pname</code>	string	Na	Name of the positive ion
<code>-pq</code>	int	1	Charge of the positive ion
<code>-nn</code>	int	0	Number of negative ions
<code>-nname</code>	string	Cl	Name of the negative ion
<code>-nq</code>	int	-1	Charge of the negative ion
<code>-rmin</code>	real	0.6	Minimum distance between ions
<code>-random</code>	bool	yes	Use random placement of ions instead of based on potential. The <code>rmin</code> option should still work
<code>-seed</code>	int	1993	Seed for random number generator
<code>-scale</code>	real	0.001	Scaling factor for the potential for <code>-pot</code>
<code>-conc</code>	real	0	Specify salt concentration (mol/liter). This will add sufficient ions to reach up to the specified concentration as computed from the volume of the cell in the input <code>tpr</code> file. Overrides the <code>-np</code> and <code>nn</code> options.
<code>-neutral</code>	bool	no	This option will add enough ions to neutralize the system. In combination with the concentration option a neutral system at a given salt concentration will be generated.

- Calculation of the potential is not reliable, therefore the `-random` option is now turned on by default.
- If you specify a salt concentration existing ions are not taken into account. In effect you therefore specify the amount of salt to be added.

D.60 genpr

genpr produces an include file for a topology containing a list of atom numbers and three force constants for the X, Y and Z direction. A single isotropic force constant may be given on the command line instead of three components.

WARNING: genpr only works for the first molecule. Position restraints are interactions within molecules, therefore they should be included within the correct [`moleculetype`] block in the topology. Since the atom numbers in every moleculetype in the topology start at 1 and the numbers in the input file for genpr number consecutively from 1, genpr will only produce a useful file for the first molecule.

The `-of` option produces an index file that can be used for freezing atoms. In this case the input file must be a `pdb` file.

Files

<code>-f</code>	<code>conf.gro</code>	Input	Generic structure: <code>gro g96 pdb tpr tpb tpa xml</code>
<code>-n</code>	<code>index.ndx</code>	Input, Opt.	Index file
<code>-o</code>	<code>posre.itp</code>	Output	Include file for topology
<code>-of</code>	<code>freeze.ndx</code>	Output, Opt.	Index file

Other options

<code>-h</code>	<code>bool</code>	<code>no</code>	Print help info and quit
<code>-nice</code>	<code>int</code>	<code>0</code>	Set the nicelevel
<code>-fc</code>	<code>vector</code>		
	<code>1000 1000 1000</code>		force constants (kJ mol ⁻¹ nm ⁻²)
<code>-freeze</code>	<code>real</code>	<code>0</code>	if the <code>-of</code> option or this one is given an index file will be written containing atom numbers of all atoms that have a B-factor less than the level given here

D.61 gmxcheck

gmxcheck reads a trajectory (`.trj`, `.trr` or `.xtc`) or an energy file (`.ene` or `.edr`) and prints out useful information about them.

Option `-c` checks for presence of coordinates, velocities and box in the file, for close contacts (smaller than `-vdwfac` and not bonded, i.e. not between `-bonlo` and `-bonhi`, all relative to the sum of both Van der Waals radii) and atoms outside the box (these may occur often and are no problem). If velocities are present, an estimated temperature will be calculated from them.

If an index file is given it's contents will be summarized.

If both a trajectory and a `tpr` file are given (with `-s1`) the program will check whether the bond lengths defined in the `tpr` file are indeed correct in the trajectory. If not you may have non-matching files due to e.g. deshuffling or due to problems with virtual sites. With these flags, gmxcheck provides a quick check for such problems.

The program can compare run two input (`.tpr`, `.tpb` or `.tpa`) files when both `-s1` and `-s2` are supplied. Similarly a pair of trajectory files can be compared (using the `-f2` option), or a pair of energy files (using the `-e2` option).

For free energy simulations the A and B state topology from one run input file can be compared with options `-s1` and `-ab`.

Files

<code>-f</code>	<code>traj.xtc</code>	Input, Opt.	Generic trajectory: <code>xtc trr trj gro g96 pdb</code>
<code>-f2</code>	<code>traj.xtc</code>	Input, Opt.	Generic trajectory: <code>xtc trr trj gro g96 pdb</code>

-s1	top1.tpr	Input, Opt.	Generic run input: tpr tpb tpa xml
-s2	top2.tpr	Input, Opt.	Generic run input: tpr tpb tpa xml
-c	topol.tpr	Input, Opt.	Structure+mass(db): tpr tpb tpa gro g96 pdb xml
-e	ener.edr	Input, Opt.	Generic energy: edr ene
-e2	ener2.edr	Input, Opt.	Generic energy: edr ene
-n	index.ndx	Input, Opt.	Index file

Other options

-h	bool	no	Print help info and quit
-nice	int	0	Set the nicelevel
-vdwfac	real	0.8	Fraction of sum of VdW radii used as warning cutoff
-bonlo	real	0.4	Min. fract. of sum of VdW radii for bonded atoms
-bonhi	real	0.7	Max. fract. of sum of VdW radii for bonded atoms
-tol	real	0.001	Relative tolerance for comparing real values defined as $2*(a-b)/(oraor+orbor)$
-ab	bool	no	Compare the A and B topology from one file
-lastener	string		Last energy term to compare (if not given all are tested). It makes sense to go up until the Pressure.

D.62 gmxdump

gmxdump reads a run input file (.tpa/.tpr/.tpb), a trajectory (.trj/.trr/.xtc) or an energy file (.ene/.edr) and prints that to standard output in a readable format. This program is essential for checking your run input file in case of problems.

Files

-s	topol.tpr	Input, Opt.	Generic run input: tpr tpb tpa xml
-f	traj.xtc	Input, Opt.	Generic trajectory: xtc trr trj gro g96 pdb
-e	ener.edr	Input, Opt.	Generic energy: edr ene

Other options

-h	bool	no	Print help info and quit
-nice	int	0	Set the nicelevel
-nr	bool	yes	Show index numbers in output (leaving them out makes comparison easier, but creates a useless topology)

D.63 grompp

The gromacs preprocessor reads a molecular topology file, checks the validity of the file, expands the topology from a molecular description to an atomic description. The topology file contains information about molecule types and the number of molecules, the preprocessor copies each molecule as needed. There is no limitation on the number of molecule types. Bonds and bond-angles can be converted into constraints, separately for hydrogens and heavy atoms. Then a coordinate file is read and velocities can be generated from a Maxwellian distribution if requested. grompp also reads parameters for the mdrun (eg. number of MD steps, time step, cut-off), and others such as NEMD parameters, which are corrected so that the net acceleration is zero. Eventually a binary file is produced that can serve as the sole input file for the MD program.

grompp uses the atom names from the topology file. The atom names in the coordinate file (option -c) are only read to generate warnings when they do not match the atom names in the topology. Note that

the atom names are irrelevant for the simulation as only the atom types are used for generating interaction parameters.

grompp calls a preprocessor to resolve includes, macros etcetera. By default we use the `cpp` in your path. To specify a different macro-preprocessor (e.g. `m4`) or alternative location you can put a line in your parameter file specifying the path to that program. Specifying `-pp` will get the pre-processed topology file written out.

If your system does not have a `c`-preprocessor, you can still use `grompp`, but you do not have access to the features from the `cpp`. Command line options to the `c`-preprocessor can be given in the `.mdp` file. See your local manual (`man cpp`).

When using position restraints a file with restraint coordinates can be supplied with `-r`, otherwise restraining will be done with respect to the conformation from the `-c` option. For free energy calculation the the coordinates for the B topology can be supplied with `-rb`, otherwise they will be equal to those of the A topology.

Starting coordinates can be read from trajectory with `-t`. The last frame with coordinates and velocities will be read, unless the `-time` option is used. Note that these velocities will not be used when `gen_vel = yes` in your `.mdp` file. An energy file can be supplied with `-e` to have exact restarts when using pressure and/or Nose-Hoover temperature coupling. For an exact restart do not forget to turn off velocity generation and turn on unconstrained starting when constraints are present in the system. If you want to continue a crashed run, it is easier to use `tpbconv`.

When preparing an input file for parallel `mdrun` it may be advantageous to partition the simulation system over the nodes in a way in which each node has a similar amount of work. The `-shuffle` option does just that. For a single protein in water this does not make a difference, however for a system where you have many copies of different molecules (e.g. liquid mixture or membrane/water system) the option is definitely a must. The output trajectories will also be shuffled. `grompp` writes an index file (option `-deshuf`) which can be used with `trjconv` to deshuffle the trajectories.

A further optimization for parallel systems is the `-sort` option which sorts molecules according to coordinates. This must always be used in conjunction with `-shuffle`, however sorting also works when you have only one molecule type.

Using the `-morse` option `grompp` can convert the harmonic bonds in your topology to morse potentials. This makes it possible to break bonds. For this option to work you need an extra file in your `$GMXLIB` with dissociation energy. Use the `-debug` option to get more information on the workings of this option (look for `MORSE` in the `grompp.log` file using `less` or something like that).

By default all bonded interactions which have constant energy due to virtual site constructions will be removed. If this constant energy is not zero, this will result in a shift in the total energy. All bonded interactions can be kept by turning off `-rmvsbds`. Additionally, all constraints for distances which will be constant anyway because of virtual site constructions will be removed. If any constraints remain which involve virtual sites, a fatal error will result.

To verify your run input file, please make notice of all warnings on the screen, and correct where necessary. Do also look at the contents of the `mdout.mdp` file, this contains comment lines, as well as the input that `grompp` has read. If in doubt you can start `grompp` with the `-debug` option which will give you more information in a file called `grompp.log` (along with real debug info). Finally, you can see the contents of the run input file with the `gmxdump` program.

Files

<code>-f</code>	<code>grompp.mdp</code>	Input, Opt.	<code>grompp</code> input file with MD parameters
<code>-po</code>	<code>mdout.mdp</code>	Output	<code>grompp</code> input file with MD parameters
<code>-c</code>	<code>conf.gro</code>	Input	Generic structure: <code>gro g96 pdb tpr tpb tpa xml</code>
<code>-r</code>	<code>conf.gro</code>	Input, Opt.	Generic structure: <code>gro g96 pdb tpr tpb tpa xml</code>
<code>-rb</code>	<code>conf.gro</code>	Input, Opt.	Generic structure: <code>gro g96 pdb tpr tpb tpa xml</code>

-n	index.ndx	Input, Opt.	Index file
-deshuf	deshuf.ndx	Output, Opt.	Index file
-p	topol.top	Input	Topology file
-pp	processed.top	Output, Opt.	Topology file
-o	topol.tpr	Output	Generic run input: tpr tpb tpa xml
-t	traj.trr	Input, Opt.	Full precision trajectory: trr trj
-e	ener.edr	Input, Opt.	Generic energy: edr ene

Other options

-h	bool	no	Print help info and quit
-nice	int	0	Set the nicelevel
-v	bool	yes	Be loud and noisy
-time	real	-1	Take frame at or first after this time.
-np	int	1	Generate statusfile for # nodes
-shuffle	bool	no	Shuffle molecules over nodes
-sort	bool	no	Sort molecules according to X coordinate
-rmvsbds	bool	yes	Remove constant bonded interactions with virtual sites
-load	string		Relative load capacity of each node on a parallel machine. Be sure to use quotes around the string, which should contain a number for each node
-maxwarn	int	10	Number of warnings after which input processing stops
-check14	bool	no	Remove 1-4 interactions without Van der Waals
-renum	bool	yes	Renumber atomtypes and minimize number of atomtypes

D.64 highway

highway is the gromacs highway simulator. It is an X-windows gadget that shows a (periodic) Autobahn with a user defined number of cars. Fog can be turned on or off to increase the number of crashes. Nice for a background CPU-eater. A sample input file is in \$GMXDATA/top/highway.dat

Files

-f	highway.dat	Input	Generic data file
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Other options

-h	bool	no	Print help info and quit
-nice	int	0	Set the nicelevel

D.65 make_edi

make_edi generates an ED-sampling input file to be used with mdrun based on eigenvectors of a covariance matrix (`g_covar`) or from a Normal Modes analysis (`g_nmeig`). ED-sampling can be used to manipulate the position along collective coordinates (eigenvectors) of (biological) macromolecules during a simulation. Particularly, it may be used to enhance the sampling efficiency of MD simulations by stimulating the system to explore new regions along these collective coordinates. A number of different algorithms are implemented to drive the system along the eigenvectors (`-linfix`, `-linacc`, `-radfix`, `-radacc`, `-radcon`), to keep the position along a certain (set of) coordinate(s) fixed (`-linfix`), or to only monitor the projections of the positions, velocities and forces onto these coordinates (`-mon`).

References:

A. Amadei, A.B.M. Linssen, B.L. de Groot, D.M.F. van Aalten and H.J.C. Berendsen; An efficient method

for sampling the essential subspace of proteins., J. Biomol. Struct. Dyn. 13:615-626 (1996)

B.L. de Groot, A. Amadei, D.M.F. van Aalten and H.J.C. Berendsen; Towards an exhaustive sampling of the configurational spaces of the two forms of the peptide hormone guanylin, J. Biomol. Struct. Dyn. 13 : 741-751 (1996)

B.L. de Groot, A. Amadei, R.M. Scheek, N.A.J. van Nuland and H.J.C. Berendsen; An extended sampling of the configurational space of HPr from E. coli PROTEINS: Struct. Funct. Gen. 26: 314-322 (1996)

You will be prompted for one or more index groups that correspond to the eigenvectors, reference structure, target positions, etc.

-mon: monitor projections of x, v and f onto selected eigenvectors.

-linfix: perform fixed-step linear expansion along selected eigenvectors.

-linacc: perform acceptance linear expansion along selected eigenvectors. (steps in the desired directions will be accepted, others will be rejected).

-radfix: perform fixed-step radius expansion along selected eigenvectors.

-radacc: perform acceptance radius expansion along selected eigenvectors. (steps in the desired direction will be accepted, others will be rejected). Note: by default the starting MD structure will be taken as origin of the first expansion cycle for radius expansion. If -ori is specified, you will be able to read in a structure file that defines an external origin.

-radcon: perform acceptance radius contraction along selected eigenvectors towards a target structure specified with -tar. NOTE: each eigenvector can be selected only once.

-outfrq: frequency (in steps) of writing out projections etc.

-logfrq: frequency (in steps) of writing out statistics to log file.

-slope: minimal slope in acceptance radius expansion. A new expansion cycle will be started if the spontaneous increase of the radius (in nm/step) is less than the value specified.

-maxedsteps: maximum number of steps per cycle in radius expansion before a new cycle is started.

Note on the parallel implementation: since ED sampling is a 'global' thing (collective coordinates etc), at least on the 'protein' side, ED sampling is not very parallel-friendly from an implementation point of view (it would require much extra communication to fully parallelize the algorithms). Fortunately, however, a typical parallel protein simulation in gromacs has most or all protein coordinates on one processor (the master) and has only other atoms (solvent, lipid, ions etc) on the other processors. With such a setup, ED sampling will still work. If the atoms over which ED sampling should be performed are spread over multiple processors, a fatal error will result.

All output of mdrun (specify with -eo) is written to a .edo file (some extra information is written to the log file of mdrun too, actually). The .edo format is a simple ASCII file that should be easy to parse with standard unix tools like awk. A script (parse_edo) can be downloaded from contribution section at www.gromacs.org to extract information from the .edo files for your convenience. In short, the header defines which information can be expected in the rest of the .edo file. After the header, per step the following information is present:

* the step number

RMSD (for atoms in fitting prior to calculating ED constr.)

projections of the positions onto selected eigenvectors

projections of the velocities onto selected eigenvectors

projections of the forces onto selected eigenvectors

All projections are in the same order as in the header, so if you have e.g. 2 groups (say one group over which acceptance radius expansion is performed, and another for which the projections are merely monitored) then you first get the position projections for each of the 2 groups, then the velocities and then the forces. Radii are not explicitly written to the .edo file, as they can be readily projected back from the positions.

Alternatively, radii may be 'grep'ped from the log file.

FLOODING:

with `-flood` you can specify which eigenvectors are used to compute a flooding potential, which will lead to extra forces expelling the structure out of the region described by the covariance matrix. if you switch `-restrain` the potential is inverted and the structure is kept in that region

the origin is normally the average structure stored in the `eigvec.trr` file it can be changed with `-ori` to an arbitrary position in configurational space with `-tau`, `-deltaF0` and `-Eflnull` you control the flooding strength `Efl` is the flooding strength, it is updated according to the rule of adaptive flooding `tau` is the time constant of adaptive flooding, high `tau` means slow adaption (i.e. growth) `deltaF0` is the flooding strength you want to reach after `tau` ps of simulation to use constant `Efl` set `-tau` to zero

`-alpha` is a fudge parameter to control the width of the flooding potential. A value of 2 has been found to give good results for most standard cases in flooding of proteins `alpha` basically accounts for incomplete sampling, if you sampled further the width of the ensemble would increase, this is mimicked by `alpha > 1` for restraining `alpha < 1` can give you smaller width in the restraining potential **RESTART and FLOODING:** If you want to restart a crashed flooding simulation please find the values `deltaF` and `Efl` in the output file and write them with your texteditor into the `.edi` file under `DELTA_F0` and `EFL_NULL`

Files

<code>-f</code>	<code>eigvec.trr</code>	Input	Full precision trajectory: <code>trr trj</code>
<code>-eig</code>	<code>eigenval.xvg</code>	Input, Opt.	<code>xvgr/xmgr</code> file
<code>-s</code>	<code>topol.tpr</code>	Input	Structure+mass(db): <code>tpr tpb tpa gro g96 pdb xml</code>
<code>-n</code>	<code>index.ndx</code>	Input, Opt.	Index file
<code>-tar</code>	<code>target.gro</code>	Input, Opt.	Generic structure: <code>gro g96 pdb tpr tpb tpa xml</code>
<code>-ori</code>	<code>origin.gro</code>	Input, Opt.	Generic structure: <code>gro g96 pdb tpr tpb tpa xml</code>
<code>-o</code>	<code>sam.edi</code>	Output	ED sampling input

Other options

<code>-h</code>	bool	no	Print help info and quit
<code>-nice</code>	int	0	Set the nicelevel
<code>-xvgr</code>	bool	yes	Add specific codes (legends etc.) in the output <code>xvg</code> files for the <code>xmgrace</code> program
<code>-mon</code>	string		Indices of eigenvectors for projections of <code>x</code> , <code>v</code> and <code>f</code> (e.g. 1,2-5,9) or 1-100:10 means 1 11 21 31 ... 91
<code>-linfix</code>	string		Indices of eigenvectors for fixed increment linear sampling
<code>-linacc</code>	string		Indices of eigenvectors for acceptance linear sampling
<code>-radfix</code>	string		Indices of eigenvectors for fixed increment radius expansion
<code>-radacc</code>	string		Indices of eigenvectors for acceptance radius expansion
<code>-radcon</code>	string		Indices of eigenvectors for acceptance radius contraction
<code>-flood</code>	string		Indices of eigenvectors for flooding
<code>-outfrq</code>	int	100	frequency (in steps) of writing output in <code>.edo</code> file
<code>-logfrq</code>	int	100	frequency (in steps) of writing to log
<code>-slope</code>	real	0	minimal slope in acceptance radius expansion
<code>-maxedsteps</code>	int	0	max nr of steps per cycle
<code>-deltaF0</code>	real	150	target destabilization energy - used for flooding
<code>-deltaF</code>	real	0	start <code>deltaF</code> with this parameter - default 0, i.g. nonzero values only needed for restart
<code>-tau</code>	real	0.1	coupling constant for adaption of flooding strength according to <code>deltaF0</code> , 0 = infinity i.e. constant flooding strength
<code>-eqsteps</code>	int	0	number of steps to run without any perturbations
<code>-Eflnull</code>	real	0	this is the starting value of the flooding strength. The flooding strength is updated according to the adaptive flooding scheme. To use a constant flooding strength use <code>-tau 0</code> .
<code>-T</code>	real	300	<code>T</code> is temperature, the value is needed if you want to do flooding

-alpha	real	1	scale width of gaussian flooding potential with α^2
-linstep	string		Stepsize (nm/step) for fixed increment linear sampling (put in quotes! "1.0 2.3 5.1 -3.1")
-accdir	string		Directions for acceptance linear sampling - only sign counts! (put in quotes! "-1 +1 -1.1")
-radstep	real	0	Stepsize (nm/step) for fixed increment radius expansion
-restrain	bool	no	use the flooding potential with inverted sign -> effects as quasiharmonic restraining potential
-hesse	bool	no	the eigenvectors and eigenvalues are from a Hesse matrix
-harmonic	bool	no	the eigenvalues are interpreted as spring constant

D.66 make_ndx

Index groups are necessary for almost every gromacs program. All these programs can generate default index groups. You ONLY have to use make_ndx when you need SPECIAL index groups. There is a default index group for the whole system, 9 default index groups are generated for proteins, a default index group is generated for every other residue name.

When no index file is supplied, also make_ndx will generate the default groups. With the index editor you can select on atom, residue and chain names and numbers. When a run input file is supplied you can also select on atom type. You can use NOT, AND and OR, you can split groups into chains, residues or atoms. You can delete and rename groups.

The atom numbering in the editor and the index file starts at 1.

Files

-f	conf.gro	Input, Opt.	Generic structure: gro g96 pdb tpr tpb tpa xml
-n	index.ndx	Input, Opt., Index	Index file
-o	index.ndx	Output	Index file

Other options

-h	bool	no	Print help info and quit
-nice	int	0	Set the nicelevel
-natoms	int	0	set number of atoms (default: read from coordinate or index file)

D.67 mdrun

The mdrun program is the main computational chemistry engine within GROMACS. Obviously, it performs Molecular Dynamics simulations, but it can also perform Brownian Dynamics and Langevin Dynamics as well as Conjugate Gradient or Steepest Descents energy minimization. Normal mode analysis is another option. In this case mdrun builds a Hessian matrix from single conformation. For usual Normal Modes-like calculations, make sure that the structure provided is properly energy-minimised. The generated matrix can be diagonalized by g_nmeig.

The mdrun program reads the run input file (-s) and distributes the topology over nodes if needed. The coordinates are passed around, so that computations can begin. First a neighborlist is made, then the forces are computed. The forces are globally summed, and the velocities and positions are updated. If necessary shake is performed to constrain bond lengths and/or bond angles. Temperature and Pressure can be controlled using weak coupling to a bath.

mdrun produces at least three output file, plus one log file (-g) per node. The trajectory file (-o), contains coordinates, velocities and optionally forces. The structure file (-c) contains the coordinates and velocities

of the last step. The energy file (-e) contains energies, the temperature, pressure, etc, a lot of these things are also printed in the log file of node 0. Optionally coordinates can be written to a compressed trajectory file (-x).

When running in parallel with PVM or an old version of MPI the `-np` option must be given to indicate the number of nodes.

The option `-dgd1` is only used when free energy perturbation is turned on.

With `-rerun` an input trajectory can be given for which forces and energies will be (re)calculated. Neighbor searching will be performed for every frame, unless `nstlist` is zero (see the `.mdp` file).

ED (essential dynamics) sampling is switched on by using the `-ei` flag followed by an `.edi` file. The `.edi` file can be produced using options in the `essdyn` menu of the WHAT IF program. `mdrun` produces a `.edo` file that contains projections of positions, velocities and forces onto selected eigenvectors.

When user-defined potential functions have been selected in the `.mdp` file the `-table` option is used to pass `mdrun` a formatted table with potential functions. The file is read from either the current directory or from the `GMXLIB` directory. A number of preformatted tables are presented in the `GMXLIB` dir, for 6-8, 6-9, 6-10, 6-11, 6-12 Lennard Jones potentials with normal Coulomb. When pair interactions are present a separate table for pair interaction functions is read using the `-tablep` option.

The options `-pi`, `-po`, `-pd`, `-pn` are used for potential of mean force calculations and umbrella sampling. See manual.

With `-multi` multiple systems are simulated in parallel. As many (single node) input files are required as the number of nodes. The node number is appended to the run input and each output filename, for instance `topol.tpr` becomes `topol0.tpr`, `topol1.tpr` etc. The main use of this option is for NMR refinement: when distance or orientation restraints are present these can be ensemble averaged over all the systems.

With `-replex` replica exchange is attempted every given number of steps. This option implies `-multi`, see above. All run input files should use a different coupling temperature, the order of the files is not important. The random seed is set with `-reseed`. The velocities are scaled and neighbor searching is performed after every exchange.

Finally some experimental algorithms can be tested when the appropriate options have been given. Currently under investigation are: polarizability, glass simulations and X-Ray bombardments.

When `mdrun` receives a `TERM` signal, it will set `nsteps` to the current step plus one. When `mdrun` receives a `USR1` signal, it will set `nsteps` to the next multiple of `nstxout` after the current step. In both cases all the usual output will be written to file. When running with MPI, a signal to one of the `mdrun` processes is sufficient, this signal should not be sent to `mpirun` or the `mdrun` process that is the parent of the others.

Files

<code>-s</code>	<code>topol.tpr</code>	Input	Generic run input: tpr tpb tpa xml
<code>-o</code>	<code>traj.trr</code>	Output	Full precision trajectory: trr trj
<code>-x</code>	<code>traj.xtc</code>	Output, Opt.	Compressed trajectory (portable xdr format)
<code>-c</code>	<code>confout.gro</code>	Output	Generic structure: gro g96 pdb xml
<code>-e</code>	<code>ener.edr</code>	Output	Generic energy: edr ene
<code>-g</code>	<code>md.log</code>	Output	Log file
<code>-dgd1</code>	<code>dgdl.xvg</code>	Output, Opt.	xvgr/xmgr file
<code>-field</code>	<code>field.xvg</code>	Output, Opt.	xvgr/xmgr file
<code>-table</code>	<code>table.xvg</code>	Input, Opt.	xvgr/xmgr file
<code>-tablep</code>	<code>tablep.xvg</code>	Input, Opt.	xvgr/xmgr file
<code>-rerun</code>	<code>rerun.xtc</code>	Input, Opt.	Generic trajectory: xtc trr trj gro g96 pdb
<code>-tpi</code>	<code>tpi.xvg</code>	Output, Opt.	xvgr/xmgr file
<code>-ei</code>	<code>sam.edi</code>	Input, Opt.	ED sampling input
<code>-eo</code>	<code>sam.edo</code>	Output, Opt.	ED sampling output
<code>-j</code>	<code>wham.gct</code>	Input, Opt.	General coupling stuff

-jo	bam.gct	Output, Opt.	General coupling stuff
-ffout	gct.xvg	Output, Opt.	xvgr/xmgr file
-devout	deviatie.xvg	Output, Opt.	xvgr/xmgr file
-runav	runaver.xvg	Output, Opt.	xvgr/xmgr file
-pi	pull.ppa	Input, Opt.	Pull parameters
-po	pullout.ppa	Output, Opt.	Pull parameters
-pd	pull.pdo	Output, Opt.	Pull data output
-pn	pull.ndx	Input, Opt.	Index file
-mtx	nm.mtx	Output, Opt.	Hessian matrix
-dn	dipole.ndx	Output, Opt.	Index file

Other options

-h	bool	no	Print help info and quit
-nice	int	19	Set the nicelevel
-deffnm	string		Set the default filename for all file options
-xvgr	bool	yes	Add specific codes (legends etc.) in the output xvg files for the xmgrace program
-np	int	1	Number of nodes, must be the same as used for grompp
-nt	int	1	Number of threads to start on each node
-v	bool	no	Be loud and noisy
-compact	bool	yes	Write a compact log file
-sepdvdl	bool	no	Write separate V and dVdl terms for each interaction type and node to the log file(s)
-multi	bool	no	Do multiple simulations in parallel (only with -np > 1)
-replex	int	0	Attempt replica exchange every # steps
-reseed	int	-1	Seed for replica exchange, -1 is generate a seed
-glas	bool	no	Do glass simulation with special long range corrections
-ionize	bool	no	Do a simulation including the effect of an X-Ray bombardment on your system

D.68 mk_angndx

mk_angndx makes an index file for calculation of angle distributions etc. It uses a run input file (.tpr) for the definitions of the angles, dihedrals etc.

Files

-s	topol.tpr	Input	Generic run input: tpr tpb tpa xml
-n	angle.ndx	Output	Index file

Other options

-h	bool	no	Print help info and quit
-nice	int	0	Set the nicelevel
-type	enum	angle	Type of angle: angle, g96-angle, dihedral, improper, ryckaert-bellemans or phi-psi

D.69 ngmx

ngmx is the Gromacs trajectory viewer. This program reads a trajectory file, a run input file and an index file and plots a 3D structure of your molecule on your standard X Window screen. No need for a high end graphics workstation, it even works on Monochrome screens.

The following features have been implemented: 3D view, rotation, translation and scaling of your molecule(s), labels on atoms, animation of trajectories, hardcopy in PostScript format, user defined atom-filters runs on MIT-X (real X), open windows and motif, user friendly menus, option to remove periodicity, option to show computational box.

Some of the more common X command line options can be used:

-bg, -fg change colors, -font fontname, changes the font.

Files

-f	traj.xtc	Input	Generic trajectory: xtc trr trj gro g96 pdb
-s	topol.tpr	Input	Generic run input: tpr tpb tpa xml
-n	index.ndx	Input, Opt.	Index file

Other options

-h	bool	no	Print help info and quit
-nice	int	0	Set the nicelevel
-b	time	0	First frame (ps) to read from trajectory
-e	time	0	Last frame (ps) to read from trajectory
-dt	time	0	Only use frame when $t \text{ MOD } dt = \text{first time (ps)}$

- Balls option does not work
- Some times dumps core without a good reason

D.70 *pdb2gmx*

This program reads a pdb file, reads some database files, adds hydrogens to the molecules and generates coordinates in Gromos (Gromos) format and a topology in Gromacs format. These files can subsequently be processed to generate a run input file.

The force fields in the distribution are currently:

oplsaa OPLS-AA/L all-atom force field (2001 aminoacid dihedrals)
 G43b1 GROMOS96 43b1 Vacuum Forcefield
 G43a1 GROMOS96 43a1 Forcefield
 G43a2 GROMOS96 43a2 Forcefield (improved alkane dihedrals)
 G45a3 GROMOS96 45a3 Forcefield
 G53a5 GROMOS96 53a5 Forcefield
 G53a6 GROMOS96 53a6 Forcefield
 gmX Gromacs Forcefield (a modified GROMOS87, see manual)
 encads Encad all-atom force field, using scaled-down vacuum charges
 encadv Encad all-atom force field, using full solvent charges

The corresponding data files can be found in the library directory with names like ffXXXX.YYY. Check chapter 5 of the manual for more information about file formats. By default the forcefield selection is interactive, but you can use the `-ff` option to specify one of the short names above on the command line instead. In that case *pdb2gmx* just looks for the corresponding file.

Note that a pdb file is nothing more than a file format, and it need not necessarily contain a protein structure. Every kind of molecule for which there is support in the database can be converted. If there is no support in the database, you can add it yourself.

The program has limited intelligence, it reads a number of database files, that allow it to make special bonds (Cys-Cys, Heme-His, etc.), if necessary this can be done manually. The program can prompt the user to select which kind of LYS, ASP, GLU, CYS or HIS residue she wants. For LYS the choice is

between LYS (two protons on NZ) or LYSH (three protons, default), for ASP and GLU unprotonated (default) or protonated, for HIS the proton can be either on ND1 (HISA), on NE2 (HISB) or on both (HISH). By default these selections are done automatically. For His, this is based on an optimal hydrogen bonding conformation. Hydrogen bonds are defined based on a simple geometric criterium, specified by the maximum hydrogen-donor-acceptor angle and donor-acceptor distance, which are set by `-angle` and `-dist` respectively.

Option `-merge` will ask if you want to merge consecutive chains into one molecule definition, this can be useful for connecting chains with a disulfide bridge or intermolecular distance restraints.

`pdb2gmx` will also check the occupancy field of the `pdb` file. If any of the occupancies are not one, indicating that the atom is not resolved well in the structure, a warning message is issued. When a `pdb` file does not originate from an X-Ray structure determination all occupancy fields may be zero. Either way, it is up to the user to verify the correctness of the input data (read the article!).

During processing the atoms will be reordered according to Gromacs conventions. With `-n` an index file can be generated that contains one group reordered in the same way. This allows you to convert a Gromos trajectory and coordinate file to Gromos. There is one limitation: reordering is done after the hydrogens are stripped from the input and before new hydrogens are added. This means that you should not use `-ignh`.

The `.gro` and `.g96` file formats do not support chain identifiers. Therefore it is useful to enter a `pdb` file name at the `-o` option when you want to convert a multichain `pdb` file.

The option `-vsite` removes hydrogen and fast improper dihedral motions. Angular and out-of-plane motions can be removed by changing hydrogens into virtual sites and fixing angles, which fixes their position relative to neighboring atoms. Additionally, all atoms in the aromatic rings of the standard amino acids (i.e. PHE, TRP, TYR and HIS) can be converted into virtual sites, eliminating the fast improper dihedral fluctuations in these rings. Note that in this case all other hydrogen atoms are also converted to virtual sites. The mass of all atoms that are converted into virtual sites, is added to the heavy atoms.

Also slowing down of dihedral motion can be done with `-heavyh` done by increasing the hydrogen-mass by a factor of 4. This is also done for water hydrogens to slow down the rotational motion of water. The increase in mass of the hydrogens is subtracted from the bonded (heavy) atom so that the total mass of the system remains the same.

Files

<code>-f</code>	<code>eiwit.pdb</code>	Input	Generic structure: <code>gro g96 pdb tpr tpb tpa xml</code>
<code>-o</code>	<code>conf.gro</code>	Output	Generic structure: <code>gro g96 pdb xml</code>
<code>-p</code>	<code>topol.top</code>	Output	Topology file
<code>-i</code>	<code>posre.itp</code>	Output	Include file for topology
<code>-n</code>	<code>clean.ndx</code>	Output, Opt.	Index file
<code>-q</code>	<code>clean.pdb</code>	Output, Opt.	Generic structure: <code>gro g96 pdb xml</code>

Other options

<code>-h</code>	bool	no	Print help info and quit
<code>-nice</code>	int	0	Set the nicelevel
<code>-merge</code>	bool	no	Merge chains into one molecule definition
<code>-ff</code>	string	select	Force field, interactive by default. Use <code>-h</code> for information.
<code>-water</code>	enum	<code>spc</code>	Water model to use: with GROMOS we recommend SPC, with OPLS, TIP4P: <code>spc, spce, tip3p, tip4p, tip5p</code> or <code>f3c</code>
<code>-inter</code>	bool	no	Set the next 6 options to interactive
<code>-ss</code>	bool	no	Interactive SS bridge selection
<code>-ter</code>	bool	no	Interactive termini selection, iso charged
<code>-lys</code>	bool	no	Interactive Lysine selection, iso charged
<code>-asp</code>	bool	no	Interactive Aspartic Acid selection, iso charged
<code>-glu</code>	bool	no	Interactive Glutamic Acid selection, iso charged
<code>-his</code>	bool	no	Interactive Histidine selection, iso checking H-bonds

-angle	real	135	Minimum hydrogen-donor-acceptor angle for a H-bond (degrees)
-dist	real	0.3	Maximum donor-acceptor distance for a H-bond (nm)
-una	bool	no	Select aromatic rings with united CH atoms on Phenylalanine, Tryptophane and Tyrosine
-ignh	bool	no	Ignore hydrogen atoms that are in the pdb file
-missing	bool	no	Continue when atoms are missing, dangerous
-v	bool	no	Be slightly more verbose in messages
-posrefc	real	1000	Force constant for position restraints
-vsite	enum	none	Convert atoms to virtual sites: none, hydrogens or aromatics
-heavyh	bool	no	Make hydrogen atoms heavy
-deuterate	bool	no	Change the mass of hydrogens to 2 amu

D.71 protonate

`protonate` reads (a) conformation(s) and adds all missing hydrogens as defined in `ffgm2.hdb`. If only `-s` is specified, this conformation will be protonated, if also `-f` is specified, the conformation(s) will be read from this file which can be either a single conformation or a trajectory.

If a `pdb` file is supplied, residue names might not correspond to the GROMACS naming conventions, in which case these residues will probably not be properly protonated.

If an index file is specified, please note that the atom numbers should correspond to the **protonated** state.

Files

-s	<code>topol.tpr</code>	Input	Structure+mass(db): <code>tpr tpb tpa gro g96 pdb xml</code>
-f	<code>traj.xtc</code>	Input, Opt.	Generic trajectory: <code>xtc trr trj gro g96 pdb</code>
-n	<code>index.ndx</code>	Input, Opt.	Index file
-o	<code>protonated.xtc</code>	Output	Generic trajectory: <code>xtc trr trj gro g96 pdb</code>

Other options

-h	bool	no	Print help info and quit
-nice	int	0	Set the nicelevel
-b	time	0	First frame (ps) to read from trajectory
-e	time	0	Last frame (ps) to read from trajectory
-dt	time	0	Only use frame when $t \text{ MOD } dt = \text{first time (ps)}$

D.72 tpbconv

`tpbconv` can edit run input files in three ways.

1st. by creating a run input file for a continuation run when your simulation has crashed due to e.g. a full disk, or by making a continuation run input file. Note that a frame with coordinates and velocities is needed, which means that when you never write velocities, you can not use `tpbconv` and you have to start the run again from the beginning. When pressure and/or Nose-Hoover temperature coupling is used an energy file can be supplied to get an exact continuation of the original run.

2nd. by creating a `tpx` file for a subset of your original `tpx` file, which is useful when you want to remove the solvent from your `tpx` file, or when you want to make e.g. a pure Ca `tpx` file. **WARNING: this `tpx` file is not fully functional.** **3rd.** by setting the charges of a specified group to zero. This is useful when doing free energy estimates using the LIE (Linear Interaction Energy) method.

Files

-s	topol.tpr	Input	Generic run input: tpr tpb tpa xml
-f	traj.trr	Input, Opt.	Full precision trajectory: trr trj
-e	ener.edr	Input, Opt.	Generic energy: edr ene
-n	index.ndx	Input, Opt.	Index file
-o	tpxout.tpr	Output	Generic run input: tpr tpb tpa xml

Other options

-h	bool	no	Print help info and quit
-nice	int	0	Set the nicelevel
-time	real	-1	Continue from frame at this time (ps) instead of the last frame
-extend	real	0	Extend runtime by this amount (ps)
-until	real	0	Extend runtime until this ending time (ps)
-zeroq	bool	no	Set the charges of a group (from the index) to zero
-unconstrained	bool	yes	For a continuous trajectory, the constraints should not be solved before the first step (default)

D.73 trjcat

trjcat concatenates several input trajectory files in sorted order. In case of double time frames the one in the later file is used. By specifying `-settime` you will be asked for the start time of each file. The input files are taken from the command line, such that a command like `trjcat -o fixed.trr *.trr` should do the trick. Using `-cat` you can simply paste several files together without removal of frames with identical time stamps.

One important option is inferred when the output file is amongst the input files. In that case that particular file will be appended to which implies you do not need to store double the amount of data. Obviously the file to append to has to be the one with lowest starting time since one can only append at the end of a file.

If the `-demux` option is given, the N trajectories that are read, are written in another order as specified in the xvg file. The xvg file should contain something like:

```
0 0 1 2 3 4 5
2 1 0 2 3 5 4
```

Where the first number is the time, and subsequent numbers point to trajectory indices. The frames corresponding to the numbers present at the first line are collected into the output trajectory. If the number of frames in the trajectory does not match that in the xvg file then the program tries to be smart. Beware.

Files

-f	traj.xtc	Input, Mult.	Generic trajectory: xtc trr trj gro g96 pdb
-o	trajout.xtc	Output, Mult.	Generic trajectory: xtc trr trj gro g96 pdb
-n	index.ndx	Input, Opt.	Index file
-demux	remd.xvg	Input, Opt.	xvgr/xmgr file

Other options

-h	bool	no	Print help info and quit
-nice	int	19	Set the nicelevel
-tu	enum	ps	Time unit: ps, fs, ns, us, ms or s
-xvgr	bool	yes	Add specific codes (legends etc.) in the output xvg files for the xmgrace program
-b	time	-1	First time to use (ps)
-e	time	-1	Last time to use (ps)
-dt	time	0	Only write frame when $t \text{ MOD } dt = \text{first time (ps)}$
-prec	int	3	Precision for .xtc and .gro writing in number of decimal places

<code>-vel</code>	bool	yes	Read and write velocities if possible
<code>-settime</code>	bool	no	Change starting time interactively
<code>-sort</code>	bool	yes	Sort trajectory files (not frames)
<code>-keepplast</code>	bool	no	keep overlapping frames at end of trajectory
<code>-cat</code>	bool	no	do not discard double time frames

D.74 *trjconv*

trjconv can convert trajectory files in many ways:

1. from one format to another
2. select a subset of atoms
3. remove periodicity from molecules
4. keep multimeric molecules together
5. center atoms in the box
6. fit atoms to reference structure
7. reduce the number of frames
8. change the timestamps of the frames (`-t0` and `-timestep`)
9. cut the trajectory in small subtrajectories according to information in an index file. This allows subsequent analysis of the subtrajectories that could, for example be the result of a cluster analysis. Use option `-sub`. This assumes that the entries in the index file are frame numbers and dumps each group in the index file to a separate trajectory file.
10. select frames within a certain range of a quantity given in an `.xvg` file.

The program *trjcat* can concatenate multiple trajectory files.

Currently seven formats are supported for input and output: `.xtc`, `.trr`, `.trj`, `.gro`, `.g96`, `.pdb` and `.g87`. The file formats are detected from the file extension. The precision of `.xtc` and `.gro` output is taken from the input file for `.xtc`, `.gro` and `.pdb`, and from the `-ndec` option for other input formats. The precision is always taken from `-ndec`, when this option is set. All other formats have fixed precision. `.trr` and `.trj` output can be single or double precision, depending on the precision of the *trjconv* binary. Note that velocities are only supported in `.trr`, `.trj`, `.gro` and `.g96` files.

Option `-app` can be used to append output to an existing trajectory file. No checks are performed to ensure integrity of the resulting combined trajectory file.

Option `-sep` can be used to write every frame to a separate `.gro`, `.g96` or `.pdb` file, default all frames all written to one file. `.pdb` files with all frames concatenated can be viewed with `rasmol -nmrpdb`.

It is possible to select part of your trajectory and write it out to a new trajectory file in order to save disk space, e.g. for leaving out the water from a trajectory of a protein in water. **ALWAYS** put the original trajectory on tape! We recommend to use the portable `.xtc` format for your analysis to save disk space and to have portable files.

There are two options for fitting the trajectory to a reference either for essential dynamics analysis or for whatever. The first option is just plain fitting to a reference structure in the structure file, the second option is a progressive fit in which the first timeframe is fitted to the reference structure in the structure file to obtain and each subsequent timeframe is fitted to the previously fitted structure. This way a continuous trajectory is generated, which might not be the case when using the regular fit method, e.g. when your protein undergoes large conformational transitions.

Option `-pbc` sets the type of periodic boundary condition treatment:

`whole` puts the atoms in the box and then makes broken molecules whole (a run input file is required). Atom number 1 of each molecule will be inside the box.

`com` puts the center of mass of all *residues* in the box. Not that this can break molecules that consist of

more than one residue (e.g. proteins).

`inbox` puts all the atoms in the box.

`nojump` checks if atoms jump across the box and then puts them back. This has the effect that all molecules will remain whole (provided they were whole in the initial conformation), note that this ensures a continuous trajectory but molecules may diffuse out of the box. The starting configuration for this procedure is taken from the structure file, if one is supplied, otherwise it is the first frame.

`cluster` clusters all the atoms in the selected index such that they are all closest to the center of mass of the cluster which is iteratively updated. Note that this will only give meaningful results if you in fact have a cluster. Luckily that can be checked afterwards using a trajectory viewer.

`-pbc` is ignored when `-fit` or `-pfit` is set, in that case molecules will be made whole.

Option `-ur` sets the unit cell representation for options `whole` and `inbox` of `-pbc`. All three options give different results for triclinic boxes and identical results for rectangular boxes. `rect` is the ordinary brick shape. `tric` is the triclinic unit cell. `compact` puts all atoms at the closest distance from the center of the box. This can be useful for visualizing e.g. truncated octahedrons. The center for options `tric` and `compact` is `tric` (see below), unless the option `-center` is set differently.

Option `-center` centers the system in the box. The user can select the group which is used to determine the geometrical center. The center options are: `tric`: half of the sum of the box vectors, `rect`: half of the box diagonal, `zero`: zero. Use option `-pbc whole` in addition to `-center` when you want all molecules in the box after the centering.

With `-dt` it is possible to reduce the number of frames in the output. This option relies on the accuracy of the times in your input trajectory, so if these are inaccurate use the `-timestep` option to modify the time (this can be done simultaneously). For making smooth movies the program `g_filter` can reduce the number of frames while using low-pass frequency filtering, this reduces aliasing of high frequency motions.

Using `-trunc trjconv` can truncate `.trj` in place, i.e. without copying the file. This is useful when a run has crashed during disk I/O (one more disk full), or when two contiguous trajectories must be concatenated without have double frames.

`trjcat` is more suitable for concatenating trajectory files.

Option `-dump` can be used to extract a frame at or near one specific time from your trajectory.

Option `-drop` reads an `.xvg` file with times and values. When options `-dropunder` and/or `-droptover` are set, frames with a value below and above the value of the respective options will not be written.

Files

<code>-f</code>	<code>traj.xtc</code>	Input	Generic trajectory: <code>xtc trr trj gro g96 pdb</code>
<code>-o</code>	<code>trajout.xtc</code>	Output	Generic trajectory: <code>xtc trr trj gro g96 pdb</code>
<code>-s</code>	<code>topol.tpr</code>	Input, Opt.	Structure+mass(db): <code>tpr tpb tpa gro g96 pdb xml</code>
<code>-n</code>	<code>index.ndx</code>	Input, Opt.	Index file
<code>-fr</code>	<code>frames.ndx</code>	Input, Opt.	Index file
<code>-sub</code>	<code>cluster.ndx</code>	Input, Opt.	Index file
<code>-drop</code>	<code>drop.xvg</code>	Input, Opt.	<code>xvgr/xmgr</code> file

Other options

<code>-h</code>	<code>bool</code>	<code>no</code>	Print help info and quit
<code>-nice</code>	<code>int</code>	<code>19</code>	Set the nicelevel
<code>-b</code>	<code>time</code>	<code>0</code>	First frame (ps) to read from trajectory
<code>-e</code>	<code>time</code>	<code>0</code>	Last frame (ps) to read from trajectory
<code>-tu</code>	<code>enum</code>	<code>ps</code>	Time unit: <code>ps, fs, ns, us, ms</code> or <code>s</code>
<code>-w</code>	<code>bool</code>	<code>no</code>	View output <code>xvg, xpm, eps</code> and <code>pdb</code> files
<code>-xvgr</code>	<code>bool</code>	<code>yes</code>	Add specific codes (legends etc.) in the output <code>xvg</code> files for the <code>xmgrace</code> program
<code>-skip</code>	<code>int</code>	<code>1</code>	Only write every <code>nr</code> -th frame

-dt	time	0	Only write frame when $t \text{ MOD } dt = \text{first time (ps)}$
-dump	time	-1	Dump frame nearest specified time (ps)
-t0	time	0	Starting time (ps) (default: don't change)
-timestep	time	0	Change time step between input frames (ps)
-pbc	enum	none	PBC treatment (see help text for full description): none, whole, inbox, nojump, cluster or com
-ur	enum	rect	Unit-cell representation: rect, tric or compact
-center	enum	no	Center atoms in box: no, tric, rect or zero
-box	vector	0 0 0	Size for new cubic box (default: read from input)
-shift	vector	0 0 0	All coordinates will be shifted by $\text{framenr} * \text{shift}$
-fit	enum	none	Fit molecule to ref structure in the structure file: none, rot+trans, translation or progressive
-ndec	int	3	Precision for .xtc and .gro writing in number of decimal places
-vel	bool	yes	Read and write velocities if possible
-force	bool	no	Read and write forces if possible
-trunc	time	-1	Truncate input trj file after this time (ps)
-exec	string		Execute command for every output frame with the frame number as argument
-app	bool	no	Append output
-split	time	0	Start writing new file when $t \text{ MOD } \text{split} = \text{first time (ps)}$
-sep	bool	no	Write each frame to a separate .gro, .g96 or .pdb file
-ter	bool	no	Use 'TER' in pdb file as end of frame in stead of default 'ENDMDL'
-dropunder	real	0	Drop all frames below this value
-dropover	real	0	Drop all frames above this value

D.75 trjorder

trjorder orders molecules according to the smallest distance to atoms in a reference group. It will ask for a group of reference atoms and a group of molecules. For each frame of the trajectory the selected molecules will be reordered according to the shortest distance between atom number `-da` in the molecule and all the atoms in the reference group. All atoms in the trajectory are written to the output trajectory.

trjorder can be useful for e.g. analyzing the n waters closest to a protein. In that case the reference group would be the protein and the group of molecules would consist of all the water atoms. When an index group of the first n waters is made, the ordered trajectory can be used with any Gromacs program to analyze the n closest waters.

If the output file is a pdb file, the distance to the reference target will be stored in the B-factor field in order to color with e.g. rasmol.

Files

-f	traj.xtc	Input	Generic trajectory: xtc trr trj gro g96 pdb
-s	topol.tpr	Input	Structure+mass(db): tpr tpb tpa gro g96 pdb xml
-n	index.ndx	Input, Opt.	Index file
-o	ordered.xtc	Output	Generic trajectory: xtc trr trj gro g96 pdb
-nshell	nshell.xvg	Output, Opt.	xvgr/xmgr file

Other options

-h	bool	no	Print help info and quit
-nice	int	19	Set the nicelevel
-b	time	0	First frame (ps) to read from trajectory
-e	time	0	Last frame (ps) to read from trajectory
-dt	time	0	Only use frame when $t \text{ MOD } dt = \text{first time (ps)}$

<code>-xvgr</code>	<code>bool</code>	<code>yes</code>	Add specific codes (legends etc.) in the output xvg files for the xmgrace program
<code>-na</code>	<code>int</code>	<code>3</code>	Number of atoms in a molecule
<code>-da</code>	<code>int</code>	<code>1</code>	Atom used for the distance calculation
<code>-com</code>	<code>bool</code>	<code>no</code>	Use the distance to the center of mass of the reference group
<code>-r</code>	<code>real</code>	<code>0</code>	Cutoff used for the distance calculation when computing the number of molecules in a shell around e.g. a protein

D.76 wheel

wheel plots a helical wheel representation of your sequence. The input sequence is in the .dat file where the first line contains the number of residues and each consecutive line contains a residuename.

Files

<code>-f</code>	<code>nnnice.dat</code>	Input	Generic data file
<code>-o</code>	<code>plot.eps</code>	Output	Encapsulated PostScript (tm) file

Other options

<code>-h</code>	<code>bool</code>	<code>no</code>	Print help info and quit
<code>-nice</code>	<code>int</code>	<code>19</code>	Set the nicelevel
<code>-r0</code>	<code>int</code>	<code>1</code>	The first residue number in the sequence
<code>-rot0</code>	<code>real</code>	<code>0</code>	Rotate around an angle initially (90 degrees makes sense)
<code>-T</code>	<code>string</code>		Plot a title in the center of the wheel (must be shorter than 10 characters, or it will overwrite the wheel)
<code>-nn</code>	<code>bool</code>	<code>yes</code>	Toggle numbers

D.77 x2top

x2top generates a primitive topology from a coordinate file. The program assumes all hydrogens are present when defining the hybridization from the atom name and the number of bonds. The program can also make an rtp entry, which you can then add to the rtp database.

When `-param` is set, equilibrium distances and angles and force constants will be printed in the topology for all interactions. The equilibrium distances and angles are taken from the input coordinates, the force constant are set with command line options. The force fields supported currently are:

G43a1 GROMOS96 43a1 Forcefield (official distribution)

oplsaa OPLS-AA/L all-atom force field (2001 aminoacid dihedrals)

G43b1 GROMOS96 43b1 Vacuum Forcefield (official distribution)

gmx Gromacs Forcefield (a modified GROMOS87, see manual)

G43a2 GROMOS96 43a2 Forcefield (development) (improved alkane dihedrals)

The corresponding data files can be found in the library directory with names like `ffXXXX.YYY`. Check chapter 5 of the manual for more information about file formats. By default the forcefield selection is interactive, but you can use the `-ff` option to specify one of the short names above on the command line instead. In that case `pdb2gmx` just looks for the corresponding file.

Files

<code>-f</code>	<code>conf.gro</code>	Input	Generic structure: gro g96 pdb tpr tpb tpa xml
<code>-o</code>	<code>out.top</code>	Output, Opt.	Topology file
<code>-r</code>	<code>out.rtp</code>	Output, Opt.	Residue Type file used by <code>pdb2gmx</code>

Other options

-h	bool	no	Print help info and quit
-nice	int	0	Set the nicelevel
-scale	real	1.1	Scaling factor for bonds with unknown atom types relative to atom type O
-ff	string	select	Select the force field for your simulation.
-nexcl	int	3	Number of exclusions
-H14	bool	yes	Use 3rd neighbour interactions for hydrogen atoms
-alldih	bool	no	Generate all proper dihedrals
-remdih	bool	no	Remove dihedrals on the same bond as an improper
-pairs	bool	yes	Output 1-4 interactions (pairs) in topology file
-name	string	ICE	Name of your molecule
-pbc	bool	yes	Use periodic boundary conditions.
-param	bool	no	Print parameters in the output
-round	bool	yes	Round off measured values
-kb	real	400000	Bonded force constant (kJ/mol/nm ²)
-kt	real	400	Angle force constant (kJ/mol/rad ²)
-kp	real	5	Dihedral angle force constant (kJ/mol/rad ²)

- The atom type selection is primitive. Virtually no chemical knowledge is used
- Periodic boundary conditions screw up the bonding
- No improper dihedrals are generated
- The atoms to atomtype translation table is incomplete (ffG43a1.n2t file in the \$GMXLIB directory). Please extend it and send the results back to the GROMACS crew.

D.78 xpm2ps

xpm2ps makes a beautiful color plot of an XPixelMap file. Labels and axis can be displayed, when they are supplied in the correct matrix format. Matrix data may be generated by programs such as do_dssp, g_rms or g_mdmat.

Parameters are set in the m2p file optionally supplied with -di. Reasonable defaults are provided. Settings for the y-axis default to those for the x-axis. Font names have a defaulting hierarchy: titlefont -> legendfont; titlefont -> (xfont -> yfont -> ytickfont) -> xtickfont, e.g. setting titlefont sets all fonts, setting xfont sets yfont, ytickfont and xtickfont.

When no m2p file is supplied, many settings are set by command line options. The most important option is -size which sets the size of the whole matrix in postscript units. This option can be overridden with the -bx and -by options (and the corresponding parameters in the m2p file), which set the size of a single matrix element.

With -f2 a 2nd matrix file can be supplied, both matrix files will be read simultaneously and the upper left half of the first one (-f) is plotted together with the lower right half of the second one (-f2). The diagonal will contain values from the matrix file selected with -diag. Plotting of the diagonal values can be suppressed altogether by setting -diag to none. In this case, a new color map will be generated with a red gradient for negative numbers and a blue for positive. If the color coding and legend labels of both matrices are identical, only one legend will be displayed, else two separate legends are displayed. With -combine an alternative operation can be selected to combine the matrices. The output range is automatically set to the actual range of the combined matrix. This can be overridden with -cmin and -cmax.

`-title` can be set to `none` to suppress the title, or to `ylabel` to show the title in the Y-label position (alongside the Y-axis).

With the `-rainbow` option dull grey-scale matrices can be turned into attractive color pictures.

Merged or rainbowed matrices can be written to an XPixmap file with the `-xpm` option.

Files

<code>-f</code>	<code>root.xpm</code>	Input	X Pixmap compatible matrix file
<code>-f2</code>	<code>root2.xpm</code>	Input, Opt.	X Pixmap compatible matrix file
<code>-di</code>	<code>ps.m2p</code>	Input, Opt., Link	Input file for mat2ps
<code>-do</code>	<code>out.m2p</code>	Output, Opt.	Input file for mat2ps
<code>-o</code>	<code>plot.eps</code>	Output, Opt.	Encapsulated PostScript (tm) file
<code>-xpm</code>	<code>root.xpm</code>	Output, Opt.	X Pixmap compatible matrix file

Other options

<code>-h</code>	bool	no	Print help info and quit
<code>-nice</code>	int	0	Set the nicelevel
<code>-w</code>	bool	no	View output xvg, xpm, eps and pdb files
<code>-frame</code>	bool	yes	Display frame, ticks, labels, title and legend
<code>-title</code>	enum	top	Show title at: top, once, ylabel or none
<code>-yonce</code>	bool	no	Show y-label only once
<code>-legend</code>	enum	both	Show legend: both, first, second or none
<code>-diag</code>	enum	first	Diagonal: first, second or none
<code>-size</code>	real	400	Horizontal size of the matrix in ps units
<code>-bx</code>	real	0	Element x-size, overrides -size (also y-size when -by is not set)
<code>-by</code>	real	0	Element y-size
<code>-rainbow</code>	enum	no	Rainbow colors, convert white to: no, blue or red
<code>-gradient</code>	vector	0 0 0	Re-scale colormap to a smooth gradient from white 1,1,1 to r,g,b
<code>-skip</code>	int	1	only write out every nr-th row and column
<code>-zeroline</code>	bool	no	insert line in xpm matrix where axis label is zero
<code>-legoffset</code>	int	0	Skip first N colors from xpm file for the legend
<code>-combine</code>	enum	halves	Combine two matrices: halves, add, sub, mult or div
<code>-cmin</code>	real	0	Minimum for combination output
<code>-cmax</code>	real	0	Maximum for combination output

D.79 xrama

`xrama` shows a Ramachandran movie, that is, it shows the Phi/Psi angles as a function of time in an X-Window.

Static Phi/Psi plots for printing can be made with `g_rama`.

Some of the more common X command line options can be used:

`-bg`, `-fg` change colors, `-font` fontname, changes the font.

Files

<code>-f</code>	<code>traj.xtc</code>	Input	Generic trajectory: xtc trr trj gro g96 pdb
<code>-s</code>	<code>topol.tpr</code>	Input	Generic run input: tpr tpb tpa xml

Other options

<code>-h</code>	bool	no	Print help info and quit
<code>-nice</code>	int	0	Set the nicelevel
<code>-b</code>	time	0	First frame (ps) to read from trajectory
<code>-e</code>	time	0	Last frame (ps) to read from trajectory
<code>-dt</code>	time	0	Only use frame when t MOD dt = first time (ps)

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